

Asymptotic Safety in Quantum Gravity

(Martin Reuter)

The 4 Fundamental Interactions : theoretical status

electromagnetic }
weak }
strong } : { class.: Yang-Mills theories
quant.: perturbatively
renormalizable QFTs

gravity: { class.: General Relativity, with
 $S[g_{\mu\nu}] \sim \int d^4x \sqrt{-g} R$
quant.: ???

Quantum field theory of spacetime metric $g_{\mu\nu}(x)$ based upon $\int \sqrt{-g} R$ is not renormalizable in perturbation theory:

increasing order in pert. th. $\Rightarrow$
" number of counter terms $\Rightarrow$
" " " undetermined parameters

$\leadsto$ no / very questionable predictivity at high energies
("effective" rather than "fundamental" theory)

Standard quantization of gravity $\hat{=}$

degrees of freedom
carried by :

$$g_{\mu\nu}(x)$$

bare action:

$$\int d^4x \sqrt{-g} R$$

calculational method:

perturbation theory in G ;
infinite cutoff limit at the
trivial "Gaussian" fixed point

What should be given up in order to arrive at
a "fundamental" or "microscopic" quantum theory
of gravity?

String Theory: d.o.f., action, calc. meth.

Loop Quantum Gravity: d.o.f., calc. meth.

Asymptotic Safety: calc. meth., action

Asymptotic Safety Approach:

~> degrees of freedom carried by $g_{\mu\nu}$

~> quantization/renormalization is non-perturbative in an essential way

~> bare action Γ_* is not an ad hoc assumption, but a prediction:

$$\Gamma_* \sim \int d^4x \sqrt{-g} R + \text{"more"} \quad \text{is a}$$

non-Gaussian fixed point of the

(∞ -dimensional, non-pert.) Wilsonian

renormalization group flow

~> fixed point "controls" UV divergences

... provided it exists

Weinberg's "asymptotic safety" conjecture (1979):

Perhaps Quantum Einstein Gravity can be defined nonperturbatively at a non-Gaussian fixed point.

$d = 2 + \varepsilon$: FP known to exist

$d = 4$: progress hampered by lack of appropriate calculational scheme



Use "effective average action"

which seems ideally suited.

Wetterich 1993

Effective average action for gravity:

M.R. 1996

$$\Gamma_k [g_{\mu\nu}, \dots]$$

The Effective Average Action Γ_k

- Wilson-type (coarse grained) free energy functional
- IR cutoff at k : Γ_k contains the effect of all quantum fluctuations with momenta $p > k$, not (yet) of those with $p < k$.
- modes with $p < k$ suppressed in the path integral by $(\text{mass})^2 = R_k(p^2)$
- $\Gamma_{k \rightarrow \infty} = S$, classical (bare) action
- $\Gamma_{k \rightarrow 0} = \Gamma$, standard effective action
- Γ_k satisfies exact RG equation; symbolically:

$$k \partial_k \Gamma_k = \frac{1}{2} \text{Tr} \left[(\Gamma_k^{(2)} + R_k)^{-1} k \partial_k R_k \right]$$

- Powerful nonperturbative approximation scheme:

"truncate" the space of action functionals,
project RG flow onto finite dimensional
subspace



Construction of Γ_k for Gravity

- starting point: $\int \mathcal{D}\gamma_{\mu\nu} e^{-S[\gamma_{\mu\nu}]}$
- decompose $\gamma_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$
fixed backgrd.
metric
- add background gauge fixing $S_{gf}[h; \bar{g}]$ + ghost terms
- expand $h_{\mu\nu}$ in $\bar{D}^2$ -eigenmodes, and introduce IR cutoff k^2 : only modes with generalized momenta ($\bar{D}^2$ -eigenvalues) are integrated out.
- add sources: generating fctl. $W_k[\text{sources}; \bar{g}]$

Legendre transf.



$$g_{\mu\nu} \equiv \langle \gamma_{\mu\nu} \rangle$$

$$\Gamma_k[g_{\mu\nu}, \bar{g}_{\mu\nu}, \text{ghosts}]$$

- derive exact RG equation from path integral:

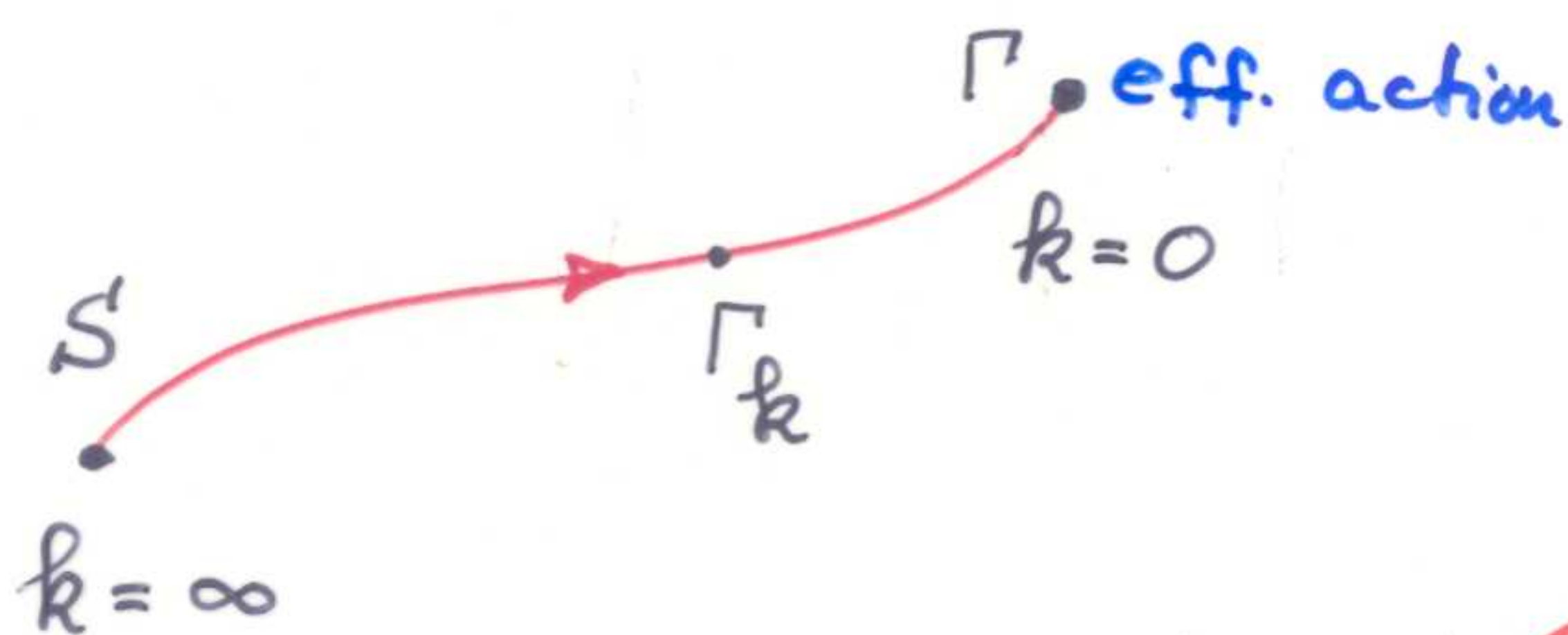
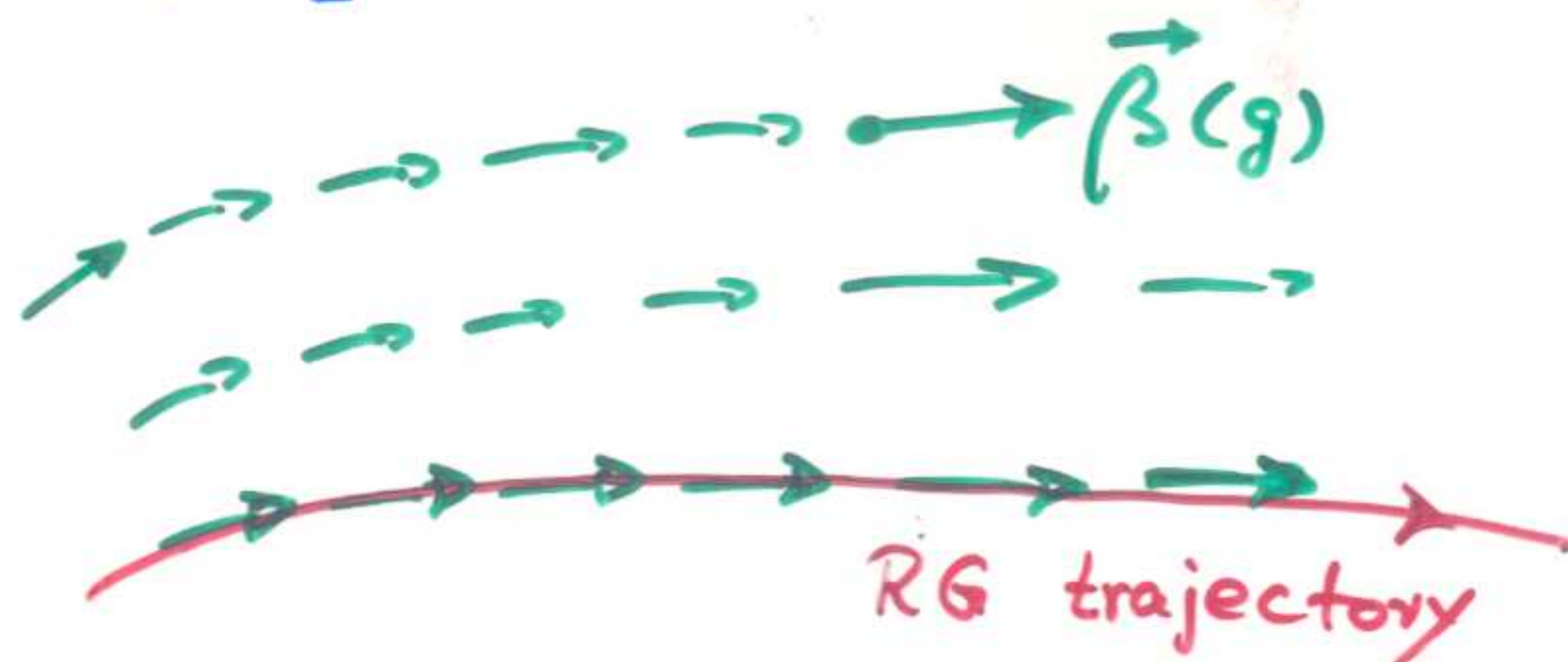
$$k \frac{\partial}{\partial k} \Gamma_k[g, \bar{g}, \dots] = \text{Tr}(\dots)$$

