Post-reionization HI 21 cm signal blogical constant A probe of negative cosm

Chandrachud B. V. Dash (BITS Pilani)

with Papomoy G Sarkar & Anjan A Sen

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OUTLINE

- \blacktriangleright Cosmological paradigm and H_0 tension Negative cosmological constant ► Observational outlook ► 21cm signal from post-reionization epoch ► Constraints
- > Summary



COSMOLOGICAL PARADIGM

Three unknown pillars:

- Inflation: an early stage of accelerated expansion which produces the initial, tiny, density perturbations, needed for structure formation.
- Dark Matter: a clustering matter component to facilitate structure formation.
- Dark Energy: an energy component to explain late time cosmic acceleration.

cosmologist's wishlist !!



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- Dark Matter: a clustering matter component to facilitate structure formation.
- Dark Energy: an energy component to explain late time cosmic acceleration.



Specific solutions for ΛCDM :

Inflation is given by a single, minimally coupled, slow-rolling scalar field;

Dark Matter is a pressureless fluid made of cold, i.e., with low momentum, and collisionless particles;

Dark Energy is a cosmological constant term.





ΛCDM

Space is homogenous and isotropic:

$$ds^{2} = c^{2}dt^{2} - a^{2}(t) \left[\frac{dr^{2}}{1 - Kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \ d\phi^{2}) \right]$$

Einstein's field equation:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda \mathbf{g}_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$$
$$\equiv H^2(a) = \frac{8\pi G}{c^4}\left(\sum_{i}\rho_i + \Lambda\right) \qquad (0)$$

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2(a) = \frac{8\pi}{c^2}$$

$$\frac{\ddot{a}}{a} = \frac{-4\pi G}{3} \left(\sum_{i} \rho_i + 3p_i \right) + \frac{\Lambda}{3}$$

(00 - Friedmann Eq)

(ij - Constraints of motion)



- Energy conservation: $\nabla_{\mu} T^{\mu}_{0} = 0$ gives $a^{-3}\frac{\partial(\rho a^3)}{\partial t} = -3\left(\frac{\dot{a}}{a}\right)p \quad ; \ \rho_i(a) \propto a^{-3(1+w_i)}$
 - Hubble expansion rate:

$$H^{2}(z) = H_{0}^{2} \left[\Omega_{m0} (1+z)^{3} + \Omega_{m0} (1+z)^{3} + \Omega_{m0} \right]$$
$$q_{0} = \frac{1}{2} \left[\Omega_{m0} + (1+z)^{3} + \Omega_{m0} \right]$$

- $\Omega_{r0}(1+z)^4 + \Omega_{de0}(1+z)^{3(1+w_{de})}$ $+ 3w)\Omega_{\Lambda 0}$, $w_i = p_i/\rho_i$
- In 90's $H_0 = 60 80$ km/s/MPc, $\Omega_{m0} = 1, \Omega_{\Lambda 0} = 0, q_0 > 0 (\Omega_{m0}/2)$



Distance modulus: $\mu = m$

Luminosity distance: $d_L(z) =$



$$e - M = H_0 \left(\frac{d_L(z)}{10Mpc} \right) + 25$$

$$\frac{c(1+z)}{H_0} \int_0^z \frac{dz}{E(z)}, \quad E(z) = \frac{H(z)}{H_0}$$
Robert P Kirshner 1999)
$$e = 0.3, \ \Omega_{\Lambda 0} = 0.7$$

 $q_0 < 0$ accelerating

Assuming w = -1













Everything seems so good so far !!!



WARNING !!!

With the improvement of the number and the accuracy of the observations, deviations from Λ CDM may be expected.

And, actually, discrepancies among key cosmological parameters of the models have emerged with different statistical significance.

While some proportion of these discrepancies may have a systematic origin, their persistence across probes should require multiple and unrelated errors, strongly hinting at cracks in the standard cosmological scenario and the necessity of new physics.

These tensions can indicate a failure of the concordance ACDM model.

HO TENSION ($\geq 5\sigma$)



(credit: LIGO)

> Planck measurement of H_0 is based on early universe physics.

