Observation and astrophysics of gravitational waves Current status and future prospects

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Anuradha Gupta The University of Mississippi

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LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars



LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars



Open Gravitational-wave Catalog 90 2020 **GWTC-3**





LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



Questions:

- How these binary mergers form?
- mergers take place?
- Is there a lower mass gap?
- Is there a higher mass gap?
- New observations, new questions!

What are the astrophysical environments in which these

• What is the rate of star formation throughout the Universe?

Quasi-circular orbits

Intrinsic parameters: $m_1, m_2, \chi_1, \theta_1, \phi_1, \theta_2, \phi_2$

Extrinsic parameters: D_L , ι , δ , α , ψ , t_c , ϕ_c



x

Quasi-circular orbits

Intrinsic parameters: $m_1, m_2, \chi_1, \theta_1, \phi_1, \theta_2$

Chirp mass: M

$$\mathscr{M}_c = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

mass ratio:

$$q = \frac{m_2}{m_1}$$

effective inspiral spin: $\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + \chi_2 \cos \theta_1 + \chi_2 \cos \theta_1 + \chi_2 \cos \theta_1 + \chi_2 \cos \theta_2 + \chi_2 \cos \theta_1 + \chi_2 \cos \theta_2 + \chi_2 \sin \theta_2 + \chi_2 \sin$ $m_1 +$

effective precessing spin: $\chi_p = \max | \chi_1 \sin \theta_1,$

$$-\frac{m_1\chi_1\cos\theta_2}{p_1}, \frac{3+4q}{q_1+3q}, q_{\chi_2}\sin\theta_2$$





Analysis Setup:

- GWTC-3:
 - events from O1, O2 and O3
 - 90 CBCs with $P_{\rm astro} > 0.5$
- With FAR $< 0.25 \text{ yr}^{-1}$ in at least one detection pipeline: 2 BNSs, 2 NSBHs and 63 BBHs
 - High purity set of events whose selection biases are understood
 - Expect only 1 event to be not of astrophysical origin
- With FAR < 1 yr^{-1} : 69 BBHs
 - relative proportion of background events remains below 10%
 - Expect only 4.6 events to be not of astrophysical origin
- NS and BH are distinguished based on EOS maximum NS mass

[More on detection and parameter estimation in Surabhi's Talk]



Population Analysis Framework

- Hierarchical Bayesian approach
- Marginalize over the uncertainty in the estimate of individual event parameters

$$\mathscr{L}(\{d\}, N_{\text{det}} | \Lambda, N_{\text{exp}}) \propto N^{N_{\text{det}}} e^{-N_{\text{exp}}} \prod_{i=1}^{N_{\text{det}}} \int \mathscr{L}(d_i | \theta) \pi$$
$$N = \frac{N_{\text{exp}}}{\xi(\Lambda)}$$

 $\xi(\Lambda)$: fraction of mergers that are detectable for a population with parameter Λ

 $\mathscr{L}(d_i \mid \theta)$: single-event likelihood described by parameter set θ

 $\pi(\theta \mid \Lambda)$: prior governing the population distribution on event parameters θ for a given value of hyperparameters Λ

Marginalizing the likelihood after imposing the a log-uniform prior on N

$$\mathscr{L}(\{d\} \mid \Lambda) \propto \frac{\prod_{i=1}^{N_{\text{det}}} \int \mathscr{L}(d_i \mid \theta) \pi(\theta \mid \Lambda) d\theta}{\xi(\Lambda)}$$

 $z(\theta \mid \Lambda)d\theta$

 $\mathscr{L}(d_i | \theta)$ are computed using default prior $\pi_{\mathcal{O}}(\theta)$, hence

$$\mathscr{L}(\{d\} \mid \Lambda) \propto \frac{\prod_{i=1}^{N_{det}} \frac{1}{\xi(\Lambda)} \left\langle \frac{\pi(\theta \mid \Lambda)}{\pi_{\emptyset}(\theta)} \right\rangle$$

Using this posterior on Λ can be computed





Population Model Assumptions



Parametric Mass Models: NS masses

- •For analyses exclusively focused on the NS-containing events
 - POWER: Power Law
 - •PEAK: Gaussian
- •With sharp minimum and maximum mass cutoffs
- •Components of BNSs are drawn independently from the common NS mass distribution
- •For NSBHs, we assume a uniform BH mass distribution and random pairing with NSs