$$\begin{aligned} \Gamma_\infty &= S \\ \Gamma_0 &= \Gamma \end{aligned}$$

- "Ordinary" diffeomorphism invariant action:

$$\Gamma_k[g] = \Gamma_k[g, \bar{g}=g, \text{ghosts}=0]$$

• $A[\cdot]$



bare action
≡ fixed point Γ_*

Theory Space

The Einstein-Hilbert Truncation

(M.R., 1996)

ansatz:

$$\Lambda_k \equiv \bar{\lambda}_k$$

$$\Gamma_k = - \frac{1}{16\pi G_k} \int d^d x \sqrt{g} \{ R - 2\Lambda_k \}$$

+ classical gauge fixing and ghost terms

two running parameters:

Newton constant G_k , dimensionless: $g(k) = k^{d-2} G_k$

cosmological constant Λ_k , dimensionless: $\lambda(k) = \Lambda_k / k^2$

insert ansatz into flow equation, expand

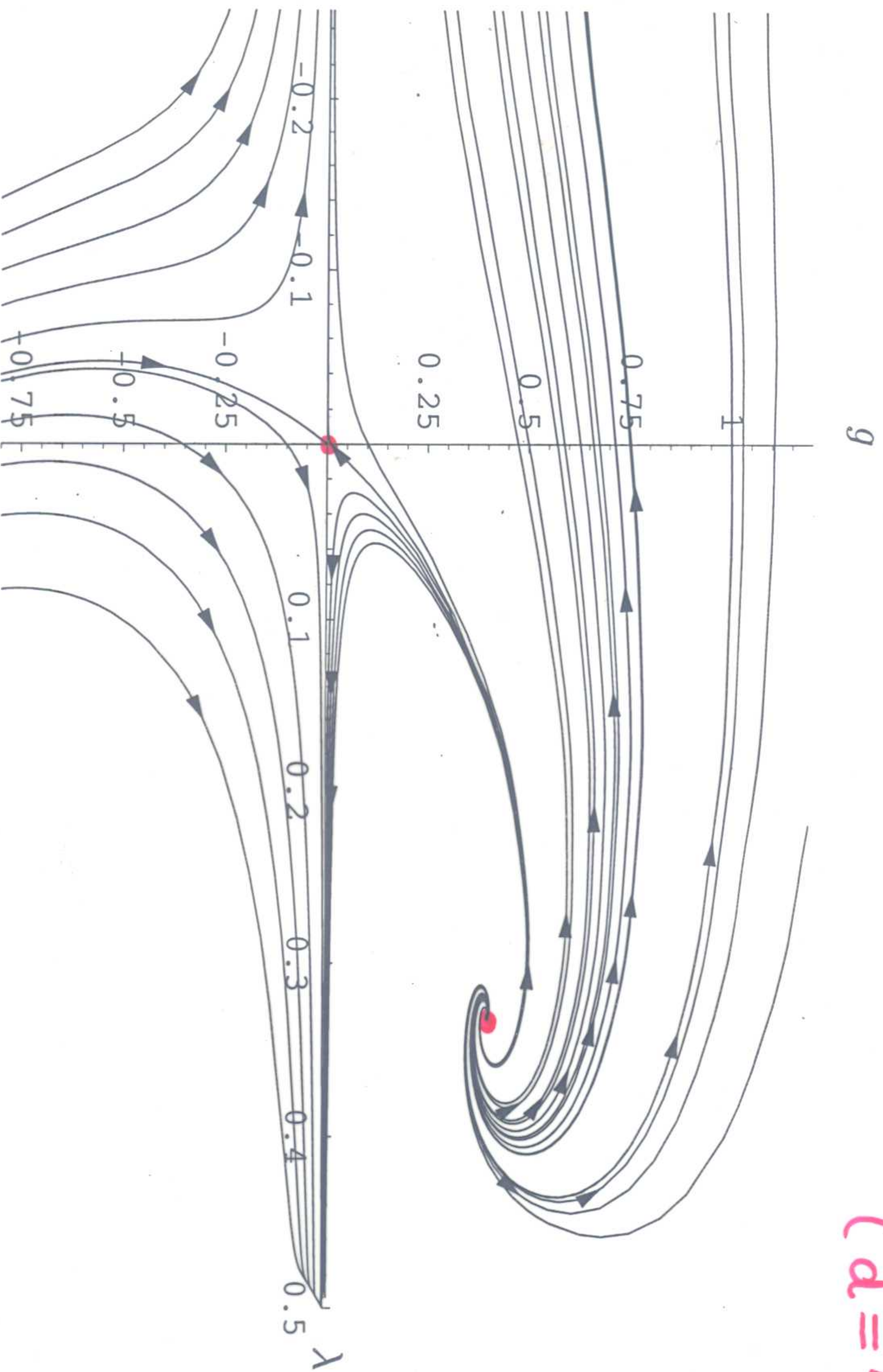
$$\text{Tr} [\dots] = (\dots) \int \sqrt{g} + (\dots) \int \sqrt{g} R + \dots$$

$$k \partial_k g(k) = \beta_g(g, \lambda)$$

$$k \partial_k \lambda(k) = \beta_\lambda(g, \lambda)$$

RG-Flow in the Einstein-Hilbert Truncation

($d=4$)



O. Lauscher, M.R.
F. Saueressig, M.R.
R. Percacci, ...
D. Litim,

Reliability checks (universality, ...),
more general truncations $\Rightarrow$

The NGFP seems to exist in
the full un-truncated theory
beyond any reasonable doubt.

Cosmological

Applications / Implications

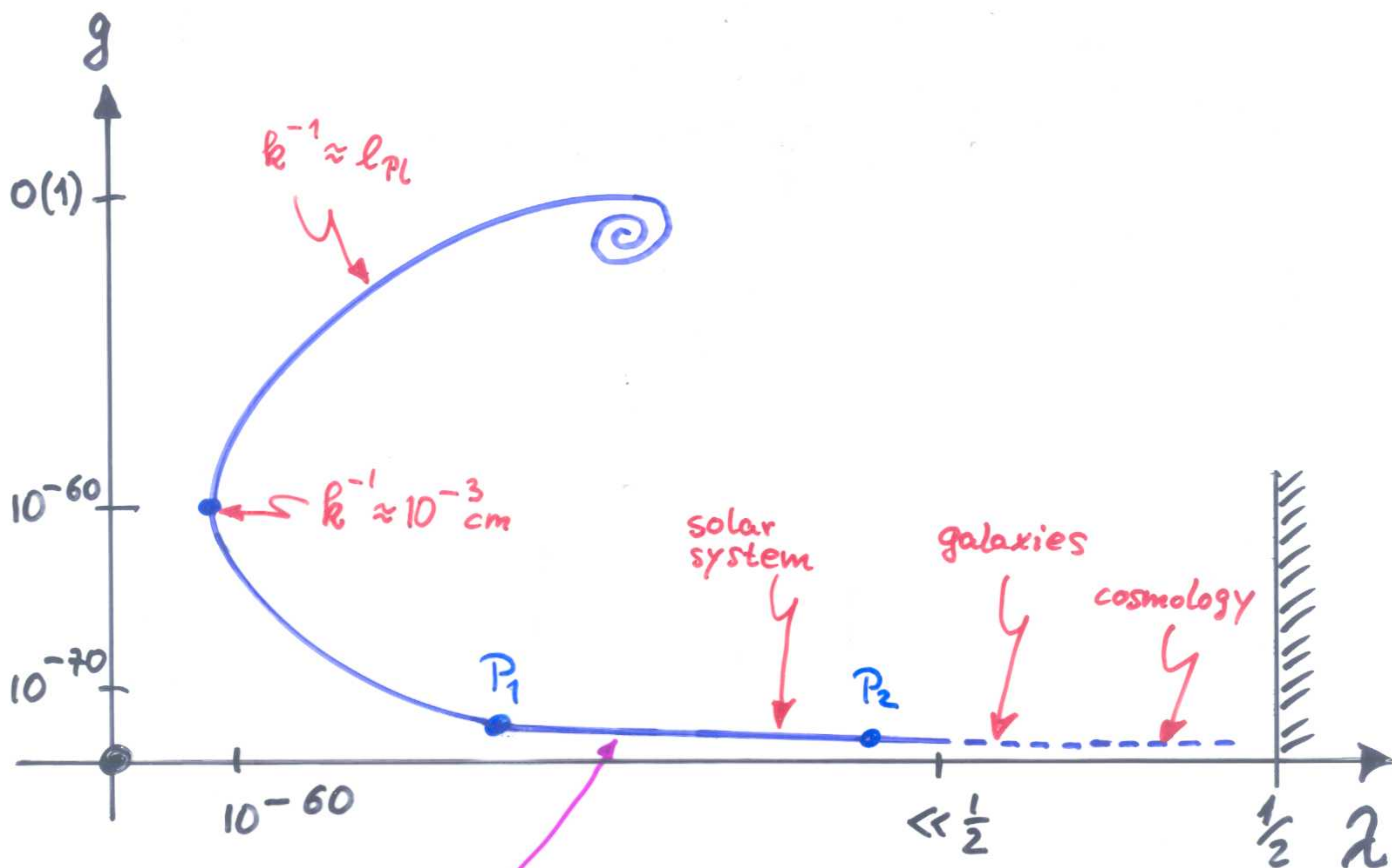
A. Bonanno, M.R. (2002)