 $H_0 = 67.5 \pm 0.5 \text{ km/s/Mpc}$

> SHoES measurement of H_0 is based on Astrophysics of stars.

 $H_0 = 73.7 \pm 1.04 \text{ km/s/Mpc}$







SHOES MEASUREMENT OF H_0



Relative Likelihood

HOW PLANCK MEASURES H_0 ??



Comoving sound horizon

$100\theta_* = 1.04109 \pm 0003$ Planck TT, TE, EE, lowE+lensing 0.03 % precision

 $H_0 = (67.3 \pm 0.6) km/s/Mpc$ (Planck, confirmed by ACT)

Comoving angular diameter distance $(D_A(z))$

HOW PLANCK MEASURES H_0 ?

Step 1: Calculate Sound horizon

$$r_s(z_r) = \int_0^{a_r} \frac{c_s da}{a^2 H(a)} ; c_s(a) = c/\sqrt{3(1+3\rho_b/4\rho_\gamma)}$$

> Step 2: Angular size of the sound horizon from the peak spacing in CMB

- Step 3: Calculate the angular diameter distance for sound horizon $D_A = -$
- ► Extrapolate H(z) to z = 0 and get H_0 (Here Planck uses Λ CDM)

$$\theta = \frac{\pi}{\Delta \ell}$$

$$\frac{1}{2} = \frac{1}{1+z_r} \int_0^{z_r} \frac{dz}{H(z)}$$

Crowded No More: The Accuracy of the Hubble Constant Tested with High Resolution Observations of Cepheids by JWST

ADAM G. RIESS,^{1,2} GAGANDEEP S. ANAND,¹ WENLONG YUAN,² STEFANO CASERTANO,¹ ANDREW DOLPHIN,³ LUCAS M. MACRI,⁴ LOUISE BREUVAL,² DAN SCOLNIC,⁵ MARSHALL PERRIN,¹ AND RICHARD I. ANDERSON⁶

Agrees with low redshift Agrees measurements Agrees measurements metric ¹Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA ²Department of Physics and Astronomy, Johns Hopkins University, Baltimore, MD 21218, USA ³Raytheon, 1151 E. Hermans Road, Tucson, AZ 85706 ⁴NSF's NOIRLab, 950 N Cherry Ave, Tucson, AZ 85719, USA ⁵Department of Physics, Duke University, Durham, NC 27708, USA ⁶Institute of Physics, Laboratory of Astrophysics, École Polytechnique Fédérale de Lausanne (EPFL), Observatoire de Sauverny, 1290 Versoix, Switzerland

High-resolution JWST observations can test confusion-limited HST observations for a photometric bias that could affect extragalactic Cepheids and the determination of the Hubble constant. We present JWST NIRCAM observations in two epochs and three filters of >320 Cepheids in NGC 4258 (which has a 1.5% maser-based geometric distance) and in NGC 5584 (host of SN Ia 2007af), near the median distance of the SH0ES HST SN Ia host sample and with the best leverage among them to detect such a bias. JWST provides far superior source separation from line-of-sight companions than HST in the NIR to largely negate confusion or crowding noise at these wavelengths, where extinction is minimal. The result is a remarkable $>2.5\times$ reduction in the dispersion of the Cepheid P-L relations, from 0.45 to 0.17 mag, improving individual Cepheid precision from 20% to 7%. Two-epoch photometry confirmed identifications, tested JWST photometric stability, and constrained Cepheid phases. The P-L relation intercepts are in very good agreement, with differences (JWST-HST) of 0.00 ± 0.03 and 0.02 ± 0.03 mag for NGC 4258 and NGC 5584, respectively. The difference in the determination of H₀ between HST and JWST from these intercepts is 0.02 ± 0.04 mag, insensitive to JWST zeropoints or count-rate non-linearity thanks to error cancellation between rungs. We explore a broad range of

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ABSTRACT





**credit: Di Valentino et al., (2021)

Cyan vertical band corresponds to SHoES Team and the pink vertical band corresponds to Planck18 team within a ΛCDM scenario.

> The high precision and consistency of the data at both ends present strong challenges to the possible solution space and demands a hypothesis with enough rigor to explain multiple observations whether these invoke new physics, unexpected large-scale structures or multiple, unrelated errors.