Parametric Mass Models: BH masses

Fiducial Power Law + Peak (PP) model + redshift evolution

$$p(m_1, q, z) \propto q^{\beta} p(m_1)(1 + z)^{\kappa - 1}$$

Merger rate normalization is chosen such that

$$\mathscr{R}(z) = \frac{dN}{dV_c dt}(z) = \mathscr{R}_0(1+z)^{\kappa}$$

The corresponding redshift distribution

$$p(z \mid \kappa) \propto \frac{1}{1+z} \frac{dV_c}{dz} (1+z)^{\kappa}$$

Parametric Mass Models: BH masses

- Power Law + Dip + Break model (PDB)
- power law with a notch filter

$$p(m \mid \lambda) = n(m \mid M_{low}^{gap}, M_{high}^{gap}, A) \times l(m \mid m_{max}, \eta)$$

$$\begin{cases} m^{\alpha_1} \text{ if } m < M_{high}^{gap} \\ m^{\alpha_2} \text{ if } m > M_{high}^{gap} \\ 0 \text{ if } m > m_{max} \text{ or } m < m_{min} \end{cases}$$

- Two pairing probability:
 - Random: $p(m_1, m_2 | \Lambda) \propto p(m = m_1 | \Lambda) p(m = m_2 | \Lambda) \Theta(m_2 < m_1)$ ullet
 - ullet

To fit the distribution of BH and NS masses, we use a parameterized model, consisting of a broken



power-law-in-mass-ratio: $p(m_1, m_2 | \Lambda) \propto p(m = m_1 | \Lambda) p(m = m_2 | \Lambda) q^{\beta} \Theta(m_2 < m_1)$



BH Spin Models

- **Default spin model**
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Spin magnitude drawn from Beta distribution
$$p(\chi_i | \alpha_{\chi}, \beta_{\chi}) \propto \chi_i^{\alpha-1} (1-\chi_i)^{\beta-1}$$

Spin tilts are mixture of isotropic and aligned
 $p(\cos \theta_i | \zeta, \sigma_i) = \frac{1}{2} (1-\zeta) + \zeta \mathcal{N}_{[-1,1]}(\cos \theta_i; 1, \sigma_i)$
Aussian spin model
Bivariate Gaussian in χ_{eff} and χ_p
 $p(\chi_{\text{eff}}, \chi_p | \mu_{\text{eff}}, \sigma_{\text{eff}}, \mu_p, \sigma_p, r) \propto \mathcal{N}(\chi_{\text{eff}}, \chi_p | \mu, \Sigma)$
 $\mu = (\mu_{\text{eff}}, \mu_p)$ and $\Sigma = \begin{pmatrix} \sigma_{\text{eff}}^2 & r\sigma_{\text{eff}}\sigma_p \\ r\sigma_{\text{eff}}\sigma_p & \sigma_p^2 \end{pmatrix}$
Sumplifying the state of the s



Nerger Rates



	BNS		NSBH	BBH	NS-Gap	BBH-gap	Full
	$m_1 \in [1, 2.5] M_{\odot}$	r	$n_1 \in [2.5, 50] M_{\odot}$	$m_1 \in [2.5, 100] M_{\odot}$	$m_1 \in [2.5, 5] M_{\odot}$	$m_1 \in [2.5, 100] M_{\odot}$	$m_1 \in [1, 100]$ M
	$m_2 \in [1, 2.5] M_{\odot}$		$m_2 \in [1, 2.5] M_{\odot}$	$m_2 \in [2.5, 100] M_{\odot}$	$m_2 \in [1, 2.5] M_{\odot}$	$m_2 \in [2.5,5] M_{\odot}$	$m_2 \in [1, 100]$ M
PDB (pair)	960^{+1740}_{-700}		59^{+81}_{-38}	$25^{+10}_{-7.0}$	41^{+69}_{-30}	$9.3^{+18.7}_{-7.6}$	1100^{+1700}_{-750}
PDB (ind)	250^{+640}_{-196}		170^{+150}_{-89}	$22^{+9.0}_{-6.0}$	29^{+55}_{-23}	$10^{+15}_{-8.0}$	470^{+830}_{-300}
\mathbf{MS}	470^{+1430}_{-413}		57^{+123}_{-42}	42^{+88}_{-20}	$3.7\substack{+20.3 \\ -3.4}$	$0.17\substack{+55.83 \\ -0.16}$	$650\substack{+1550 \\ -460}$
BGP	99^{+260}_{-86}		32^{+62}_{-25}	33^{+16}_{-10}	$2.1\substack{+33.0 \\ -2.1}$	$5.1^{+12.0}_{-4.0}$	$180\substack{+260 \\ -110}$
Merged	13-1900		7.4-320	16-130	0.029-84	0.0095-56	71-2200