M.R., H. Weyer (2005)

M.R., F. Saueressig (2005)

A. Bonanno, M.R. (2007)

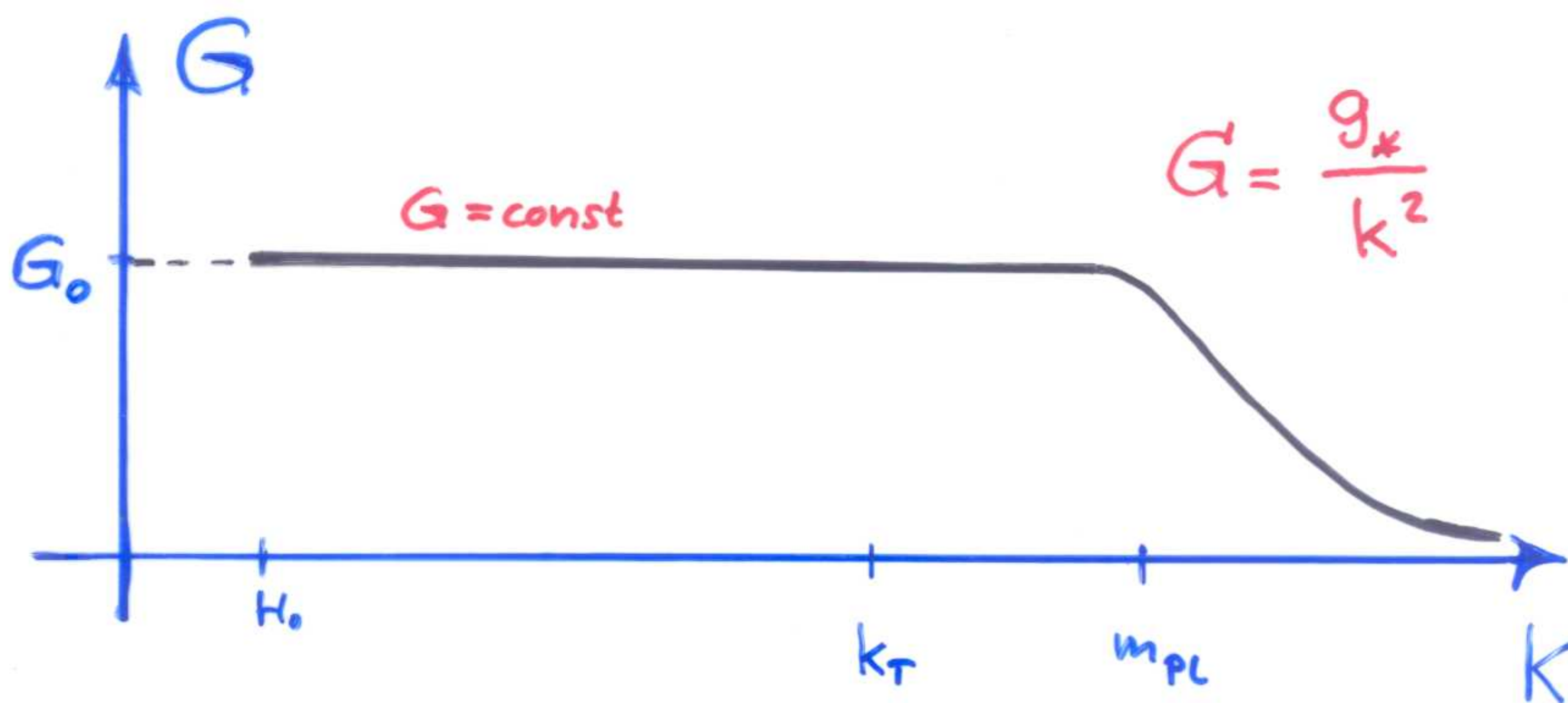
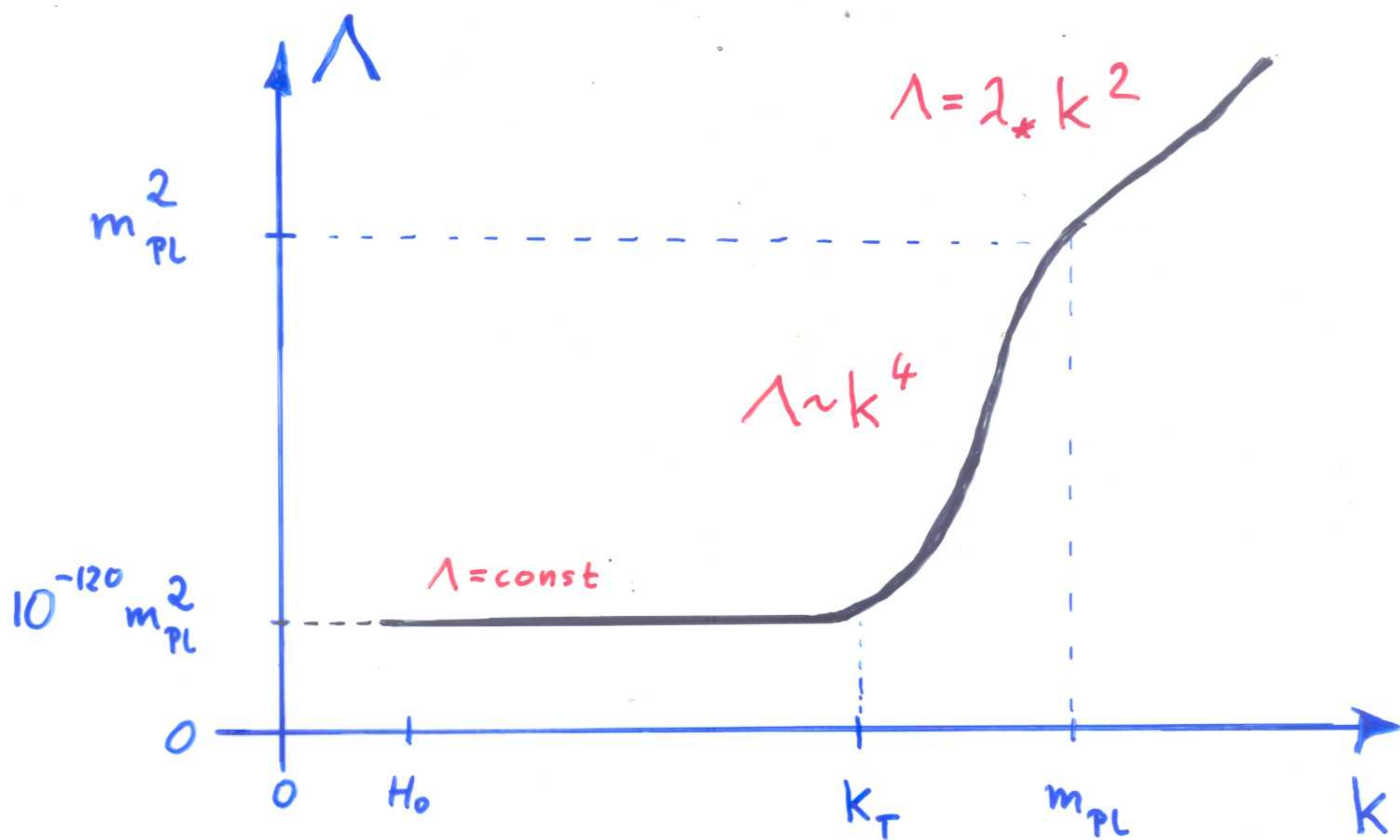
The RG trajectory "realized in Nature"



classical GR regime:
 $G, \Lambda \approx \text{const}$

"Today" in cosmology:

$$\lambda_{\text{cosmo}} \equiv \frac{\Lambda_{\text{cosmo}}}{k_{\text{cosmo}}^2} \approx H_0^2 = O(1) !!!$$



Definition of Planck scale: $m_{\text{Pl}} \equiv \ell_{\text{Pl}}^{-1} \equiv G_0^{-\frac{1}{2}}$

Are there observable / observed
physical phenomena related to
the RG-running implied by QEG ?

