## A fast method for quantifying the shape dependence of the bispectrum

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## Cosmological random fields

Large-scale cosmological maps contain random fluctuations (more or less NG)


Summary statistics (mostly power spectrum) are used to quantify the information encoded in these random fluctuations.

## Insufficient Power spectra

Power spectrum is sufficient to quantify fluctuations if the signal is a Gaussian random field.
In order to capture the non-Gaussian features of the field one need higher-order statistics (e.g. bispectrum, trispectrum etc.)




## Bispectrum (Three-point statistic)

Bispectrum is the higher order statistics next to PS
Sensitive to the non-Gaussianity in the map
Vanishes is the field is Gaussian random

Considering statistical homogeneity


$$
\delta_{\mathrm{K}}\left(\vec{k}_{1}+\vec{k}_{2}+\vec{k}_{3}\right) B\left(k_{1}, k_{2}, k_{3}\right)=V^{-1}\left\langle\Delta\left(\vec{k}_{1}\right) \Delta\left(\vec{k}_{2}\right) \Delta\left(\vec{k}_{3}\right)\right\rangle
$$



## Binned Bispectrum Estimator

$$
\hat{B}\left(k_{1}, k_{2}, k_{3}\right)=\frac{1}{V N_{\operatorname{tri}}} \sum_{\vec{k}_{a_{1}}} \sum_{\vec{k}_{a_{2}}} \sum_{\vec{k}_{a_{3}}} \Delta\left(\vec{k}_{a_{1}}\right) \Delta\left(\vec{k}_{a_{2}}\right) \Delta\left(\vec{k}_{a_{3}}\right) \delta_{\mathrm{K}}\left(\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right)
$$


$\mathrm{k}_{1}$
Binning reduces the storage cost $+$
decreases the statistical uncertainties in the estimates


## Issue

- For a cube with $N^{3}$ grids, time required for the direct search approach for closed triangles scales as $N^{6}$.
- LSS simulations/ future observations will have large data volume. A direct search method will be extremely time consuming and computationally intractable.
- Need an efficiently faster technique to estimate bispectrum accurately for all possible closed triangles of unique shapes.


## Fourier Technique

## Numerator

## Denominator

$$
N_{\text {tri }}=\sum_{\vec{k}_{a_{1}}} \sum_{\vec{k}_{a_{2}}} \sum_{\vec{k}_{a_{3}}} \delta_{\mathrm{K}}\left(\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right)
$$

Evaluate the sum using

$$
\delta_{\mathrm{K}}\left(\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right)=\frac{1}{N_{\mathrm{g}}^{3}} \sum_{\vec{x}} \exp \left(-i\left[\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right] \cdot \vec{x}\right)
$$

## Fast Fourier Transform



$$
\begin{aligned}
& D\left(k_{n}, \vec{x}\right)=\sum_{\vec{k}_{a_{n}}} \Delta\left(\vec{k}_{a_{n}}\right) \exp \left(-i \vec{k}_{a_{n}} \cdot \vec{x}\right) \\
& N_{\text {tri }}=\frac{1}{N_{\mathrm{g}}^{3}} \sum_{\vec{x}} I\left(k_{1}, \vec{x}\right) I\left(k_{2}, \vec{x}\right) I\left(k_{3}, \vec{x}\right)
\end{aligned}
$$

$$
I\left(k_{n}, \vec{x}\right)=\sum_{\vec{k}_{a_{n}}} \exp \left(-i \vec{k}_{a_{n}} \cdot \vec{x}\right)
$$

Total 6 FFTs are required.
$\mathrm{O}\left(N^{3} \log N^{3}\right)$ time-complexity

## Bispectrum Parameterization

We design the estimator to compute the bispectrum for the triangles of all possible shapes and sizes.

Size : $k_{1}$ (length of the largest side)
Shape : $\left\{\begin{array}{l}t=k_{2} / k_{1} \\ \mu=\cos \theta=-\left(\boldsymbol{k}_{1} \cdot \boldsymbol{k}_{2}\right) /\left(k_{1} k_{2}\right)\end{array}\right.$

$$
B\left(k_{1}, k_{2}, k_{3}\right) \leftrightarrow B\left(k_{1}, \mu, t\right)
$$

Unique triangles (Bharadwaj+ 20)

$$
k_{1} \geq k_{2} \geq k_{3}
$$



## Unique shaped triangles



$4^{\text {th }}$ September, 2021

## Binning scheme

Average over all possible orientation of a triangle in 3D Fourier space.
use circular rings/spherical shells in $\boldsymbol{k}$-space.