- Self consistent measure of merger rate of all detected CBCs
- Subdivided into astrophysically interesting mass ranges
- Constant co-moving volume merger rate density
- BBH merger rates are consistent with previously published estimates
- For BNS/NSBH, the inferred rate depends on the presumed mass distribution
- Different approaches arrive at different binary mass distributions between $1\,M_{\odot}$ and $2.5\,M_{\odot}$
- Highlights the importance of modeling systematics when drawing inferences about populations with few confident members





- ° NSs are expected to have a maximum mass of ~3 M_{\odot}
- $^{\circ}$ Heaviest NS observed to date has mass 2.01±0.04 M_{\odot}
- ° PSR J0740+6620 'may' host a ~2.07-2.08 M_{\odot} NS
- X-ray observations found a dearth of objects in the $3 - 5 M_{\odot}$ range
- Population synthesis models with rapid supernova instabilities can not produce BHs of mass < 5 M_{\odot}
- Thus, we don't expect to observe binaries with component masses in ~3-5 M_{\odot} in GWs
 - provided they are formed in a way similar to the binaries observed with EM radiation

Lower mass gap





- A reduction in the rate above NS masses
- Neither find evidence for nor rule out the existence of a two-sided lower mass gap
- objects
- If a lower mass gap does exists, it may not be totally empty

• Unable to confidently infer absence or presence of a subsequent rise in merger rates from lower mass gap masses

• Rates in mass gap are 1 or 2 order of magnitude different; GW data suggest two distinct populations of compact



Possible ways to populate this gap

- Remnants of binary neutron stars 0 (Gupta et al, 2020)
- Hierarchical merger of stellar few-body Systems (Safarzadeh et al, 2020, Lu et al, 2020)
- BHs of primordial origin (García-Bellido, 2019) 0
- Still a mystery that what the secondary of GW190814 is?
- Primary of GW190425 ($m_1 \in [1.61, 2.52]M_{\odot}$) could be a low-mass gap BH

LVC/@LIGO/@EGO_Virgo/@Daniel_Williams



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Masses

NS masses in binaries



- maximum NS mass > 50%
- • $M_{\text{max,TOV}} = 2.21^{+0.31}_{-0.21} M_{\odot}$ from 2106.05313 based on pulsar timing, GW and X-ray observations of NSs

Name	$FAR_{min} (yr^{-1})$	$P(m < M_{\max, TOV})$	$P(m < M_{ m low}^{ m gap})$	Classification
GW170817	$< 1 \times 10^{-5}$	0.99	0.98	BNS
GW190425	$3.38{ imes}10^{-02}$	0.68	0.73	BNS
CW190814	$< 1 \times 10^{-5}$	0.06	0.19	BBH
GW200105	$2.04{ imes}10^{-01}$	0.94	0.74	NSBH
GW200115	$< 1 \times 10^{-5}$	0.95	0.97	NSBH
GW190426	$9.12{ imes}10^{-01}$	0.82	_	NSBH
GW190917	$6.56 imes 10^{-01}$	0.56	_	NSBH

• Classification criteria: probability that at least one of the component masses is less than the







Power: $\alpha = -2.0^{+5.1}_{-7.0}$

Peak: $\sigma = 1.1^{+0.8}_{-0.8} M_{\odot} \text{ and } \mu = 1.5^{+0.4}_{-0.4} M_{\odot}$

- •GW observations to date do not support a NS mass distribution with a pronounced single peak
- In contrasts with the Galactic BNS subpopulation whose mass distribution is sharply peaked around $1.35\,M_{\odot}$
- Mass distribution of NSs observed in GWs is broader and has greater support for high-mass NSs
- Galactic NS population distribution has a double-peaked shape

$$m_{\rm min} = 1.2^{+0.1}_{-0.2} M_{\odot}$$
 $m_{\rm max} = 2.0^{+0.3}_{-0.2} M_{\odot}$







GW190814 is an outlier from the secondary masses in BNS and NSBH systems

 $\max_5(m_2)$: the largest observed secondary mass after 2 BNS observations and 3 NSBH observations







BH masses in binaries

FAR threshold of $1 \ yr^{-1}$ and redshift dependance



BH mass distribution is consistent with GWTC-2











- •The inferred mass spectrum decays more rapidly
- Expected because new observations in GWTC-3 contain a greater fraction of lower mass systems
- The inferred mass ratio distribution is less peaked towards equal mass binaries compared to GWTC-2
- •Inconclusive evidence for an upper mass gap (~50-150 M_{\odot})





High-mass gap: what theory says?