Candidates in cosmology:

a) Entropy carried by cosmological matter
(CMBR photons, ...)

$$[S_{\text{CMBR}} / \text{Hubble volume}]_{\text{today}} \approx 10^{88} \gg 1$$

most plausible initial value = $O(1)$!

b) Automatic inflation in the NGFP regime

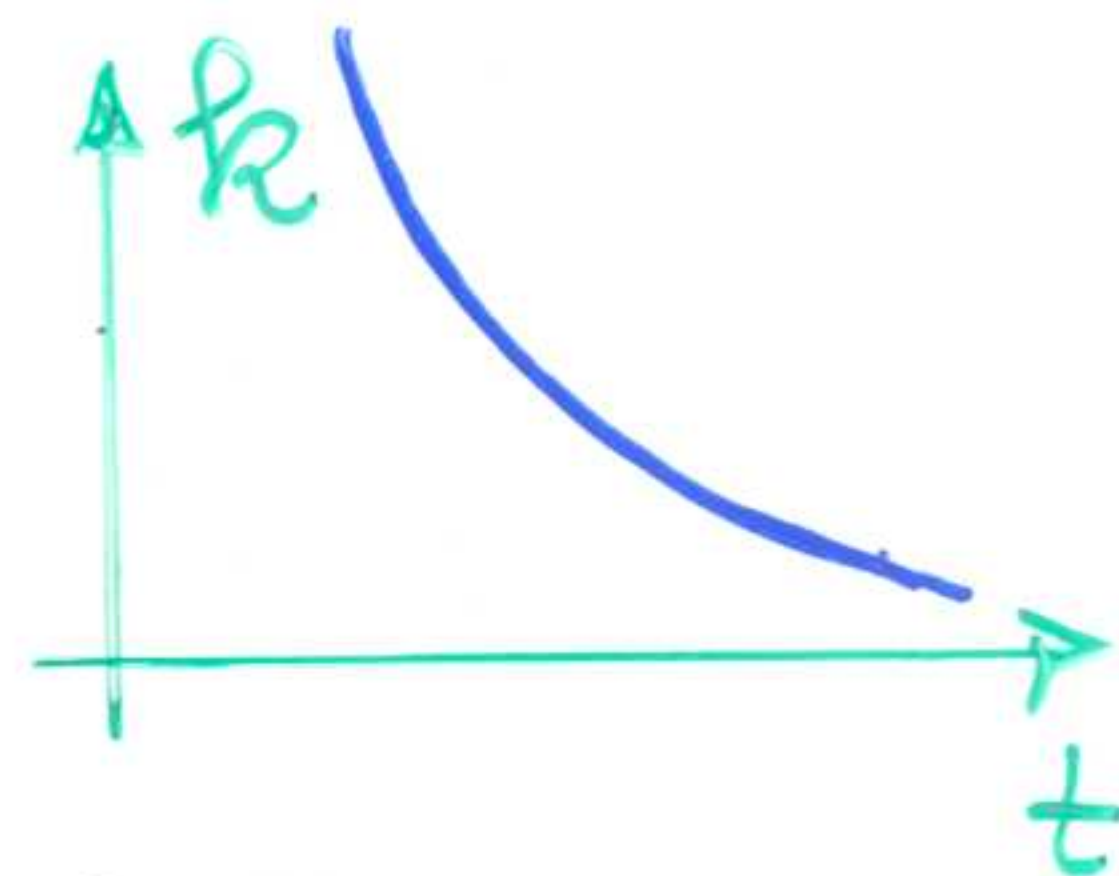
↑ no inflaton needed,
no reheating necessary;
 $\Lambda(k)$ large: cosmolog. const. drives inflation
 $\Lambda(k)$ small: inflation stops automatically

c) Generation of primordial density perturbations

spectrum scale free:

big bang = "critical phenomenon" governed
by the NGFP

For monotonic cosmological
cutoff identification $k = k(t)$



and $\odot$ below the NGFP regime:

$\Lambda(t) \equiv \Lambda(k = k(t))$ is a **positive**
and **decreasing** function of time.



Energy transfer into the matter
system ("heating up").

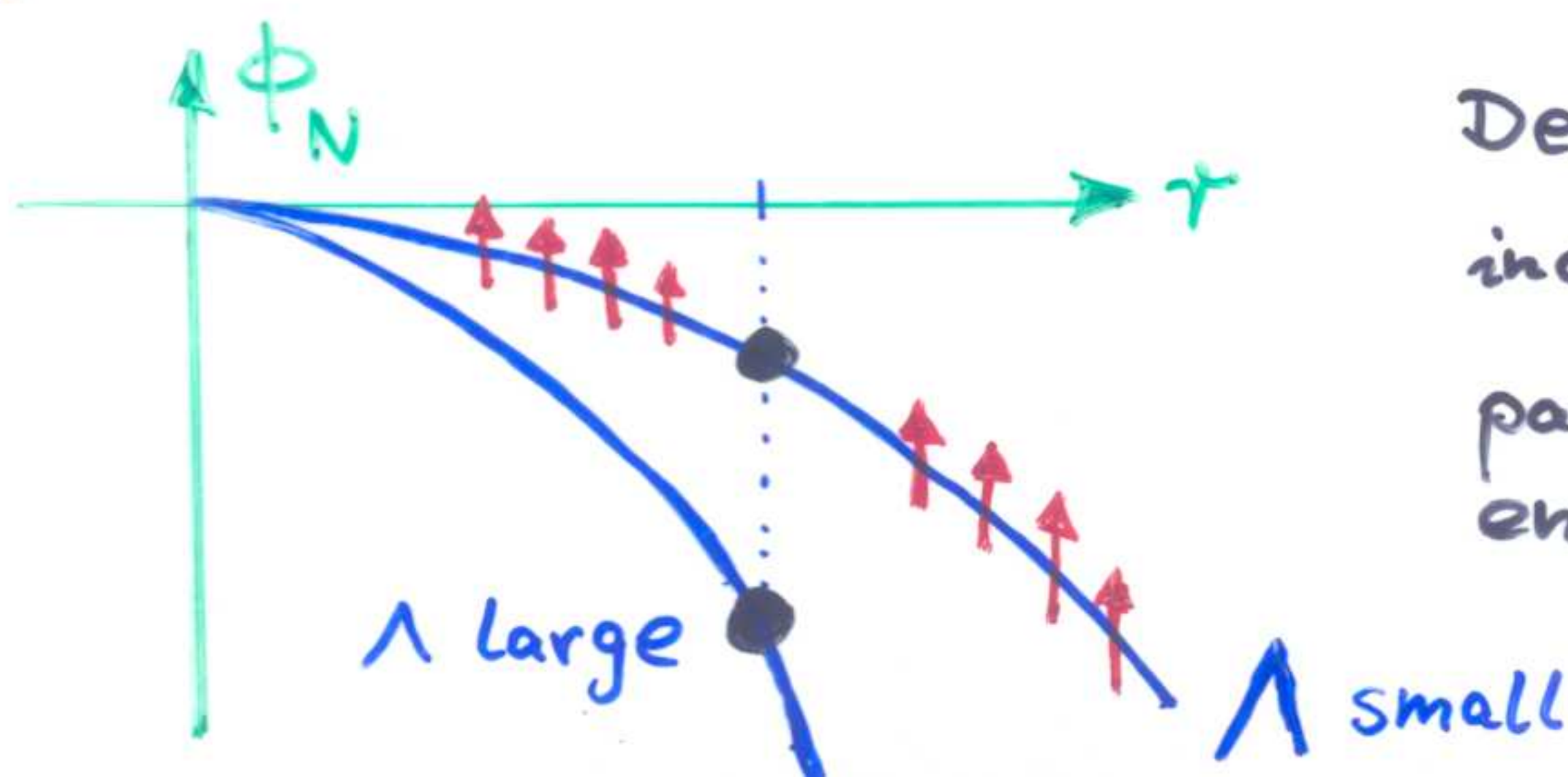
Example: de Sitter plus test particle

$$ds^2 = -[1 + 2\phi_N(r)]dt^2 + [1 + 2\phi_N(r)]^{-1}dr^2 + r^2 d\Omega^2$$

$$\phi_N(r) = -\frac{1}{6}\Lambda r^2$$

Newtonian interpretation:

$\Lambda > 0$:



Decreasing Λ
increases the test
particle's potential
energy!

RG - Improved Cosmology

Spatially flat RW geometry:

$$ds^2 = -dt^2 + a(t)^2 [dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2)]$$

$$T_{\mu}^{\nu} = \text{diag} [-\rho(t), p(t), p(t), p(t)]$$

The scale factor is determined by

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = -\Lambda(t) g_{\mu\nu} + 8\pi G(t) T_{\mu\nu}$$

where $\Lambda(t) \equiv \Lambda(k=k(t))$, $G(t) \equiv G(k=k(t))$

"Improved" Einstein's eq. $\Leftrightarrow$

modified Friedmann equation:

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3} G(t) \rho + \frac{1}{3} \Lambda(t)$$

modified conservation law $D^{\mu} [-\Lambda(t) g_{\mu\nu} + 8\pi G(t) T_{\mu\nu}] = 0$:

$$\underbrace{\dot{\rho} + 3H(\rho + p)}_{\sim D^{\mu} T_{\mu 0}} = - \frac{\dot{\Lambda} + 8\pi \rho \dot{G}}{8\pi G}$$