We compute binned bispectra using combinations of three bins which satisfy the unique triangle conditions.


## Sampling $\mu$-t plane

$$
\begin{aligned}
t & =\frac{k_{2}}{k_{1}} \\
\mu & =\frac{1}{2}\left[\frac{k_{1}}{k_{2}}+\frac{k_{2}}{k_{1}}-\frac{k_{3}}{k_{1}} \frac{k_{3}}{k_{2}}\right]
\end{aligned}
$$

Increasing more number of shells will increase the sampling of $\mu$-t plane.

However smaller shell thickness increase the fluctuations in the estimated binned-bispectra.

The computation time will also increase due to more numbers of shells.


## Performance



Highly time efficient! $\quad N_{g}=1000--->300 \sec (F B E) \mid \quad 1.5 \times 10^{8} \sec (D B E)$

## Pros and Cons

## Bottleneck beyond N/3

If one of the three vectors crosses $N / 3$ on the Fourier grid there could be several instances where the sum of the three vectors will hit the other boundary of the box


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## Bottleneck beyond N/3

If one of the three vectors crosses $N / 3$ on the Fourier grid there could be several instances where the sum of the three vectors will hit the other boundary of the box

The sum then will wrap around the box and map back to the origin due to periodicity in the Fourier space
$\delta_{\mathrm{K}}\left(\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right)=\frac{1}{N_{\mathrm{g}}^{3}} \sum_{\vec{x}} \exp \left(-i\left[\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}\right] \cdot \vec{x}\right)$
The factor in the exponent will become integral multiple of $2 \pi$

$$
\vec{k}_{a_{1}}+\vec{k}_{a_{2}}+\vec{k}_{a_{3}}=\vec{q}\left(\frac{2 \pi N_{g}}{L}\right)
$$

## Excess triangles !

The FFT method predicts more number of triangles (due to PBC) than the true number of closed triangles

Counted exact number of triangles in equilateral triangle bin using direct method estimator.

For bispectrum the computation must be restricted within $\left(2 \pi N_{g}\right) /(3 L)$

For any $p$-th order spectrum $k<\left(2 \pi N_{\mathrm{g}}\right) /(p L)$

## Validating FBE

## Second-order biased NG field

- Generated Gaussian random field for a given power-law power spectrum $P(k)=A k^{n}$.

$$
\text { e.g. } \quad A=1, n=-2
$$

- Introduced non-Gaussianity using :

$$
\delta(\vec{x})=\delta_{\mathrm{G}}(\vec{x})+f_{\mathrm{NG}}\left[\delta_{\mathrm{G}}^{2}(\vec{x})-\left\langle\delta_{\mathrm{G}}^{2}(\vec{x})\right\rangle\right]
$$

- Compared the FBE results against the tree-level analytical predictions (for mild NG)

$$
B_{\mathrm{Ana}}\left(k_{1}, k_{2}, k_{3}\right)=2 f_{\mathrm{NG}}\left[P\left(k_{1}\right) P\left(k_{2}\right)+P\left(k_{2}\right) P\left(k_{3}\right)+P\left(k_{3}\right) P\left(k_{1}\right)\right]
$$

## Results-I (Real space)

Mean bispectrum computed over an ensemble of 100 realizations

$$
f_{\mathrm{NG}}=0.5 ; \quad N_{\mathrm{g}}=384
$$

## Results-I (Real space)

Compared with Analytical CV $\quad \sigma_{B}^{2}=\frac{1}{N_{\text {tri }}}\left[V P\left(k_{1}\right) P\left(k_{2}\right) P\left(k_{3}\right)+3 B^{2}\left(k_{1}, k_{2}, k_{3}\right)\right]$


## Results-I (Real space)

$$
\begin{aligned}
& f_{\mathrm{NG}}=0.5 \\
& N_{\mathrm{g}}=384
\end{aligned}
$$

