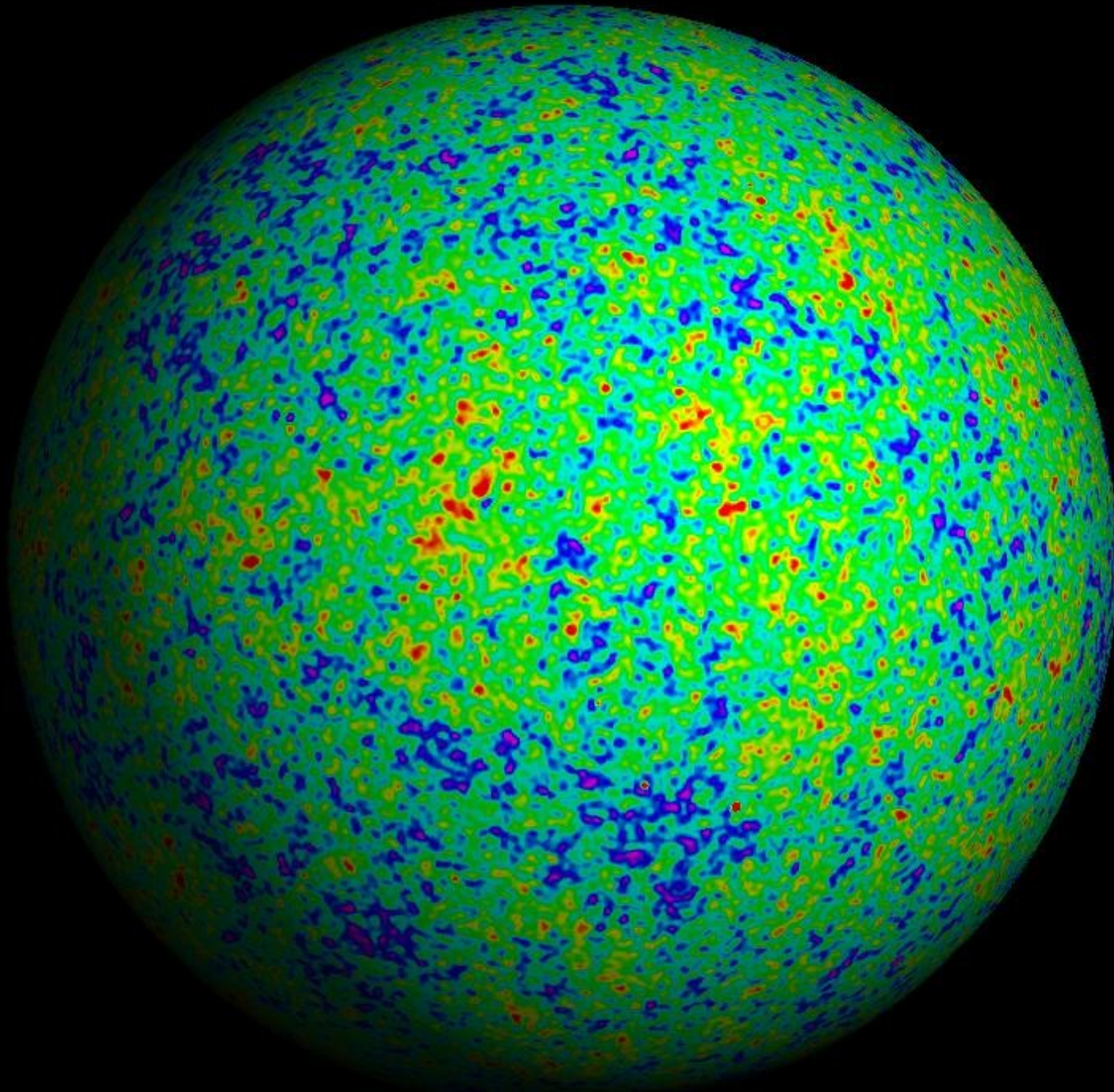


Curvature couplings of the Inflaton

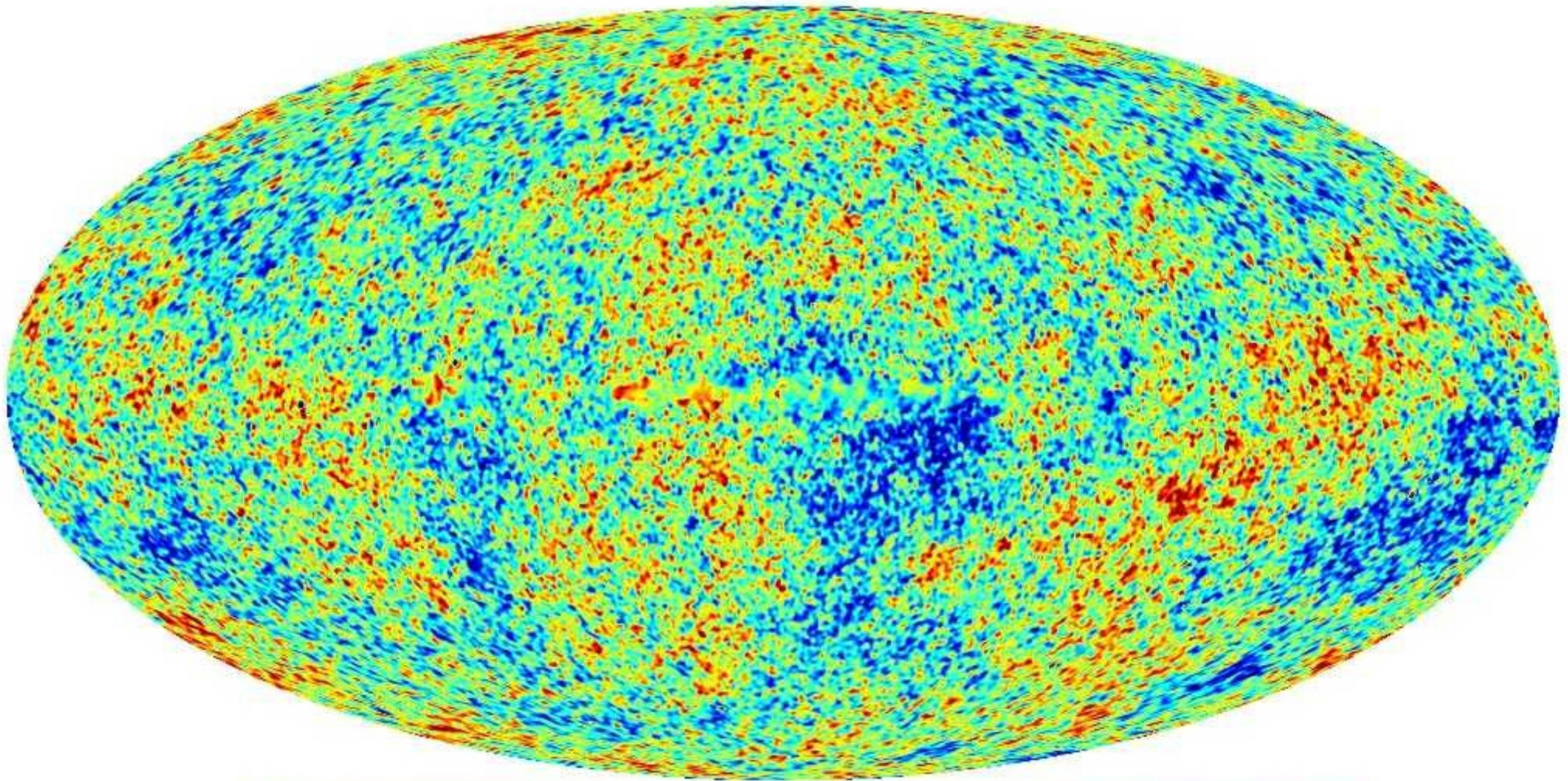
Subhendra Mohanty

Physical Research Laboratory,

Ahmedabad



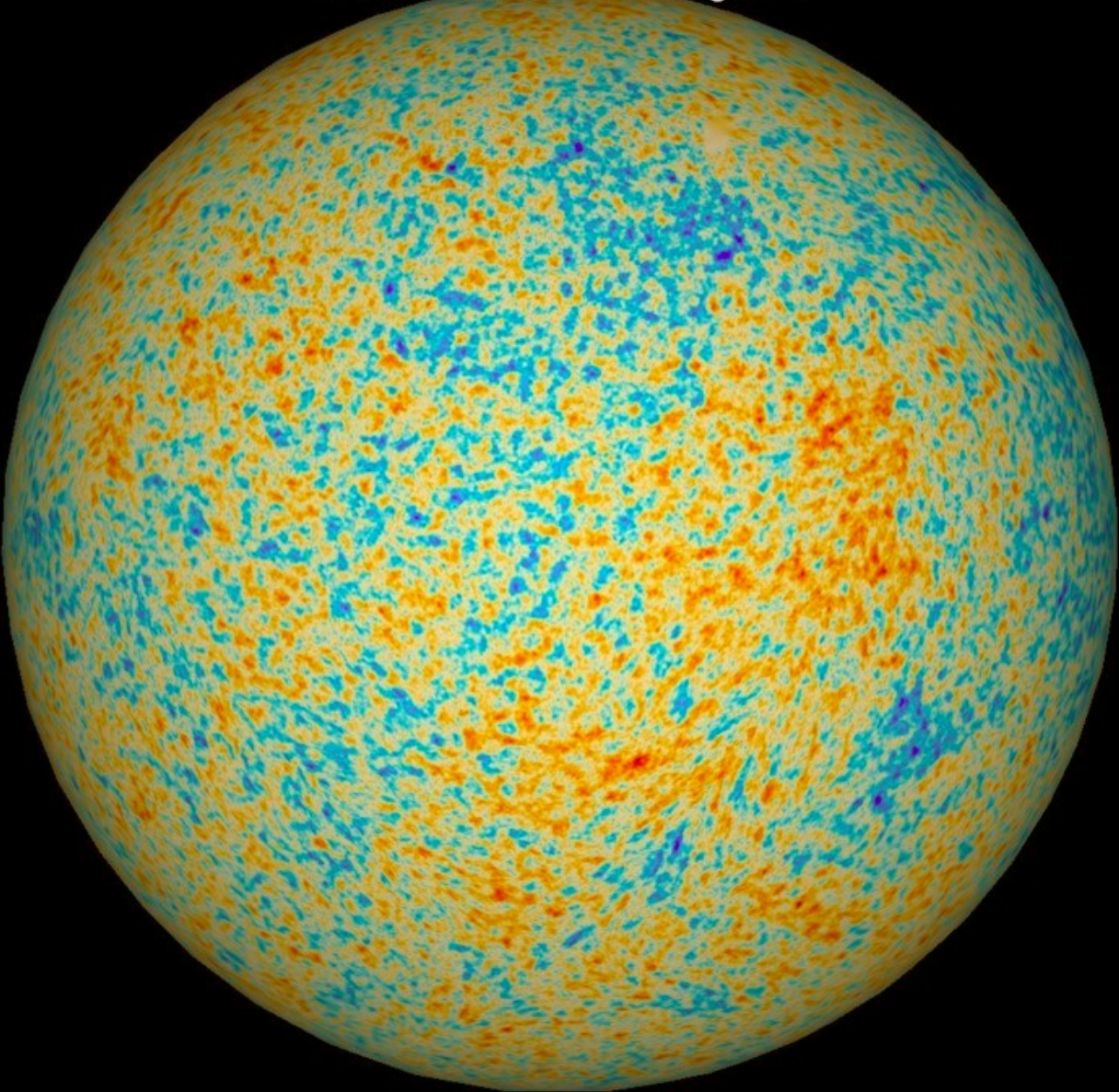
CMB Anisotropy
WMAP



WMAP 3 year

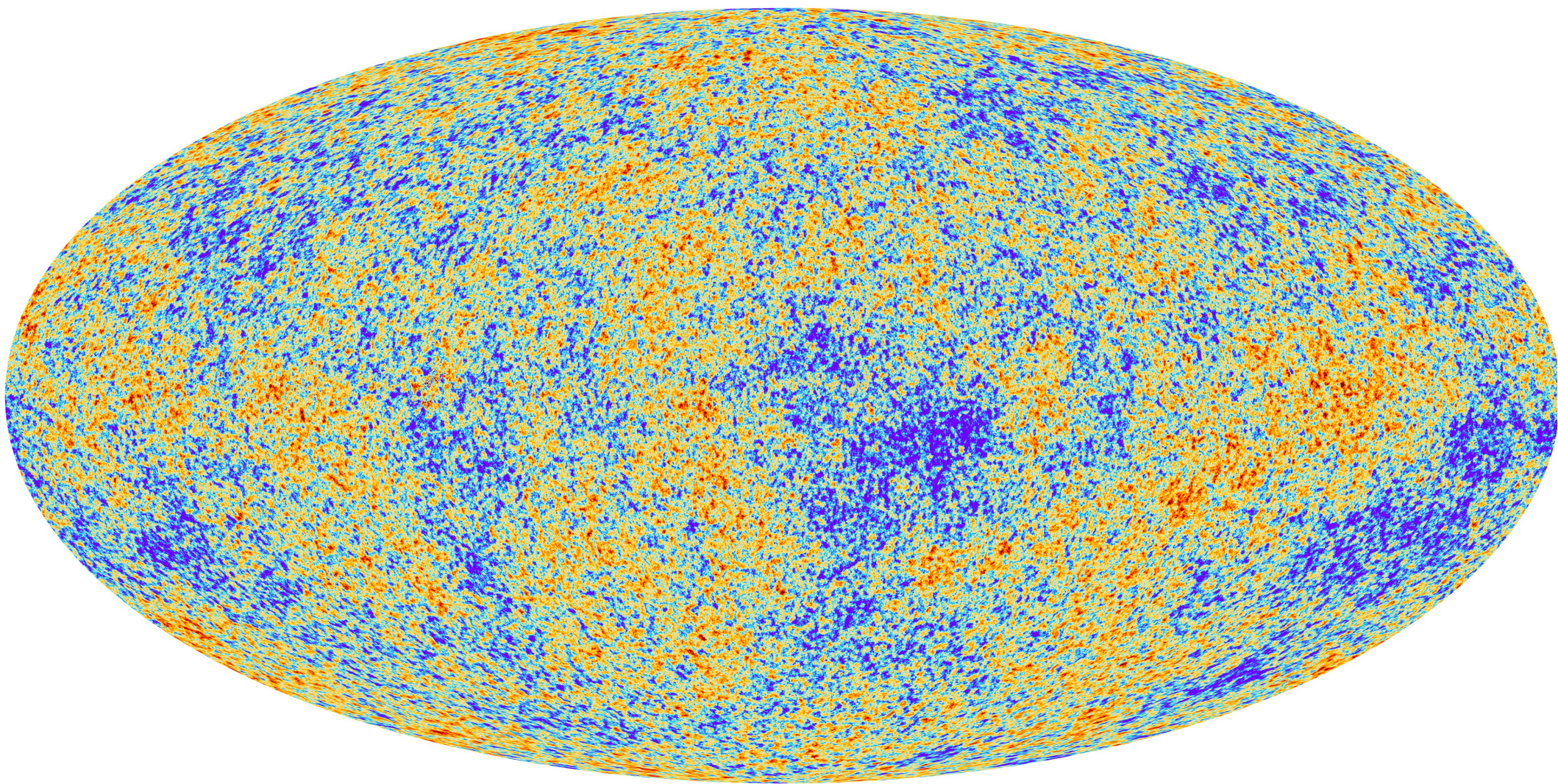