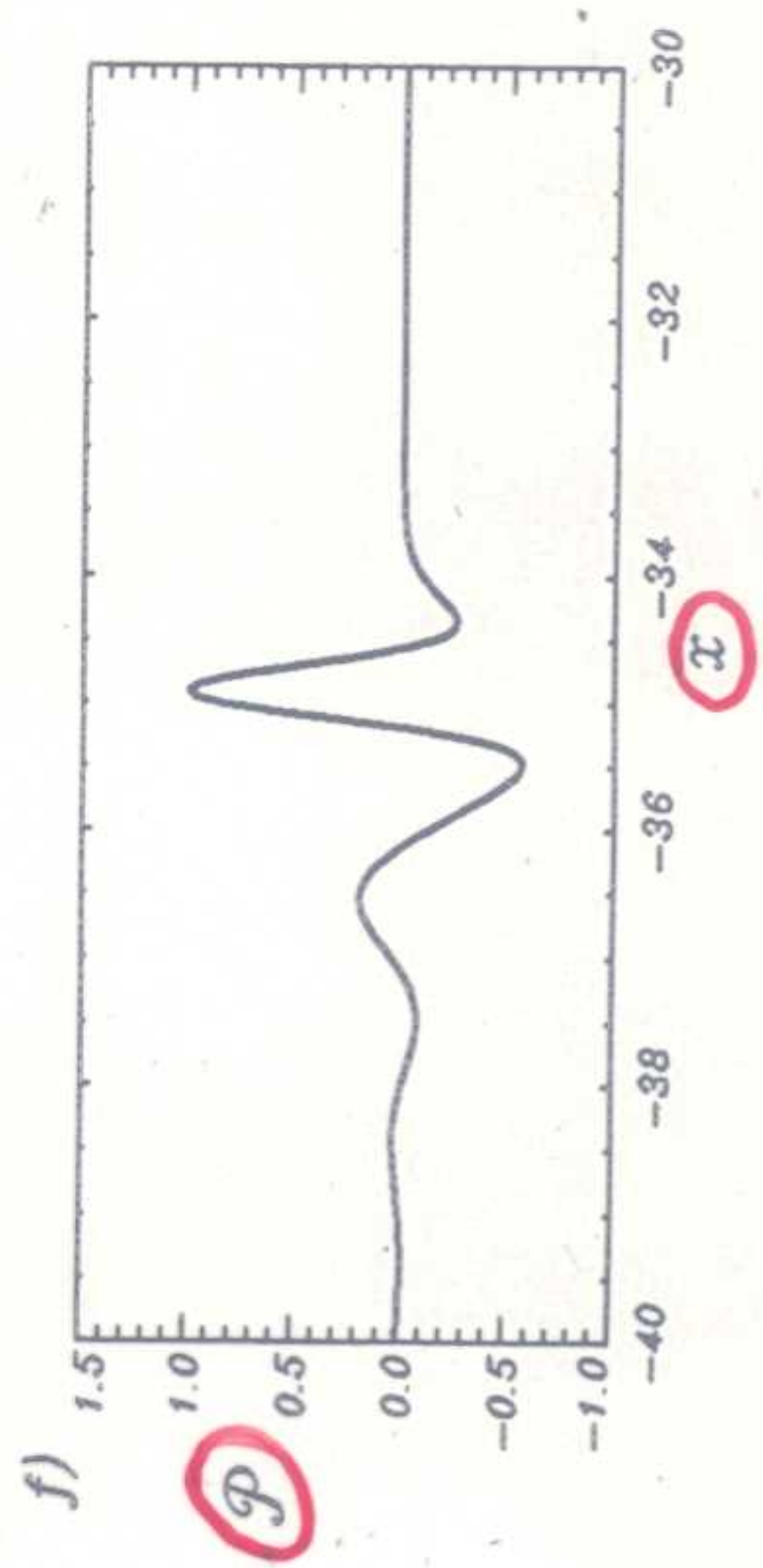
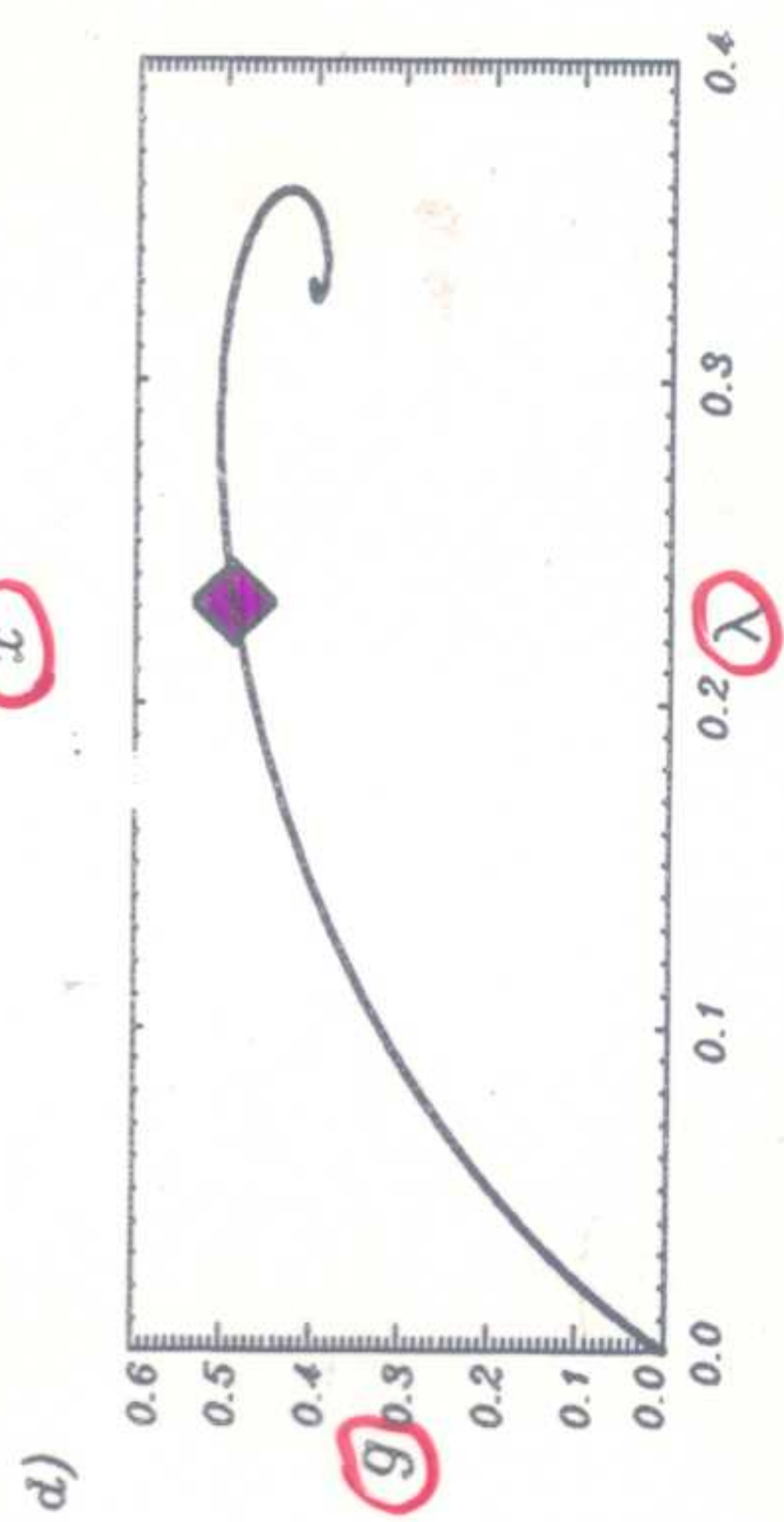
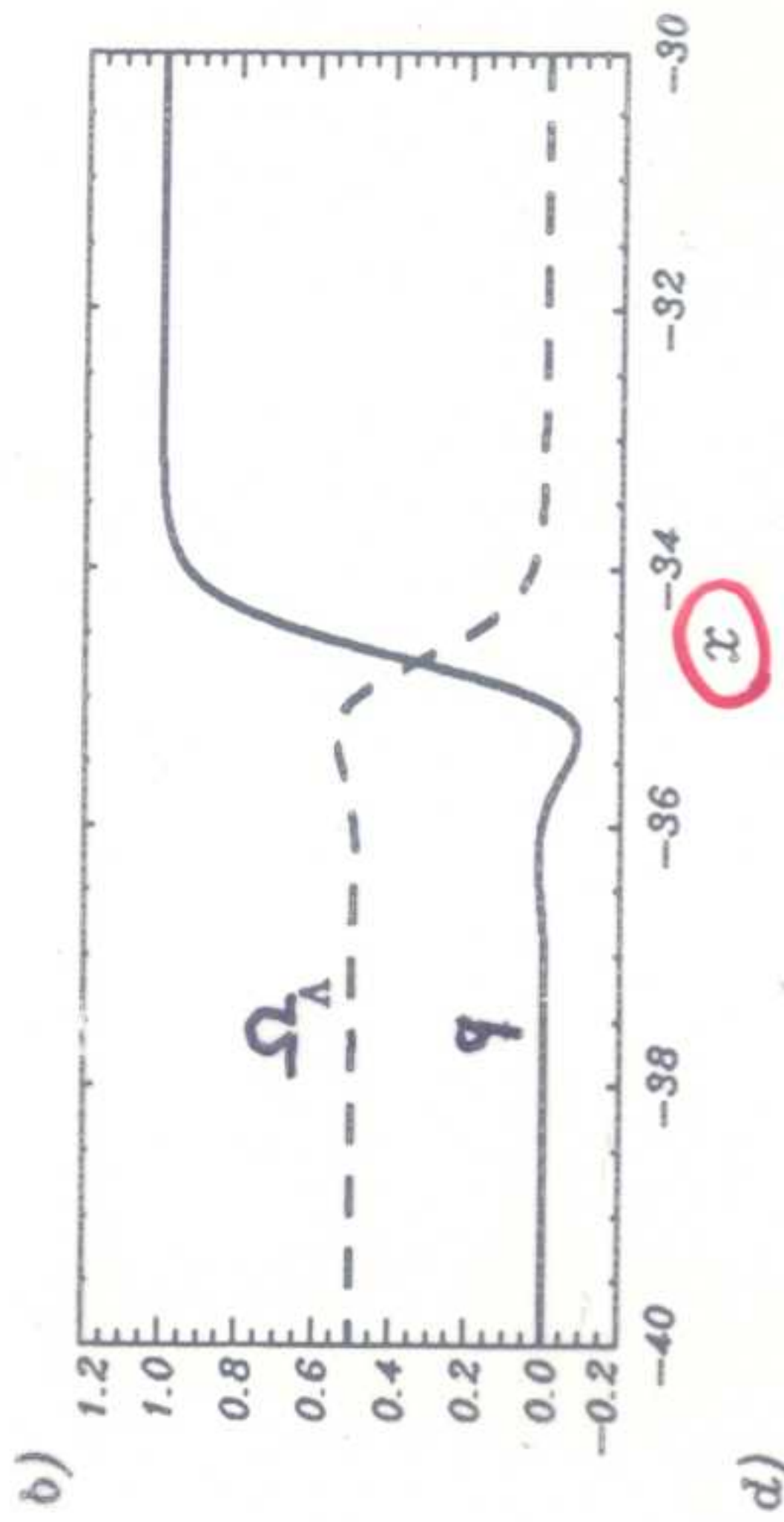
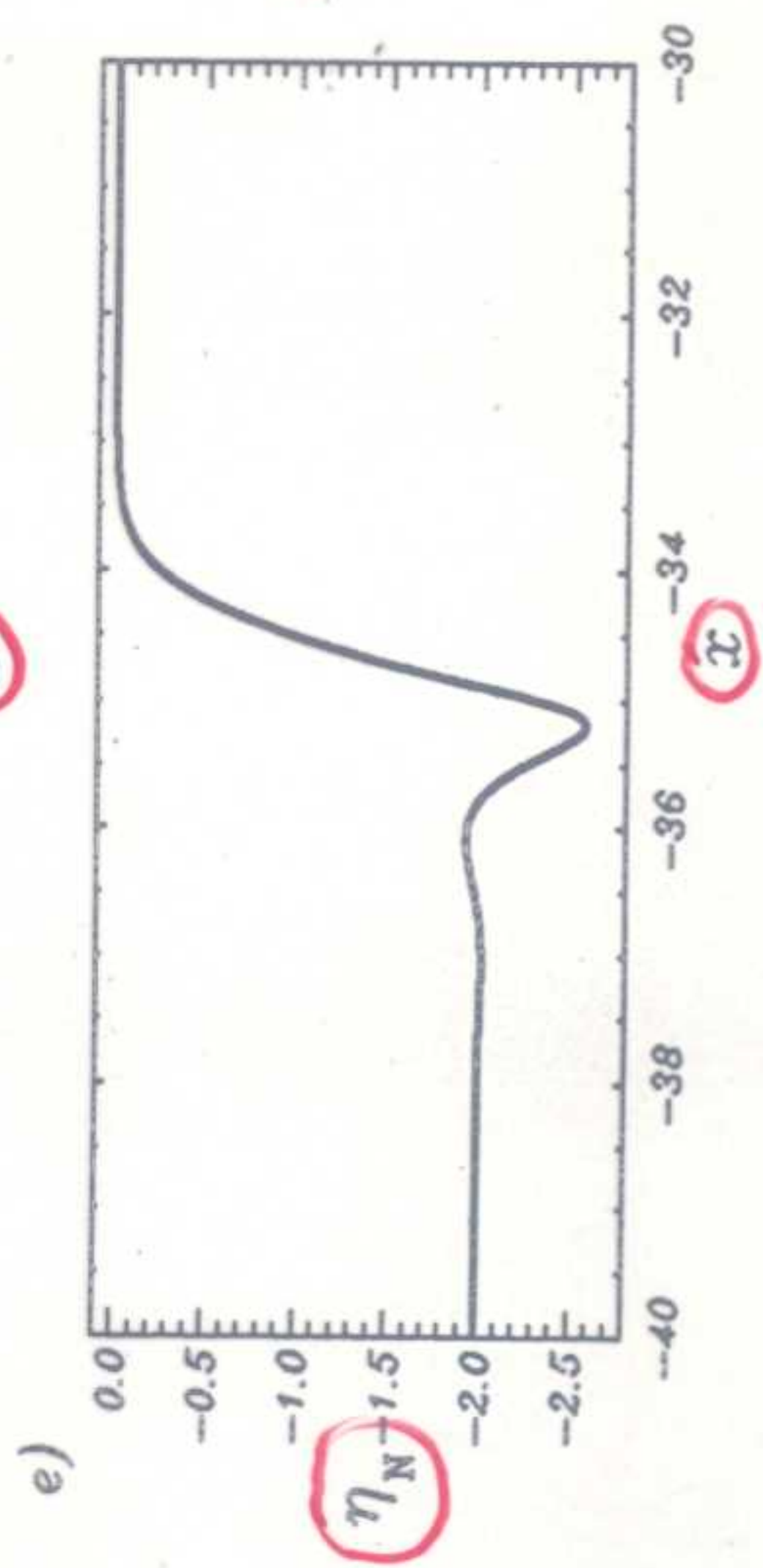