POSSIBLE SOLUTION



POSSIBLE SOLUTION (FORMALLY)

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]
Vacuum Metamorphosis [402]	Bulk viscous models [340, 341]	Decaying Ultralight Scalar [538]
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	BD- Λ CDM [852]
Sterile Neutrinos [433]	Phantom Braneworld DE [378]	Metastable Dark Energy [314]
Decaying Dark Matter [481]	PEDE [391, 392]	Self-Interacting Neutrinos [700]
Neutrino-Majoron Interactions [509]	Elaborated Vacuum Metamorphosis [401]	Dark Neutrino Interactions [716]
IDE [637, 639, 657, 661]	IDE [659,670]	IDE [634-636, 653, 656, 663, 669]
DM - Photon Coupling [685]	Interacting Dark Radiation [517]	Scalar-tensor gravity [855, 856]
$f(\mathcal{T})$ gravity theory [812]	Decaying Dark Matter [471, 474]	Galileon gravity [877,881]
BD-ΛCDM [851]	DM - Photon Coupling [686]	Nonlocal gravity [886]
Über-Gravity [59]	Self-interacting sterile neutrinos [711]	Modified recombination [986]
Galileon Gravity [875]	$f(\mathcal{T})$ gravity theory [817]	Effective Electron Rest Mass [989]
Unimodular Gravity [890]	Über-Gravity [871]	Super ACDM [1007]
Time Varying Electron Mass [990]	VCDM [893]	Axi-Higgs [991]
ACDM [995]	Primordial magnetic fields [992]	Self-Interacting Dark Matter [479]
Ginzburg-Landau theory [996]	Early modified gravity [859]	Primordial Black Holes [545]
Lorentzian Quintessential Inflation [979]	Bianchi type I spacetime [999]	
Holographic Dark Energy [351]	f(T) [818]	

**credit: Di Valentino et al., (2021)



DARK ENERGY MODEL WITH NON ZERO VACUA

- ► EoS: $w_{\phi}(a) = w_0 + w_a(1 a)$ CPL model **Chavellier, Polarski & Linder
- > This can be obtained by expanding w(a) in the Taylor series around a = 1

$$w(a) = w|_{a=1} + (a-1)\frac{dw}{da}|_{a=1}$$

where w_0 is the present value i.e $w(a_0) = w_0 + w_0(1 - a_0) = w_0$ and the value of its slope $dw(a)/da = -w_a$.

> Density evolve: $\rho_{\phi}(a) \propto a^{-3(1+w_0+w_a)} \exp^{3w_a a}$ *According to ancient Greek science it denotes a fifth cosmic element after earth, fire, water & air. In Latin "fifth element" 19

> Quintessence field*(ϕ) + Cosmological constant (Λ) Total Dark Energy

► i.e $\rho_{DE} = \rho_{\phi} + \Lambda$ with the constraint $\rho_{DE} > 0$ (for late time cosmic acceleration)

 $+\frac{1}{2}(a-1)^2\frac{d^2w}{da^2}\Big|_{a=1} + \mathcal{O}\left[(a-1)^3\right]$

QUINTESSENCE

gravity. A slowly varying field along a potential $V(\phi)$ can lead to the acceleration of the Universe.

$$S = \int d^4 \vec{x} \sqrt{-g} \left[\frac{1}{2} M_{p_l}^2 R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right] + S_m$$

> The dark energy EoS is $w_{\phi} \equiv \frac{P_{\phi}}{\rho_{\phi}} = \frac{\phi^2/2 - V(\phi)}{\dot{\phi}^2/2 + V(\phi)}$

Slow roll condition: $\dot{\phi}^2 < V(\phi) \& |\ddot{\phi}| < |V_{\phi}|$ where $V_{\phi} = dV(\phi)/d\phi$

define the following dimensionless variables
$$x = \left(\frac{d\phi}{dN}\right)/\sqrt{6}M_{pl}, y = \frac{\sqrt{V}}{\sqrt{6}HM}$$
, $\lambda = -M_{pl}\frac{V_{\phi}}{V}, \Gamma = V\frac{V_{\phi\phi}}{V_{\phi}}$, we have: $\Omega_{\phi} = x^2 + y^2$, $w_{\phi} = \frac{2x^2}{x^2 + y^2} - 1$.