The cosmic microwave background

261.71°, -55.67° galactic

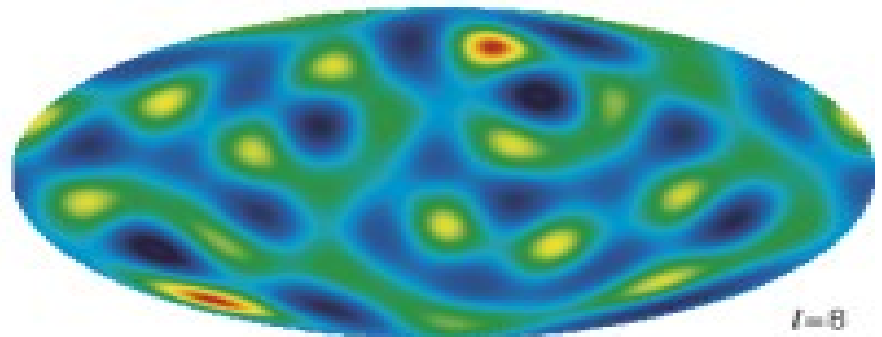
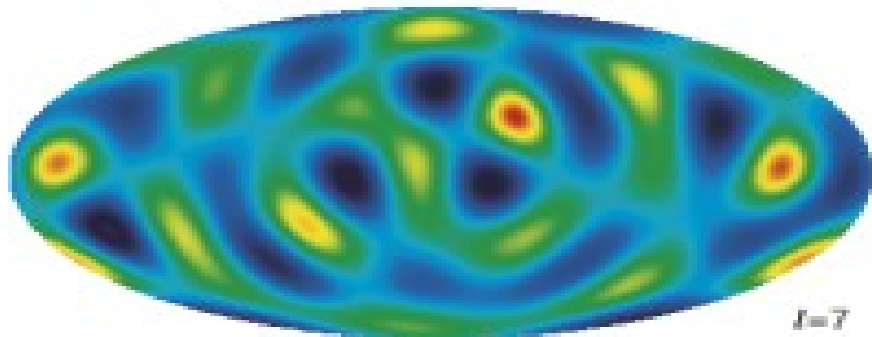
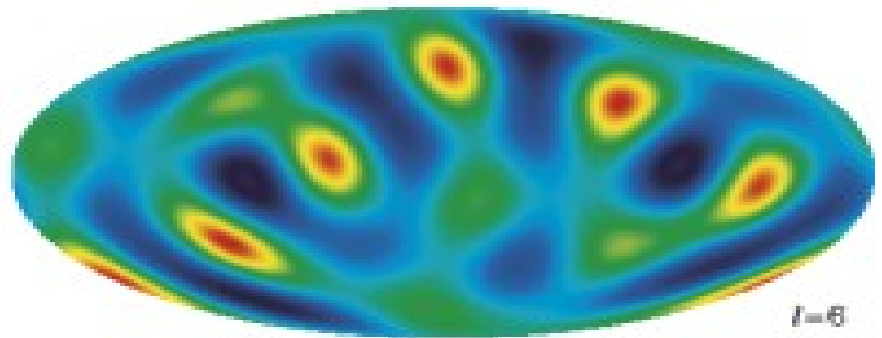
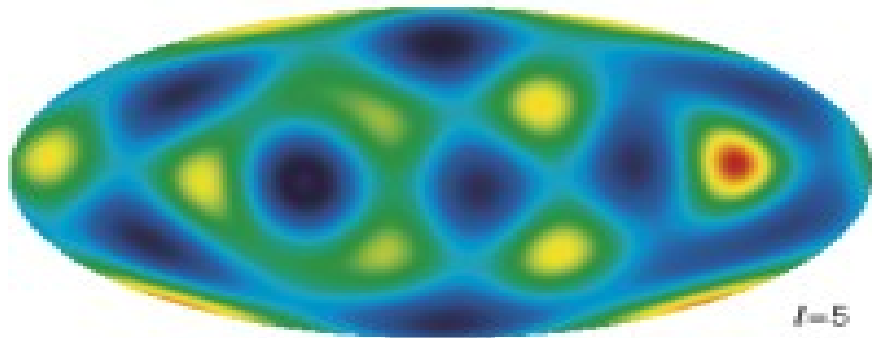
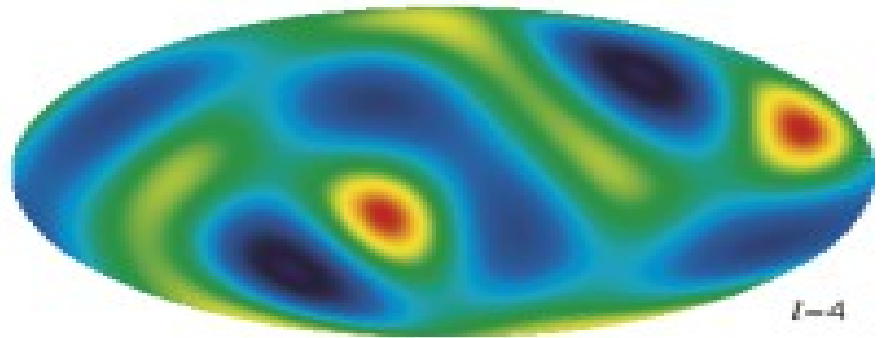