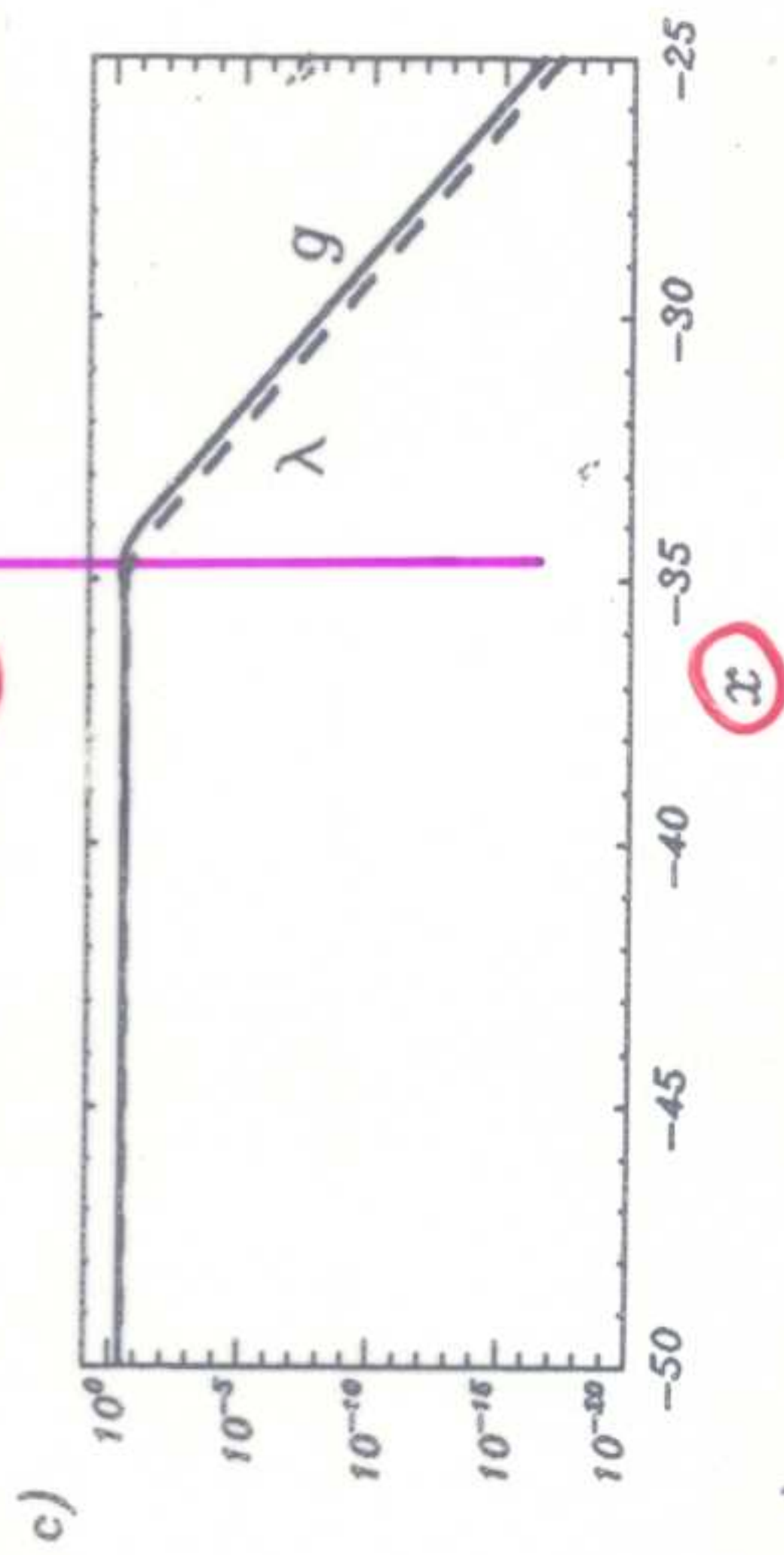
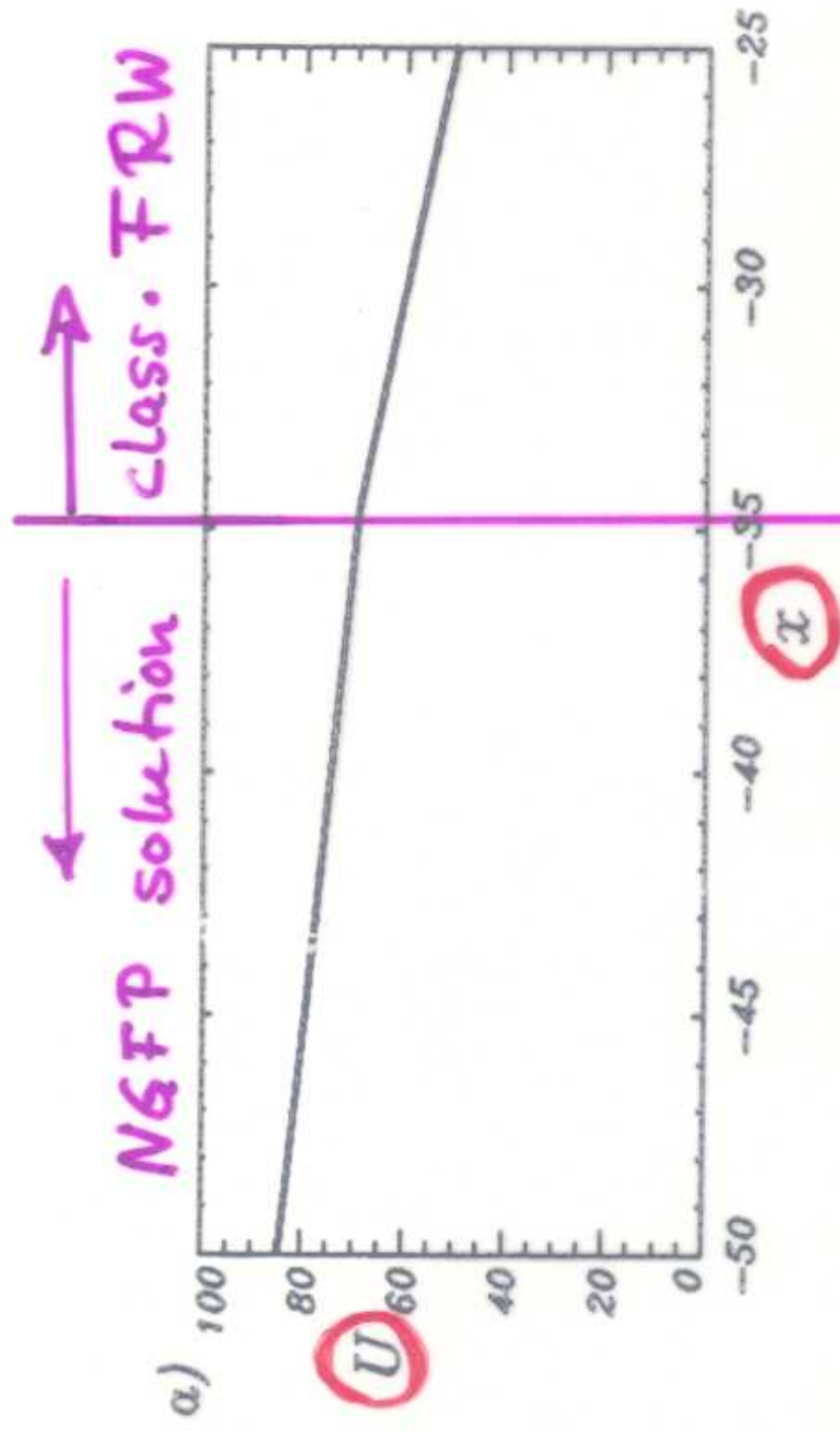
Describes energy exchange between matter and fields Λ , G ("vacuum").