Error bars are obtained from 100 realizations

Numerical predictions are consistent well within $5 \sigma$ with the analytical predictions

Anomaly is found at a bin near the squeezed limit


## LoS Anisotropy

- We introduce the LoS anisotropy (in Fourier space) to the previous NG field using

$$
\Delta^{s}(\vec{k})=\left(1+\beta_{1} \mu_{1}^{2}\right) \Delta^{r}(\vec{k}) \quad \text { where } \quad \beta_{1}=1
$$

- Compared the FBE monopole moment estimates against the analytical predictions (for mild NG)

$$
B_{0}^{0}\left(k_{1}, \mu, t\right)=B_{\mathrm{Ana}}\left(k_{1}, \mu, t\right) \times\left[1+\beta_{1}+\frac{3 \beta_{1}^{2}}{5}+\frac{\beta_{1}^{3}}{7}-\frac{4 \beta_{1}^{2}\left(3 \beta_{1}+7\right)\left(1-\mu^{2}\right)\left(t^{2}-\mu t+1\right)}{105\left(t^{2}-2 \mu t+1\right)}\right]
$$

(Bharadwaj+ 20)

## Results-I (Redshift space)

Mean bispectrum computed over an ensemble of 100 realizations

$$
f_{\mathrm{NG}}=0.5 ; \quad N_{\mathrm{g}}=384
$$

## Results-I (Redshift space)

$$
\begin{aligned}
& f_{\mathrm{NG}}=0.5 \\
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$$

Error bars are obtained from 100 realizations

Numerical predictions are consistent well within $5 \sigma$ with the analytical predictions

Anomaly is found at a bin near the squeezed limit


## Application to DM Density Field

## Gravitationally-induced NG

- Dark matter densities at $z=3,2,1$. (50 realizations)
- Box volume $\mathrm{V}=[512 \mathrm{Mpc}]^{3}$
- Grid size --> [2 Mpc]
- Computed the bispectrum using FBE and compare it against the $2^{\text {nd }}$ order perturbation theory (2PT) predictions.
(Bharadwaj+ 20)

$$
\begin{aligned}
B\left(k_{1}, k_{2}, k_{3}\right)=2[ & F_{2}\left(\boldsymbol{k}_{1}, \boldsymbol{k}_{2}\right) P\left(k_{1}\right) P\left(k_{2}\right) \\
& +F_{2}\left(\boldsymbol{k}_{2}, \boldsymbol{k}_{3}\right) P\left(k_{2}\right) P\left(k_{3}\right) \\
& \left.+F_{2}\left(\boldsymbol{k}_{3}, \boldsymbol{k}_{1}\right) P\left(k_{3}\right) P\left(k_{1}\right)\right]
\end{aligned}
$$

$$
\begin{aligned}
& F_{12}(\mu, t)=\frac{1}{14}\left(4 \mu^{2}-7 \mu t-\frac{7 \mu}{t}+10\right) \\
& F_{23}(\mu, t)=\frac{7 \mu+\left(3-10 \mu^{2}\right) t}{14 t s^{2}}, \\
& F_{31}(\mu, t)=\frac{t^{2}\left(-10 \mu^{2}+7 \mu t+3\right)}{14 s^{2}} .
\end{aligned}
$$

## Results-II (Real Space)



Qualitatively matches with the 2PT predictions and quantitatively close

## Results-II

Largely dominated by cosmic variance at small $k_{1}$ values and higher $z$.

Fluctuations are random and small wherever 2PT is expected to be valid.

There is a significant positive deviations towards larger $k_{1}$ and lower $z$.


## Comparing with the Direct Bisp. Estim.



Bispectra for the dark matter density field matches well with the DBE (Majumdar+18) within 1 $\sigma$

## Take Home

- The FBE has $\mathrm{O}\left(N^{3}\right)$ improvement in speed over the direct method.
- Considers all unique possible shapes of triangles
- Consistent results in case of 2nd order biased non-Gaussian fields and also for the dark matter density fields. Deviations from analytical predictions are within $\sim 3 \sigma$
- Can also be applied to EoR/ Post-EoR signal (Ongoing).
- With proper modifications to FBE one can also compute the different multipole moments of the bispectrum to quantify the anisotropy such as RSD effects.