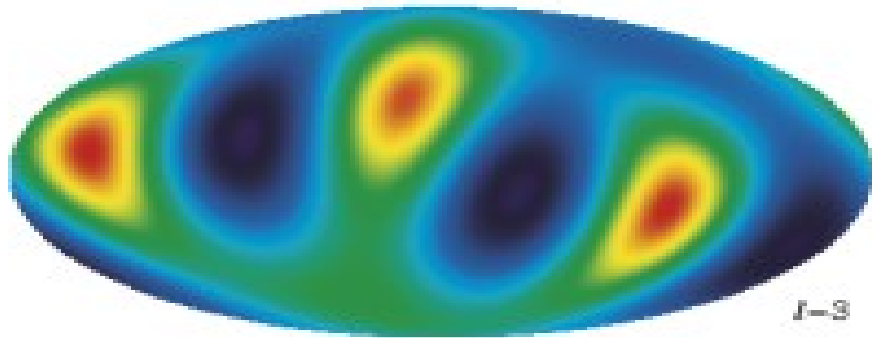
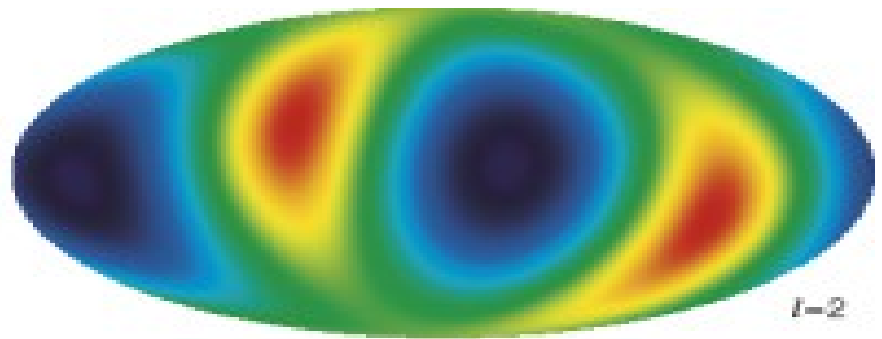
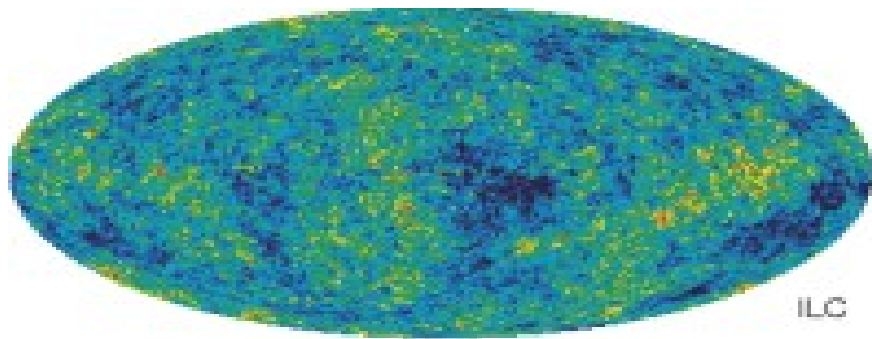


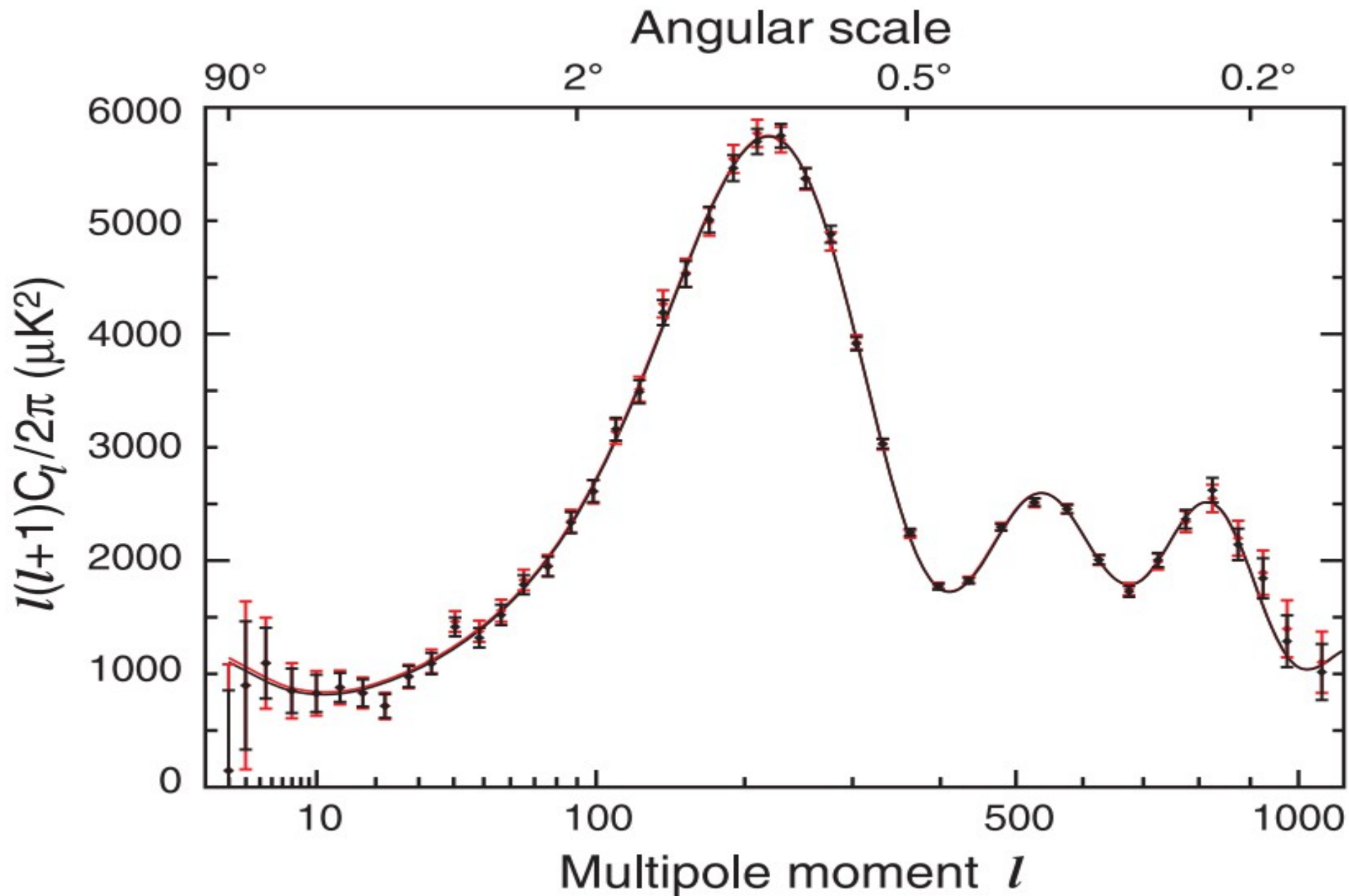
Data from the [Planck mission](#)
Toggle plane: galactic ecliptic
Resolution: low medium high
Channel: CMB or GHz 30 44 70 100 143 217 353