Star Mass

No PPISN

PPISN Pulsational pair instability supernovae (PPISN)

PISN

Pair instability supernovae (PISN)

Inconclusive evidence for upper mass gap

- A gap is defined as a rapid decline and a rapid rise in merger rate at significantly higher mass
- No evidence is found for such a gap
- The PPISN mass gap could start at higher mass than theory expects
- Or the high-mass binaries in GWTC-3 could be formed in a way that avoids **PPISN**





Possible ways to populate this gap

- Hierarchical BBH mergers
- Stellar mergers
- BH remnants of population III stars
- Stellar triples in the field of the galaxy
- Growth via accretion in an Active Galactic Nuclei (AGN)

[More in Vishal and Parthapratim's Talks]





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2.0

GW190814 and GW190917 are outliers from the secondary masses in BBH systems

 $p(m_{
m min}) \, [M_{\odot}^{-1}$ 1.50.5



m_{min}: minimum recovered BH mass



Evolution of rates with redshift

$$\mathcal{R}(z) = \frac{dN}{dV_c dt}(z) = \mathcal{R}_0(1 + dV_c dt)$$

- Initially there was a preference for a rate that increased with redshift but still consistent with a non-evolving merger rate
- But with GWTC-3 we are confident that BBH rate is evolving with redshift

$(+ z)^{\kappa}$



Spin distribution in BBHs



BBH spin distribution is consistent with GWTC-2

 ${\cal M}[M_\odot]$







favor a broad or isotropic distribution of spin tilts



Again excluded the case of perfect spin–orbit alignment and now data strongly





- vanishing spins, with a mean centered at $0.06^{+0.04}_{-0.05}$
- $\chi_p = 0$, or a narrow distribution centered at $\chi_p \approx 0.2$

• As with GWTC-2 we again infer a $\chi_{\rm eff}$ distribution compatible with small but non-

• χ_p measurements can be explained either by a broad distribution centered at





Evidence of extreme spin-orbit misalignment

- Spin tilts $> 90^{\circ}$ unlikely for BBH formation from isolated stellar progenitors (Kalogera, 2000)
- $\chi_{\rm eff} < 0$ would serve as a strong indicator of dynamical interaction during BBH evolution
- Extended the Gaussian model to truncate χ_{eff} on the range $\chi_{\rm eff,min} \leq \chi_{\rm eff} \leq 1$ (rather than $-1 \leq \chi_{\rm eff} \leq 1$) and hierarchically measured $\chi_{\rm eff,min}$
- $\chi_{\rm eff,min} < 0$ at 99.8% credibility
- There were objection that unless BBH spin models are expanded to allow the existence of a secondary subpopulation with vanishingly small spins and spins allowed to correlate with other BBH parameters like the mass ratio
- Repeat the inference of $\chi_{\rm eff,min}$ but under an expanded model that allows for a narrow subpopulation of BBHs with extremely small $\chi_{\rm eff}$
- Data still prefer a negative $\chi_{eff,min}$ but with lower significance, $\chi_{\rm eff,min} < 0$ at 88.4% credibility





- Mass and spin distributions are consistent with GWTC-2
- More BBHs with preferentially negatively aligned spins but it could be by chance as well
- The BBH merger rate density increases with redshift
- A relative dearth of observations with component masses between 3 and 5 M_{\odot}
- No strong evidence for lower and higher mass gaps
- The inferred NS mass distribution, albeit based on a limited sample of observations, does not exhibit a peak at $1.35 M_{\odot}$
- GW190814 is an outlier for the secondary mass of BNS/NSBH and BBH population

Summary







- LIGO-Virgo-KAGRA
 - Stellar mass NS/BH merger up to redshift $z\sim 1$
- 2G: Voyager (Livingston, Hanford, Pisa, Hingoli, Kamioka)
 - 4-5 times better sensitivity, detections up tp redshift $z\sim 10$
- 3G: Einstein Telescope, Cosmic Explorer
 - 20 times better sensitivity, detections up tp redshift $z\sim 30$
- Other: LISA, TianQin, DECIGO, Pulsar Timing array
 - Explore lower frequency regime and new sources



The properties of the second second