 \blacktriangleright Quintessence is described by a canonical scalar field ϕ minimally coupled to

$$-V(\phi)$$







> In spatially flat universe, evolution of H(a)

$$\frac{H(a)}{H_0} = \sqrt{\Omega_{m_0} a^{-3} + \Omega_{\phi 0} \exp\left[-3 \int_1^a da' \frac{1 + w_{\phi}(a')}{a'}\right] + \Omega_{\Lambda 0}}$$

This is same as adding a non-zero cosmological constant for the scalar field potential $V(\phi) = F(\phi) + V_0$

with $\Omega_{m0} + \Omega_{\phi0} + \Omega_{\Lambda} = 1$ (CPL- Λ CDM)

 V_0 can be positive (dS) or negative (AdS)



OBSERVATIONAL OUTLOOK

$$\mathbf{r} \theta_s(z) = \frac{s(z_d)}{(1+z)D_A(z)} \qquad \delta z_s = \frac{s(z_d)H}{c}$$

► Related effective distance: $D_V(z) = (1)$

► Define a dimensionless quantity: $r_{BAO}(z) = \frac{r_s}{D_s(z)}$

- > 2df galaxy survey: $r_{BAO}(z = 0.2) = 0.1980 \pm 0.0058$ and $r_{RAO}(z = 0.35) = 0.1094 \pm 0.0033$ (Percival et. al. 2007)

H(z)

$$(+z)^2 D_A^2(z) \frac{cz}{H(z)} \Big]^{1/3}$$

(**Eisenstein et. al. 2005)

► BOSS (SDSS III) CMASS_LRG: $r_{BAO}(z = 0.57) = 0.07315 \pm 0.002$ (Anderson et. al. 2012)



$r_{BAO}(z = 0.57) = 0.07315 \pm 0.002$ (Anderson et. al. 2012)



- 0.1068	Task!!!
- 0.0996	Find out what is the possible
- 0.0924	combinations for (w_0, Ω_Λ) that
- 0.0852	give identical r_{BAO} in ΛCDM
- 0.0780	framework but with $H_0 = 72$
- 0.0708	km/s/Mpc
- 0.0636	For simplicity we kept $w_a = 0$
- 0.0564	and fix $\Omega_{m0} = 0.311$ measured
- 0.0492	by Planck.
0.0420	—observation data,·····ΛCDM
	- 0.0996 - 0.0924 - 0.0852 - 0.0780 - 0.0708 - 0.0708 - 0.0564 - 0.0564 - 0.0492















z = 0.35

z = 0.2

Same but for different redshifts



21CM SIGNAL



**credit: SKAO

- ► In H-atom electron and proton normally spin in the same direction.
- Occasionally, electron flips spin to other direction (happens only about once every 100million years for each atom)
- > Emit radiation of 1420MHz ($\lambda = 21$ cm)
- $\sum_{n=1}^{n_1} = \frac{g_1}{m_1} \exp^{T_*/T_s} ; T_* = h\nu_e/k_B = 0.68K$ n_0 g_0
- ► predicted: Henrik van de Hulst 1944 detected in 1951.









21CM SIGNAL



CMB acts as backlight

Neutral gas imprints signal $z = 1, \nu_e = 1420 \text{ MHz}$

Redshifted signal detected $z = 0, \nu_e = 710 \text{ MHz}$





21CM SIGNAL FROM POST-REIONIZATION EPOCH



**Fig. credit: S Bharadwaj & SK. Saiyad Ali, (2004)

CMB photons free streaming from the LSS undergoes a dip in its brightness temperature on passing through a HI cloud depending on its optical depth and spin temperature.