CMB Anisotropy, PLANCK



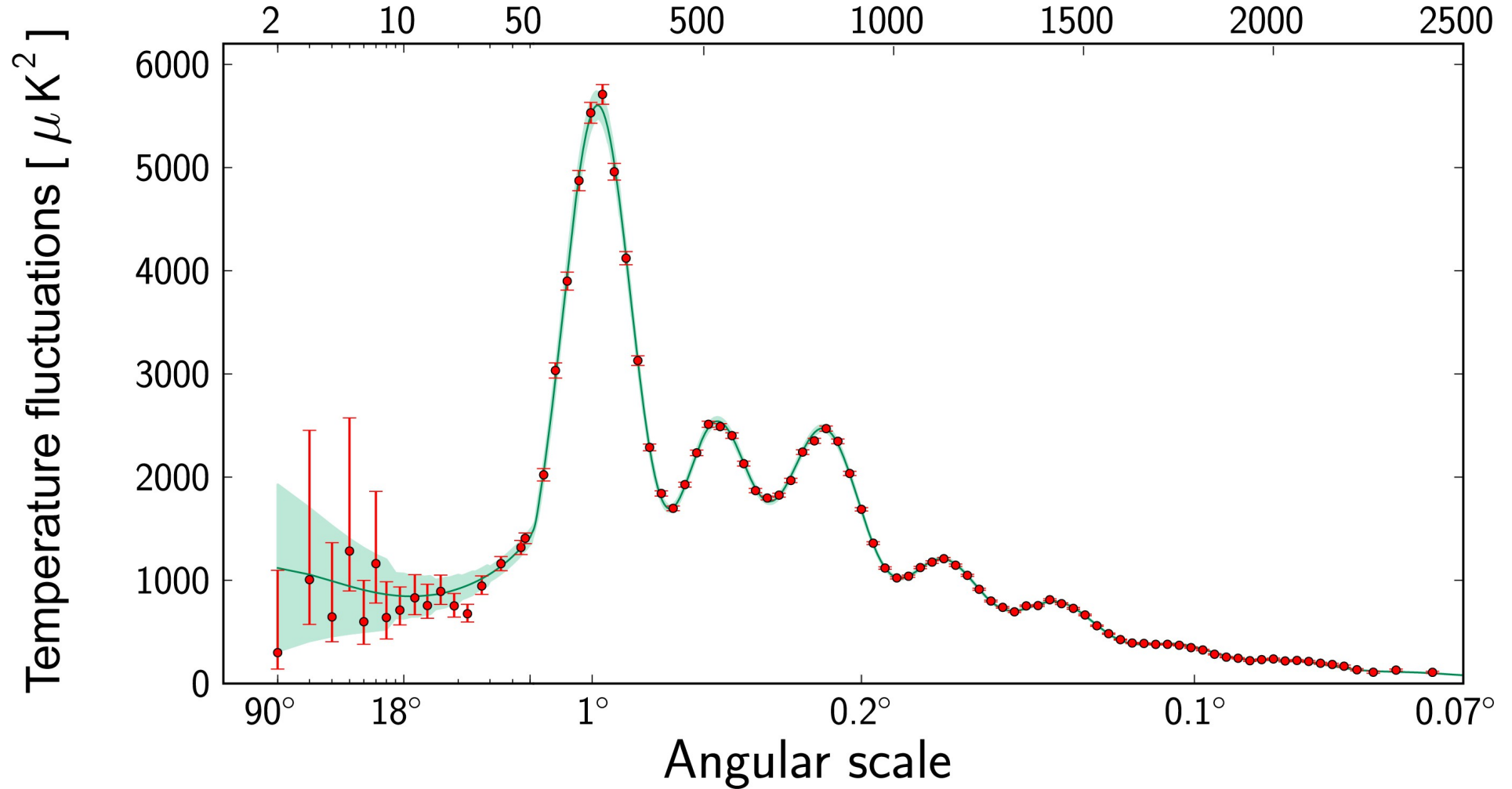
PLANCK, CMB anisotropy





WMAP 9 year data

Multipole moment, ℓ



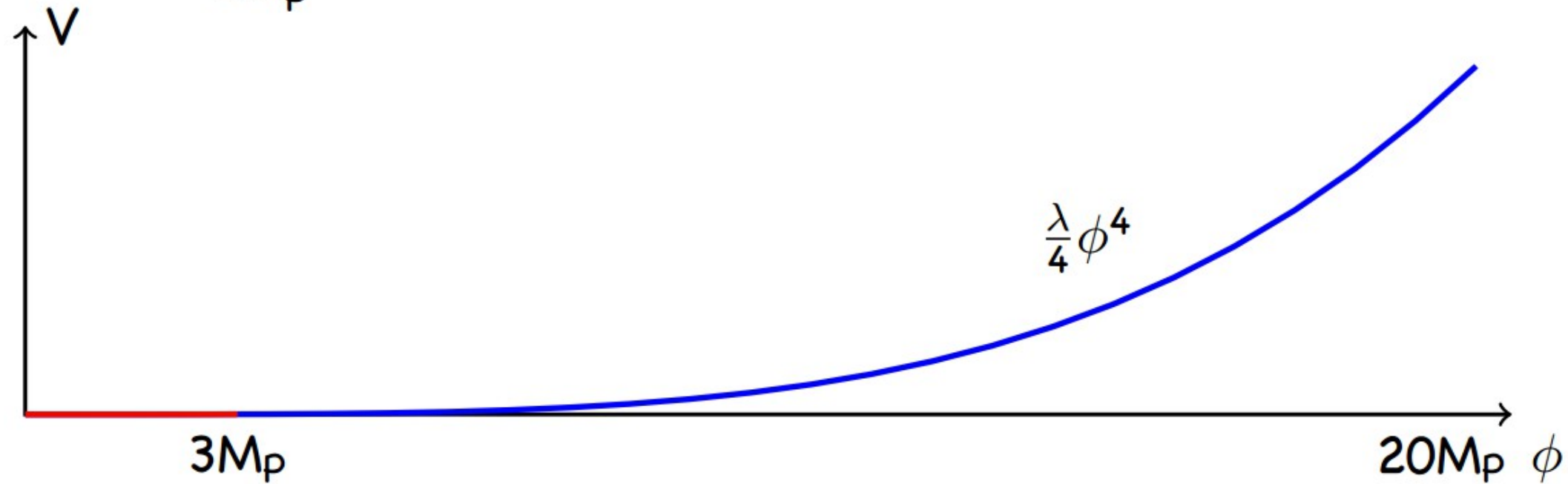
PLANCK 2013

Theory of Cosmic Inflation solves the problem of explaining super-horizon correlations.

It also provides a mechanism for generation of scale invariant density perturbations.

Slow roll inflation

$$\mathcal{H}^2 \simeq \frac{1}{3M_{\text{p}}^2} \left(V(\phi) + \dot{\phi}^2/2 \right), \quad \ddot{\phi} + 3\mathcal{H}\dot{\phi} + V'(\phi) = 0$$



Relating CMB observables to Inflation potential

$$A_s \approx \frac{V}{24\pi^2 M_{\text{pl}}^4 \epsilon_V}$$

$$A_t \approx \frac{2V}{3\pi^2 M_{\text{pl}}^4}$$

$$n_s - 1 \approx 2\eta_V - 6\epsilon_V$$

$$\epsilon_V \equiv \frac{M_{\text{pl}}^2}{2} \left(\frac{V_{,\phi}}{V} \right)^2, \quad \eta_V \equiv \frac{M_{\text{pl}}^2 V_{,\phi\phi}}{V}$$

PLANCK 2013

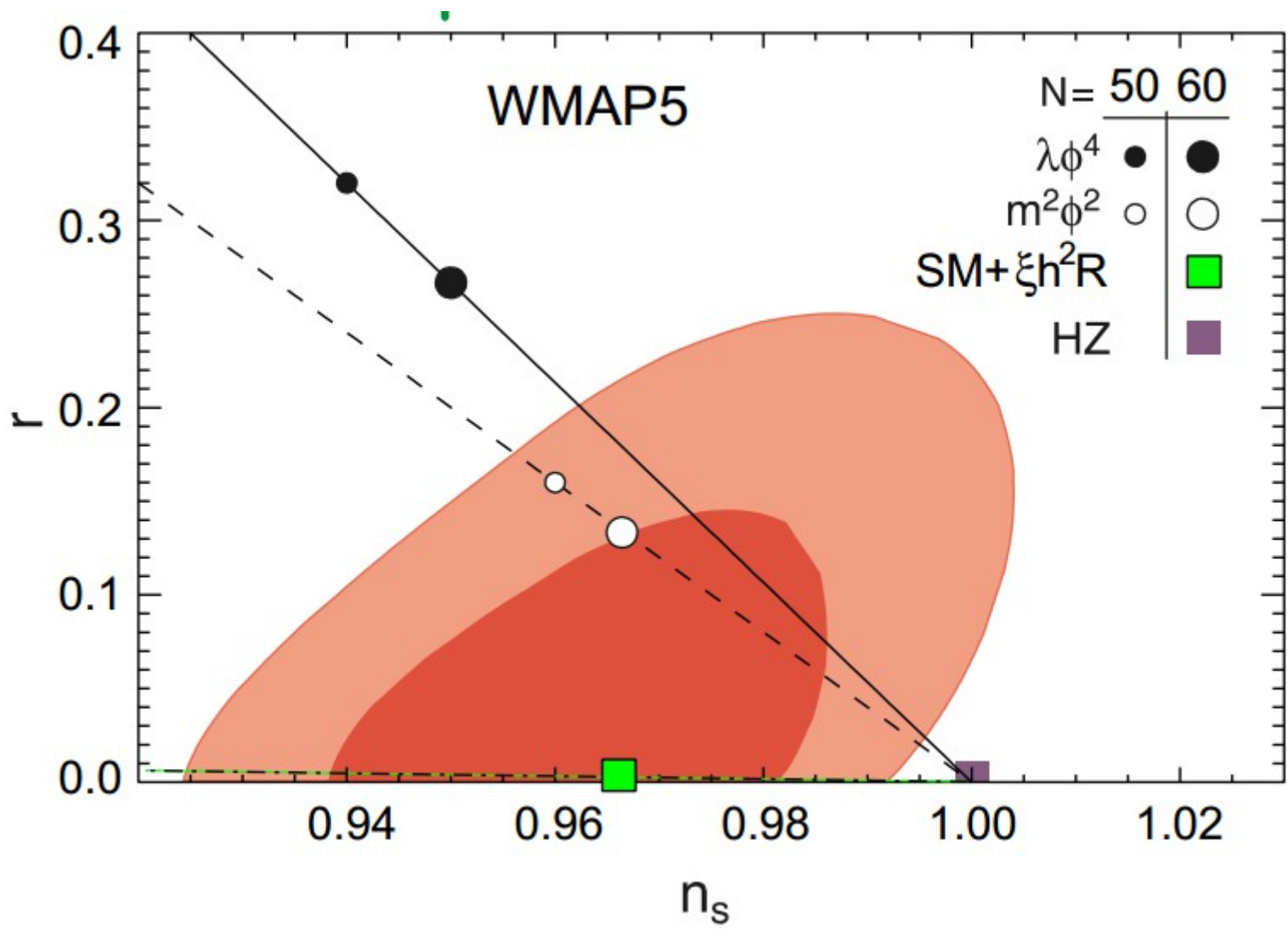
$$\Delta_{\mathcal{R}}^2 = 2.1955_{-0.585}^{+0.533} \times 10^{-9}$$

$$n_s = .9603 \pm .0073$$

$$\lambda \sim 10^{-12}$$

$$r = \frac{\mathcal{P}_t(k_*)}{\mathcal{P}_{\mathcal{R}}(k_*)} < 0.12 \text{ at } 95\% \text{ CL}$$

Energy scale of inflation: $V_* < (1.94 \times 10^{16} \text{ GeV})^4$



Inflaton with non-minimal curvature coupling

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2 + \xi\phi^2}{2} R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{\lambda}{4} (\phi^2 - \phi_0^2)^2 \right]$$

$$\hat{g}_{\mu\nu} = \Omega^2 g_{\mu\nu} , \quad \Omega^2 = 1 + \frac{\xi\phi^2}{M_P^2}$$

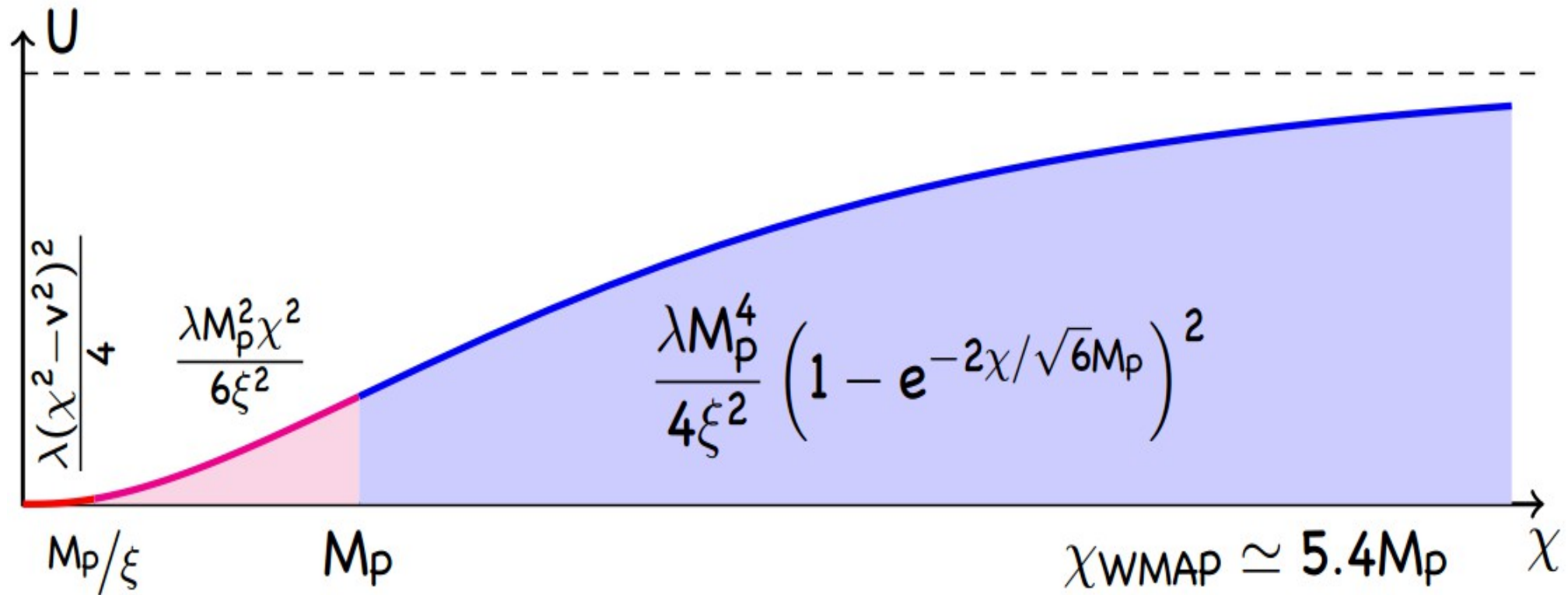
$$S = \int d^4x \sqrt{-\tilde{g}} \left(-\frac{M_P^2}{2} \tilde{R} + \frac{M_P^2 (M_P^2 + (6\xi + 1)\xi\phi^2)}{(M_P^2 + \xi\phi^2)^2} \frac{(\partial_\mu \phi)^2}{2} - \frac{\lambda M_P^4 \phi^4}{4(M_P^2 + \xi\phi^2)^2} \right)$$