 $^{\prime}HI$

CMBR





21CM POWER SPECTRUM

where
$$\mu = \frac{k_{\parallel}}{k} = \cos\theta$$
, $f = d\log D_+/d\log a$ is the growth rate a
 $\mathscr{A}_T(z) = 4.0mK \ b_T \bar{x}_{HI}(1+z)^2 \left(\frac{\Omega_{b0}h^2}{0.02}\right) \left(\frac{0.7}{h}\right) \frac{H_0}{H(z)}$

$$P_{21}(k, z, \mu) = \frac{\mathscr{A}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1 - F^2)} \right]^2 P_m \left(\frac{k}{\alpha_{\perp}} \sqrt{1 + \mu^2(F^{-2} - 1)}, z \right)$$

(Alcock-Paczynski test: $\alpha_{\parallel} = H^f/H^r, \alpha_{\perp} = D_A^r/D_A^f, F = \alpha_{\parallel}/\alpha_{\perp}$)

The 21cm power spectrum of this excess brightness temperature field is given by $P_{21}(k, z, \mu) = \mathscr{A}_T^2 (1 + f(z)\mu^2)^2 P_m(k, z)$

and A_T is given by

The bias (with respect to the dark matter field) and the mean neutral fraction completely models the 21-cm signal in the post-reionization epoch.*



ONGOING/UPCOMING EXPT.(S)



MWA (Australia)



CHIME (Canada)



SKA (South Africa)



uGMRT (India)



MeerKAT (Australia)





			3D HI 21cm power
- 1.340	- 1.0	48	spectrum at $z = 1$ in
- 1.195	- 0.9	36	$(k_{\perp}, k_{\parallel})$ space.
- 1.050	- 0.8	24	The colorbar shows the
- 0.905	- 0.7	12	value of the dimensionless
- 0.760 - U.760 - Ju	- 0.6	VCDM VCDM	quantity
- 0.615	- 0.4	88	$\Delta_{21}^2 = k^3 P_{21}(\mathbf{k}) / (2\pi^2)$
- 0.470	- 0.3	76	The asymmetry is
- 0.325	- 0.2	64	indicative of redshift space
- 0.180	- 0.1	52	distortions.
L 0.035	0.0	40	Alcock-Paczynski effect enhanced the distortions



MULTIPOLES OF 21CM PS

polynomials $\mathscr{P}_{\ell}(\mu)$ as

 $P_{21}(k,\mu,z) =$

The first few Legendre polynomials are given by $\mathscr{P}_0(\mu) = 1, \quad \mathscr{P}_2(\mu) = \frac{1}{2} \left(3\mu^2 - 1 \right), \quad \mathscr{P}_4(\mu) = \frac{1}{8}$

Coefficients of the expansion of the 21cm power spectrum, can be found by inverting

$$P_{\ell}(z,k) = \frac{(2\ell + 2)}{2}$$

➤ Redshift space 21cm power spectrum can be decomposed in the basis of Legendre

$$\sum_{\ell} P_{\ell}(z,k) \mathcal{P}_{\ell}(\mu)$$

$$\mathcal{P}_4(\mu) = \frac{1}{8}(35\mu^2 - 30\mu^2 + 3)$$

 $\frac{(+1)}{2}\int d\mu \mathcal{P}_{\ell}(\mu)P_{21}(z,k,\mu)$



SKA1-MID SPECIFICATIONS



More info: www.skatelescope.org

- ► Medium deep band-2 survey sky area coverage 5,000 deg² in frequency range $\nu = [0.95 - 1.75]GHz, z = [0 - 0.5]$
- ► Wide band-1 survey sky area coverage $20,000 \text{ deg}^2$ in the frequency range $\nu = [0.35 1.05]GHz, z = [0.35 3]$
- ➤ No. of antennae = 250, $D_{dis} = 15m$, System temperature $T_{sys} = 60K$, Antennae efficiency $\epsilon = 0.7$, Total observation time $T_0 = 500 \times 150$ with 150 independent pointings, Bandwidth B = 128MHz





z = 0.2

Monopole

z = 0.57



Quadrupole







Hexadecapole



z = 0.2

z = 0.57



BAO IMPRINTS

- > The sound horizon at the epoch of recombination is given by $s(z_d) = \int_{0}^{\alpha_r} \frac{c_s da}{a^2 H(a)}$
- the radial direction as
 - $\theta_s(z) = \frac{S(z_d)}{(1+z)I}$
- bins, in the redshift range z = 0 3 with $\Delta z = 0.5$.