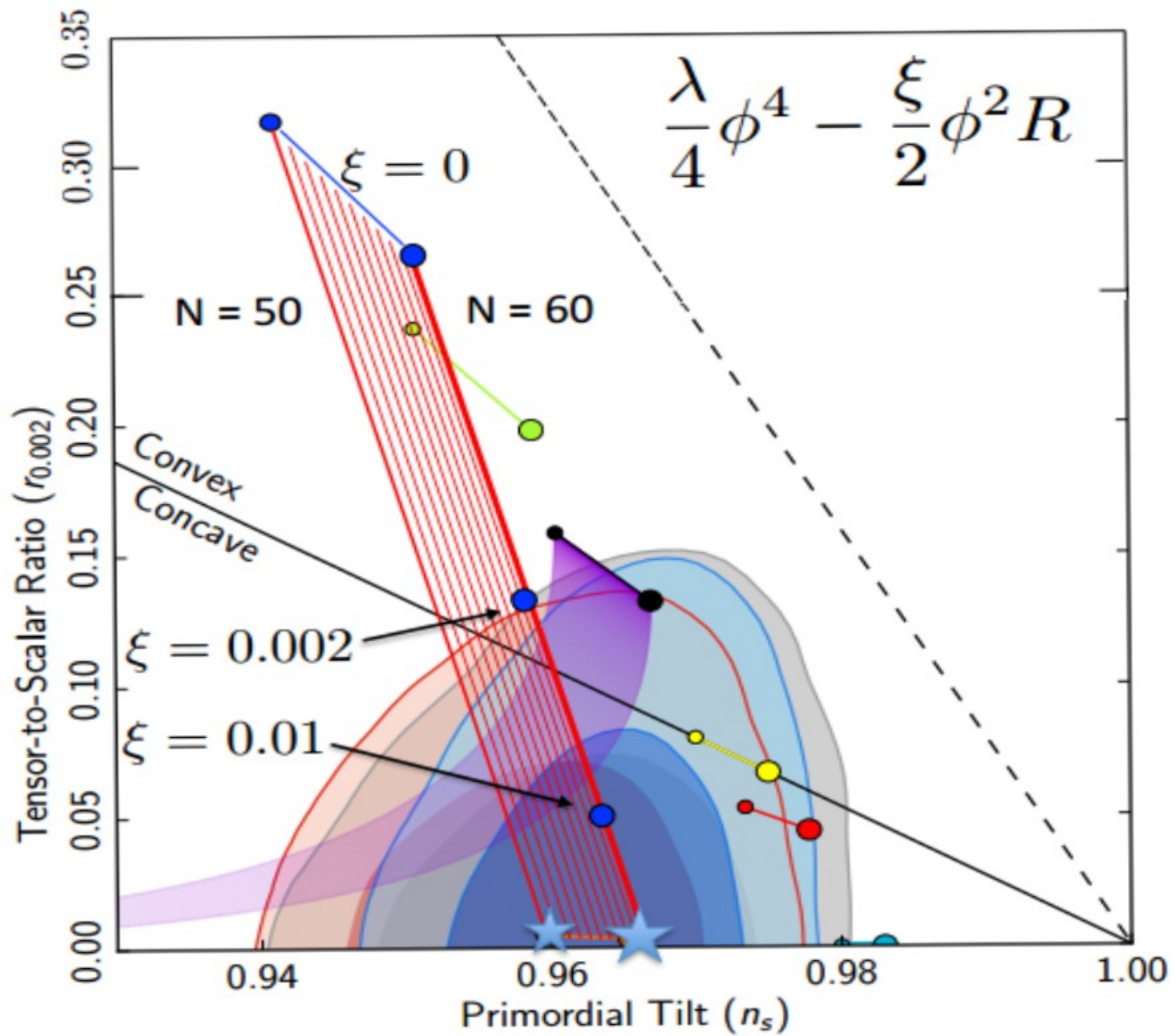
$$\frac{d\chi}{d\phi} = \frac{M_P \sqrt{M_P^2 + (6\xi + 1)\xi\phi^2}}{M_P^2 + \xi\phi^2}$$

$$S_E = \int d^4x \sqrt{-\hat{g}} \left\{ -\frac{M_P^2}{2} \hat{R} + \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) \right\}$$

$$U(\chi) = \frac{\lambda M_P^4}{4\xi^2} \left(1 - \exp \left(-\frac{2\chi}{\sqrt{6}M_P} \right) \right)^2$$