;
$$c_s(a) = c/\sqrt{3(1+3\rho_b/4\rho_\gamma)}$$

► rThe standard ruler 's' defines a transverse angular scale and a redshift interval in

$$\frac{\delta d}{\delta D_A(z)} \qquad \delta z_s = \frac{s(z_d)H(z)}{c}$$

> Measurement of θ_{s} and δz_{s} , allows the independent determination of $D_{A}(z)$ and H(z).

> We calculate the expected error projections in five evenly space non-overlapping



$$\begin{split} F_{ij} &= \left(\frac{2\ell+1}{2}\right) \int dk' \int_{-1}^{+1} d\mu \; \frac{\mathscr{A}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \frac{\partial P_b(k')}{\partial p_i} \frac{\partial P_b(k')}{\partial p_i} \frac{\partial P_b(k')}{\partial p_j} \right] \\ &= \left(\frac{2\ell+1}{2}\right) \int dk' \int_{-1}^{+1} d\mu \; \frac{\mathscr{A}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \left(\cos x - \frac{\sin x}{x} \right)^2 f_i f_j A^2 e^{-2(k' \sum_s)^{1/4}} e^{-k} d\mu \int_{-1}^{1} d\mu \; \frac{\mathscr{A}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \left(\cos x - \frac{\sin x}{x} \right)^2 f_i f_j A^2 e^{-2(k' \sum_s)^{1/4}} e^{-k} d\mu \int_{-1}^{1} d\mu \; \frac{\mathscr{A}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \left(\cos x - \frac{\sin x}{x} \right)^2 f_i f_j A^2 e^{-2(k' \sum_s)^{1/4}} e^{-k} d\mu \int_{-1}^{1} d\mu \; \frac{\mathscr{P}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \left(\cos x - \frac{\sin x}{x} \right)^2 f_i f_j A^2 e^{-2(k' \sum_s)^{1/4}} e^{-k} d\mu \int_{-1}^{1} d\mu \; \frac{\mathscr{P}_T^2}{\alpha_{\parallel} \alpha_{\perp}^2} \left[b_T + \frac{f(z)\mu^2}{F^2 + \mu^2(1-F^2)} \right]^2 \frac{\mathscr{P}_{\ell}(\mu)}{\delta P_{21}^2(k,z,\mu)} \left(\cos x - \frac{\sin x}{x} \right)^2 f_i f_j A^2 e^{-2(k' \sum_s)^{1/4}} e^{-k} d\mu \int_{-1}^{1} d\mu \int_{-1}^{1}$$

► BAO Power spectrum: $P_b(k') = A \frac{\sin x}{x}$

> Parameter(s): $p_i = \ln(s_{\perp})^{-1}$ and $p_j = \ln s_{\parallel}$ Estimated error: $\sqrt{F_{ii}^{-1}}$

$$-(k'\sum_{s})^{1.4}e^{-k'^2\sum_{nl}^2/2}$$











Fiducial values taken from Anjan A Sen et. al (2022)

$$\Omega_{\Lambda} = -0.883^{0.978}_{-2.987}$$

$$w_0 = -1.030^{0.023}_{-0.082}$$

$$w_a = -0.088^{0.162}_{-0.343}$$
(Constraints with 1 - σ uncertainty)





SUMMARY

- ► Beyond ACDM: new physics at early universe or late universe.
- > Allowing a dS/AdS minima gives a hope to alleviate tension.
- Observation of redshifted 21cm signal from post-deionization epoch can distinguishable from standard cosmological model with $> 3\sigma$ sensitivity.
- However 21cm signal alone can not put tight constraints, need to look more comprehensive approach.

> Λ CDM paradigm: tensions from direct and indirect measurement of $\geq 5\sigma$



TAKE HOME MESSAGE

- ► Dark energy is still an open problem. We have no idea (honestly !!)
- We're always hoping to find new things, but we're finding that our model is really, really good - - maybe disappointingly good. (Max Tegmark)
- Detection of 21cm signal will be a giant leap for precision cosmology.
- Upcoming SKA experiment promises to deliver qualitatively new insights on astrophysics and cosmology.





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