Effective potential in curvature coupling models





Standard model Higgs as Inflaton

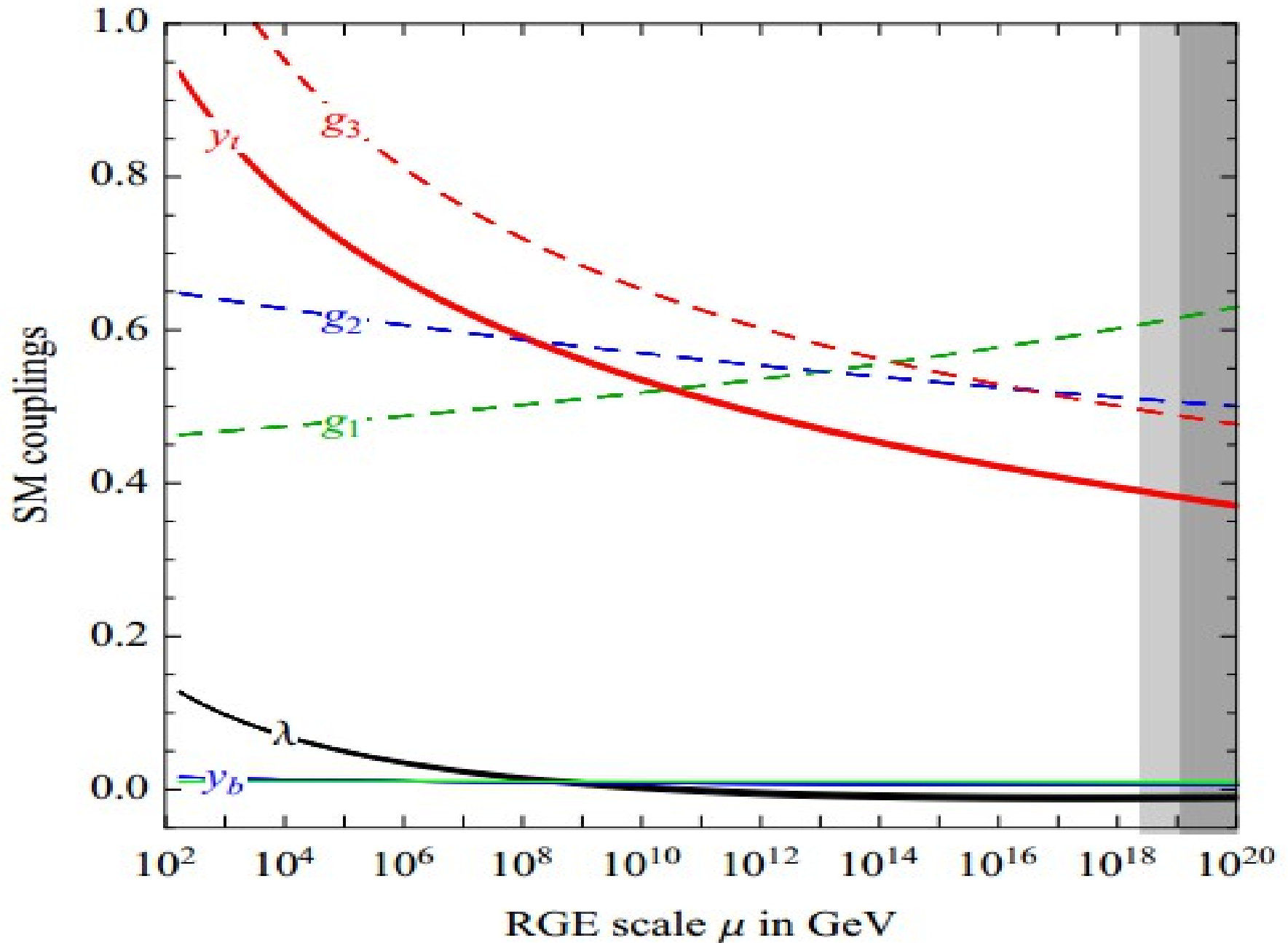
Bezrukov and Shaposhnikov 2008

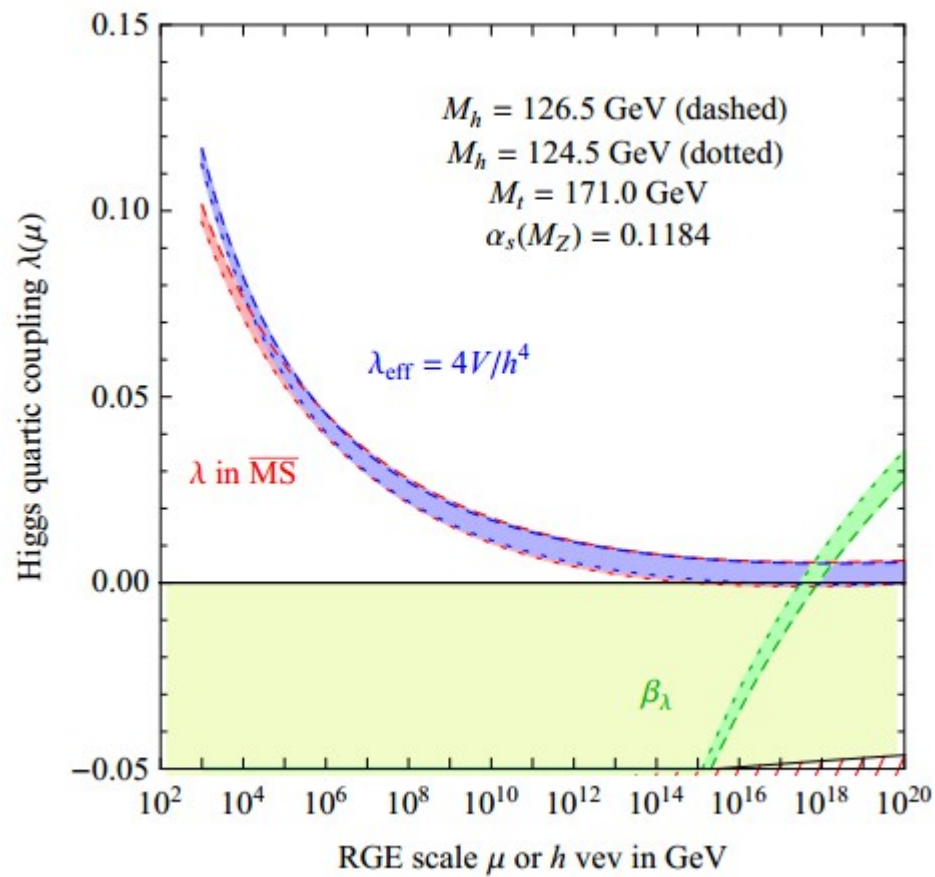
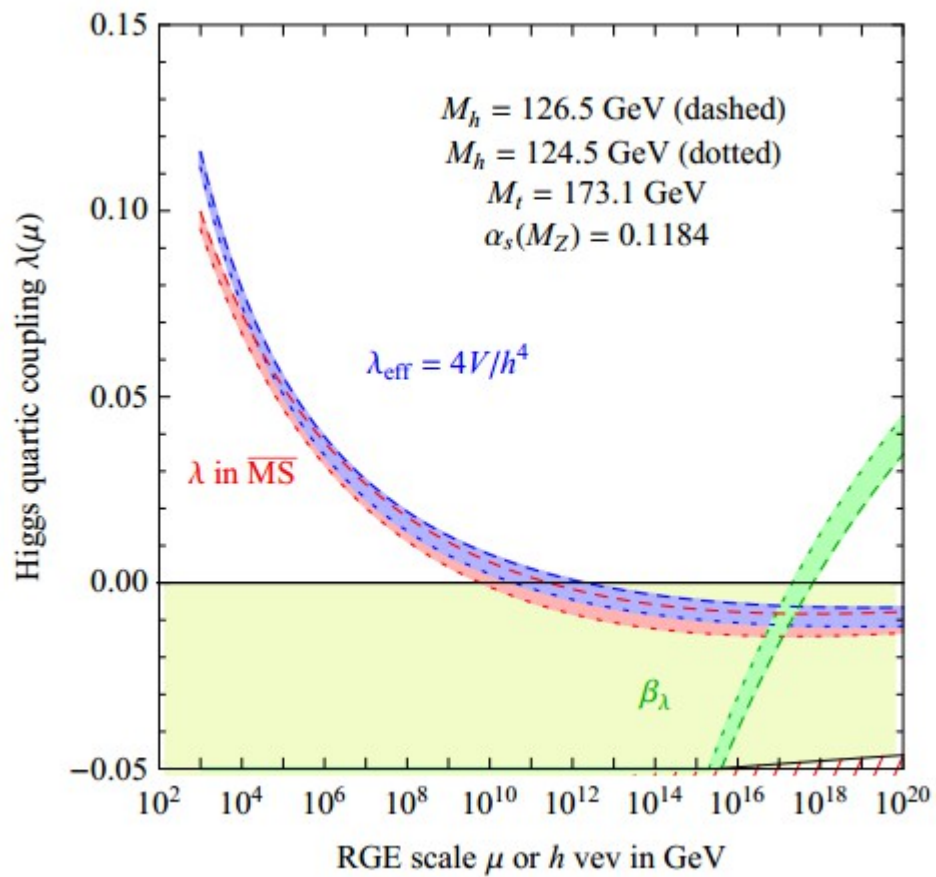
$$S_J = \int d^4x \sqrt{-g} \left\{ -\frac{M^2 + \xi h^2}{2} R + \frac{\partial_\mu h \partial^\mu h}{2} - \frac{\lambda}{4} (h^2 - v^2)^2 \right\}$$

$$\xi \simeq 49000 \sqrt{\lambda} = 49000 \frac{m_H}{\sqrt{2}v}$$

$$r = 16\epsilon \simeq 0.0033.$$

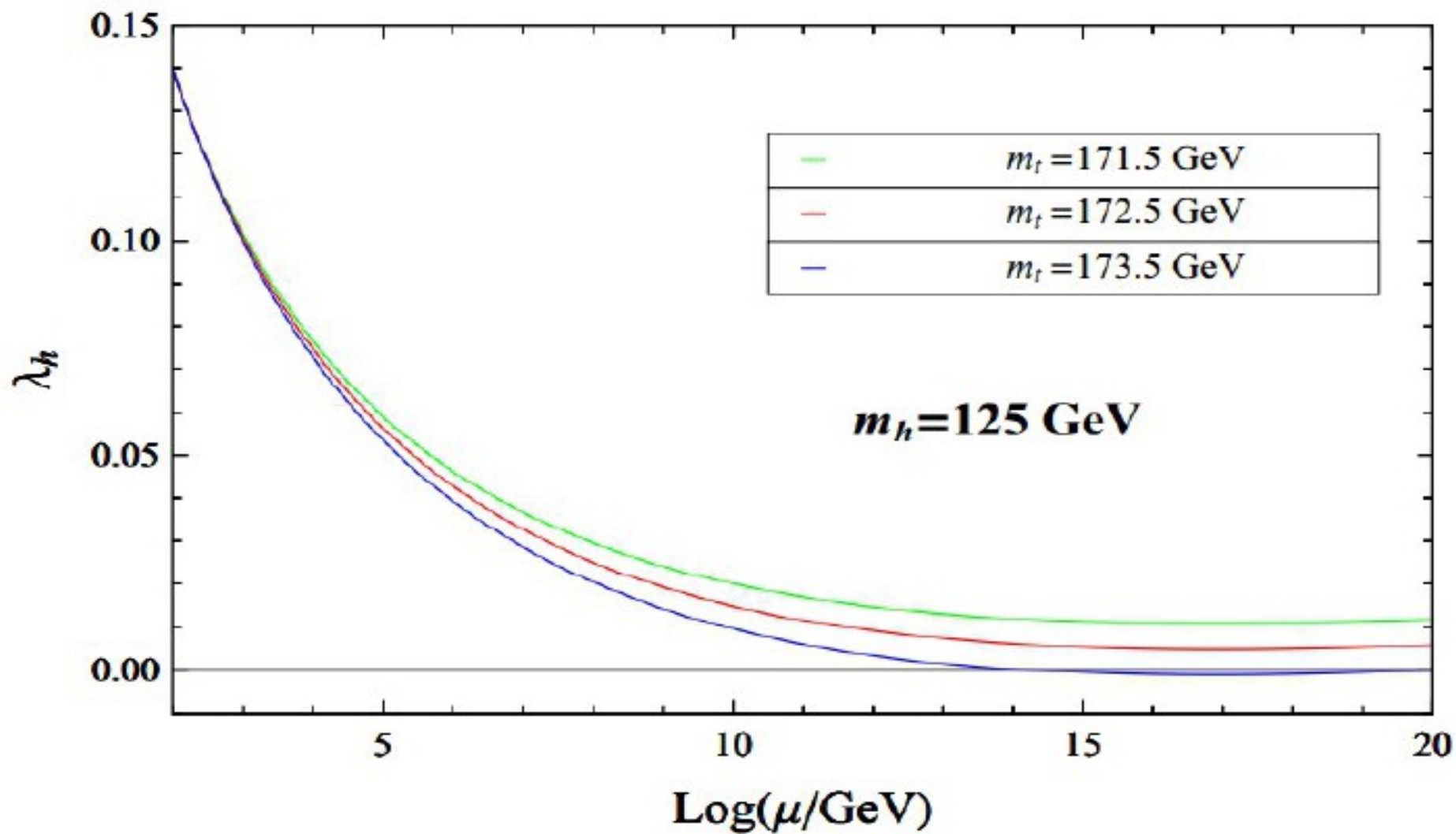
Running of the couplings in the Standard Model

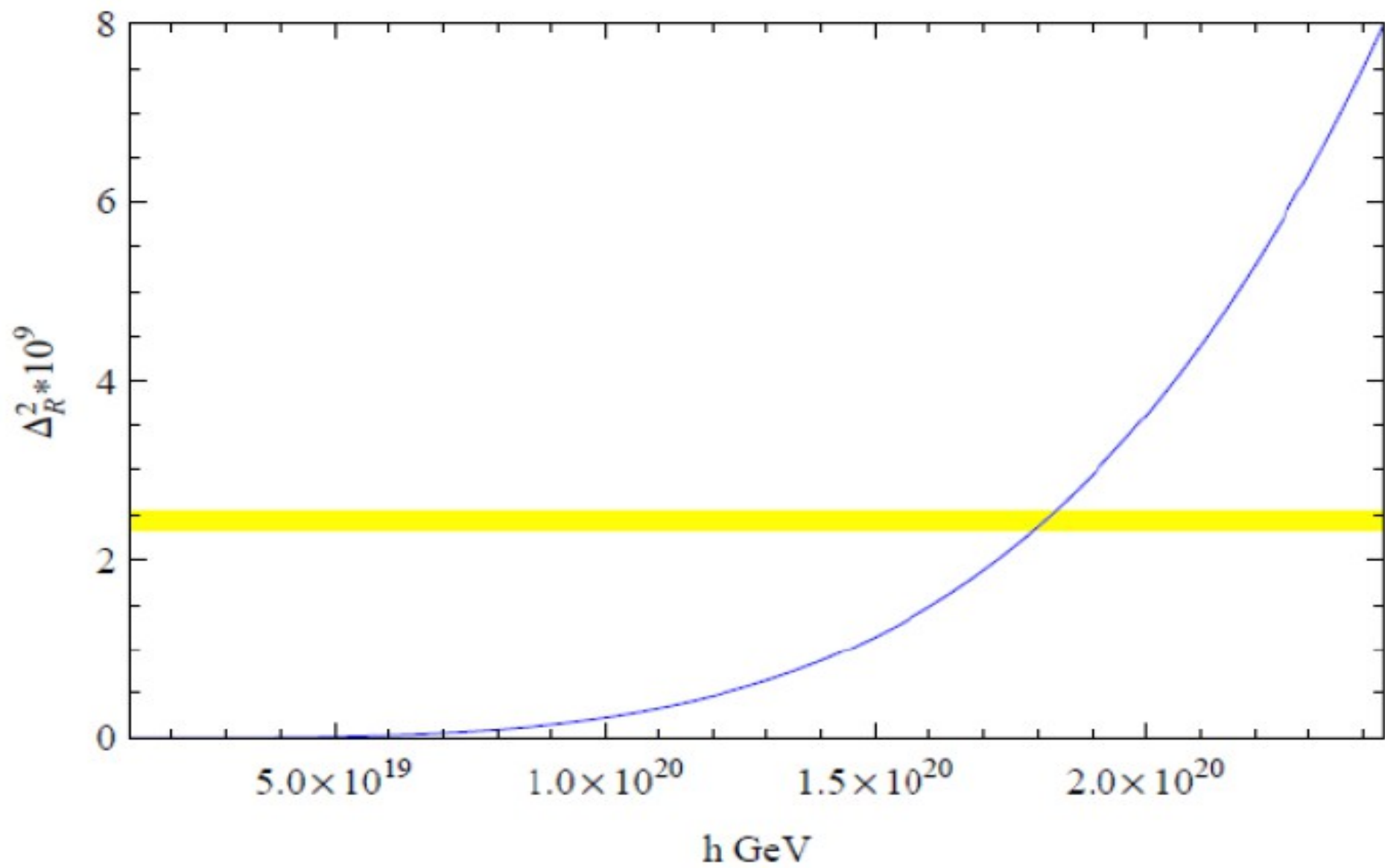




Degrassi et al 2013

Moumita Das, Joydeep Chakraborty, SM
arXiv 1207.2027





Moumita Das

f(R) theories are equivalent to scalar tensor theories

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} f(R) + S^{(m)}$$

$$\psi(\phi) \equiv f'(\phi)$$

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} [\psi(\phi)R - V(\phi)] + S^{(m)}$$

$$V(\phi) = \phi f'(\phi) - f(\phi)$$

Starobinsky model

$$S_S = \frac{1}{2} \int d^4x \sqrt{-g} \left(M_{\text{P}}^2 R + \frac{M_{\text{P}}^2}{6M^2} R^2 \right)$$

$$\tilde{g}_{\mu\nu} = (1 + \varphi/3M^2) g_{\mu\nu}$$

$$\varphi' = \sqrt{\frac{3}{2}} \ln \left(1 + \frac{\varphi}{3M^2} \right):$$

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[\tilde{R} + (\partial_\mu \varphi')^2 - \frac{3}{2} M^2 (1 - e^{-\sqrt{2/3} \varphi'})^2 \right]$$

Effective potential in the Starobinsky model

$$V = \frac{3}{4} M^2 \left(1 - e^{-\sqrt{2/3} \varphi'}\right)^2 M_P^2$$

is of the same form as the Higgs inflation model

$$U(\chi) = \frac{\lambda M_P^4}{4\xi^2} \left(1 - \exp\left(-\frac{2\chi}{\sqrt{6}M_P}\right)\right)^2$$

Distinguishing Higgs Inflation from Staroninsky Inflation from observation

$$T_H^{\text{reh}} \simeq 6 \times 10^{13} \text{ GeV} \quad , \quad T_{R^2}^{\text{reh}} = 3.1 \times 10^9 \text{ GeV}.$$

$$N_* \approx 57 - \frac{1}{3} \log \frac{10^{13} \text{ GeV}}{T_r}.$$

$$N_H = 57.66, \quad N_{R^2} = 54.37.$$

Higgs-inflation: $n_s = 0.967$, $r = 0.0032$,

R^2 -inflation: $n_s = 0.965$, $r = 0.0036$.

$$n_s = .9603 \pm .0073$$

F. L. Bezrukova, D. S. Gorbunov, 2012

Can we differentiate between HI and SI with non-Gaussinity ?

$$\langle \mathcal{R}(\mathbf{k}_1) \mathcal{R}(\mathbf{k}_2) \rangle = (2\pi)^3 \frac{2\pi^2}{k^3} \mathcal{P}_{\mathcal{R}}(k) \delta^3(\mathbf{k}_1 + \mathbf{k}_2)$$

$$\langle \mathcal{R}(\mathbf{k}_1) \mathcal{R}(\mathbf{k}_2) \mathcal{R}(\mathbf{k}_3) \rangle = (2\pi)^3 B_{\mathcal{R}}(k_1, k_2, k_3) \delta^3(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)$$

Higgs Inflation in $f(\Phi, R)$ Theory

$$S = \int d^4x \sqrt{-g} \left[-\frac{M_p^2}{2} R \left(1 + \frac{\xi \Phi^a R^{b-1}}{M_p^{a+2b-2}} \right) + \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi + \frac{\lambda \Phi^4}{4} \right]$$

Girish Chakravarty, SM, Naveen Singh ,

arXiv:1303.3870 [astro-ph.CO]

Amplitude of curvature perturbation

$$\Delta_{\mathcal{R}} = \frac{1}{\sqrt{Q_s}} \left(\frac{H}{2\pi} \right)$$

$$Q_s = \frac{\dot{\phi}^2 + 3\dot{F}^2 / (2\kappa^2 F)}{\left(H + \dot{F} / (2F) \right)^2}$$

$$F = \partial f / \partial R = 1 + \frac{\xi b \Phi^a R^{b-1}}{M_p^{a+2b-2}}$$

Tensor amplitude

$$P_T = \frac{2}{\pi^2} \left(\frac{H}{M_P} \right)^2 \frac{1}{F}$$

$$r \simeq \frac{8\kappa^2 Q_s}{F} \simeq 0.002$$

Parameters fit with the PLANCK+WMAP data

λ	0.1	10^{-2}	10^{-3}	10^{-4}	10^{-5}
$\phi_f _{(\phi_E=1M_p)}$	$0.0146M_p$	$0.0253M_p$	$0.044M_p$	$0.077M_p$	$0.134M_p$
$\phi_J _{(\phi_E=13M_p)}$	$3.566M_p$	$6.187M_p$	$10.77M_p$	$18.8M_p$	$32.77M_p$
a	3.56398962	3.27512990	3.02576940	2.80956100	2.62085100
b	0.21800513	0.36243484	0.48711456	0.59521700	0.68956620
a+2b	3.999999	3.999999	3.999998	3.999995	3.99998

TABLE II: The values of parameters (a,b) are evaluated in the Jordan frame at $\xi = 1$

$$S = \int d^4x \sqrt{-g} \left[-\frac{f(\Phi, R)}{2\kappa^2} + \frac{1}{2} g^{\mu\nu} \partial_\mu \Phi \partial_\nu \Phi + V(\Phi) \right]$$

$$\frac{1}{\kappa^2} f(\Phi, R) = \frac{1}{\kappa^2} R + \frac{\xi \Phi^a R^b}{M_p^{a+2b-4}} ; \quad V(\Phi) = \frac{\lambda \Phi^4}{4}$$

Scale invariant spectrum requires scale invariant interaction terms !

THANK YOU