Complete Cosmology:

$$U \equiv \ell_n \frac{H}{H_r}$$

$$x = \ell_n \frac{a}{a_r}$$

$$\Omega_\lambda^* = \frac{1}{2} \quad (\alpha = 1)$$

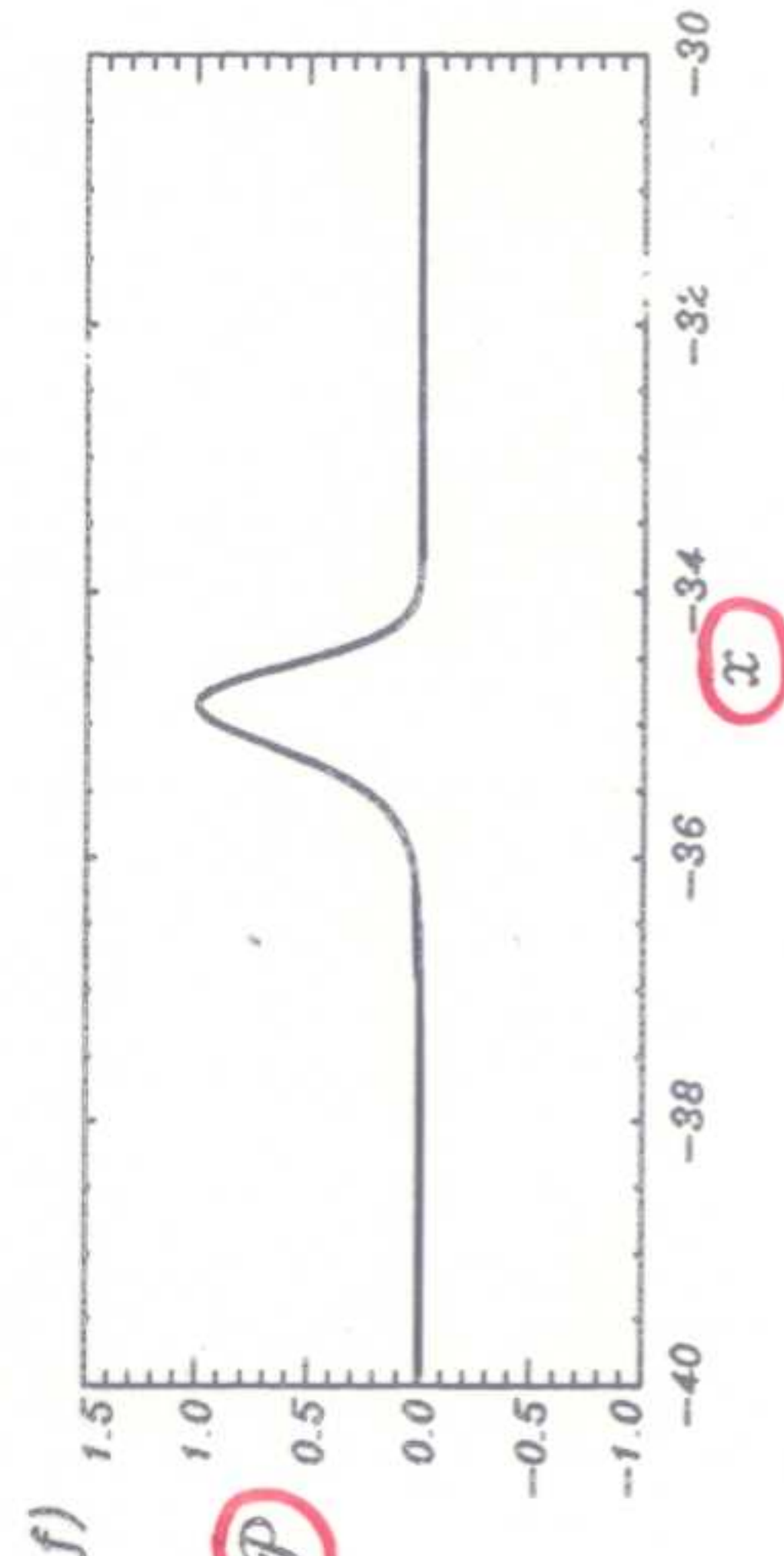
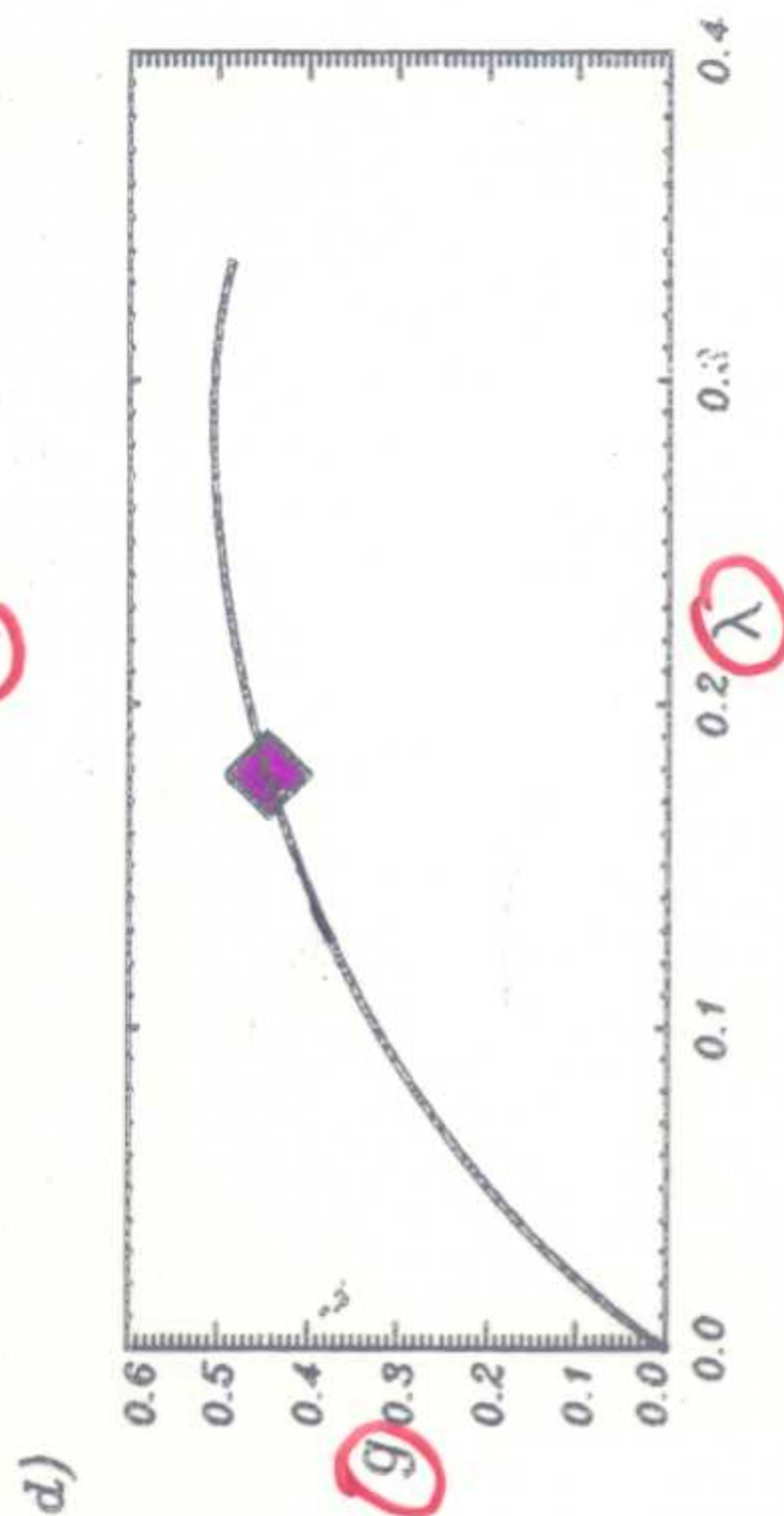
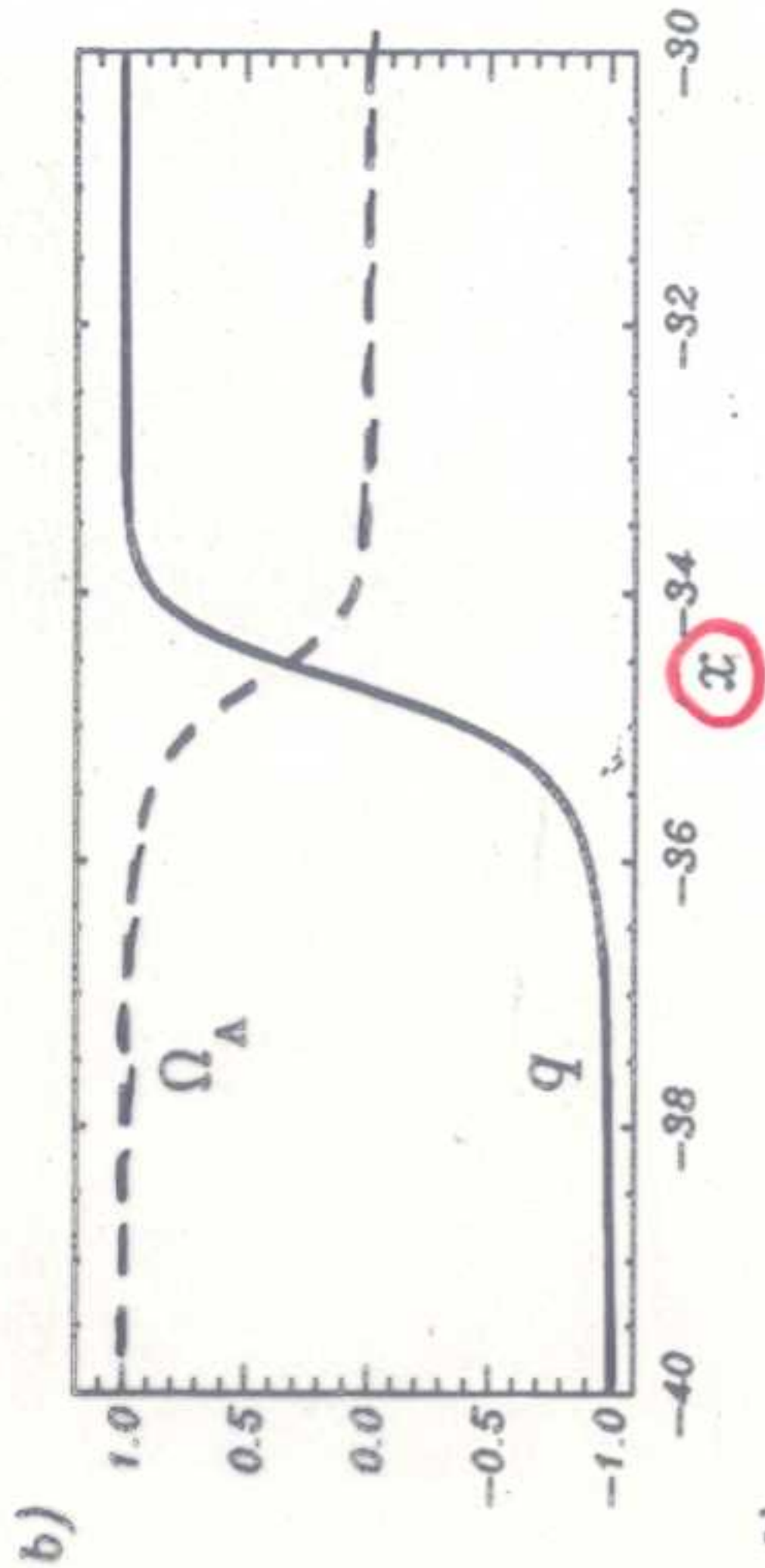
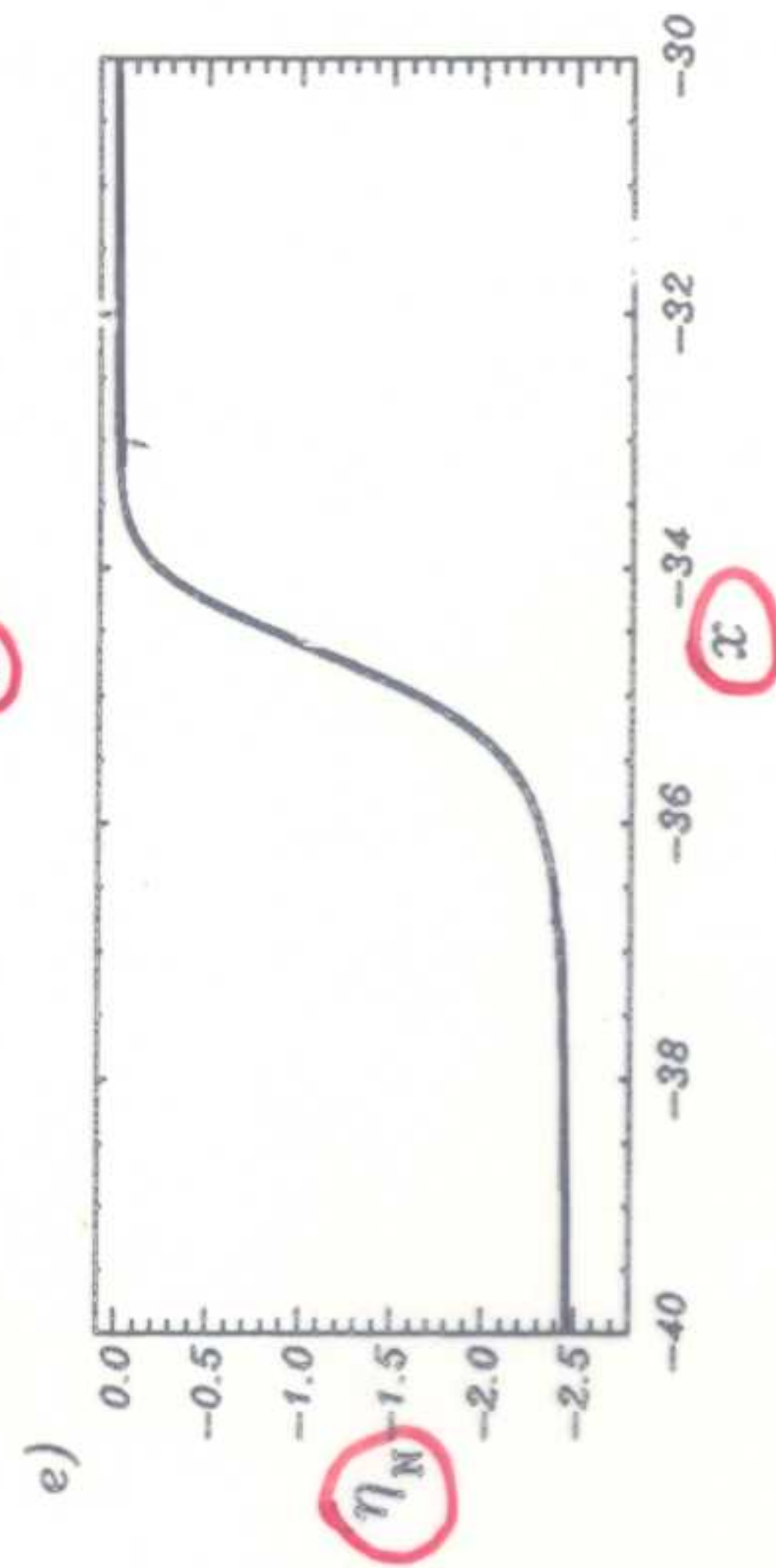
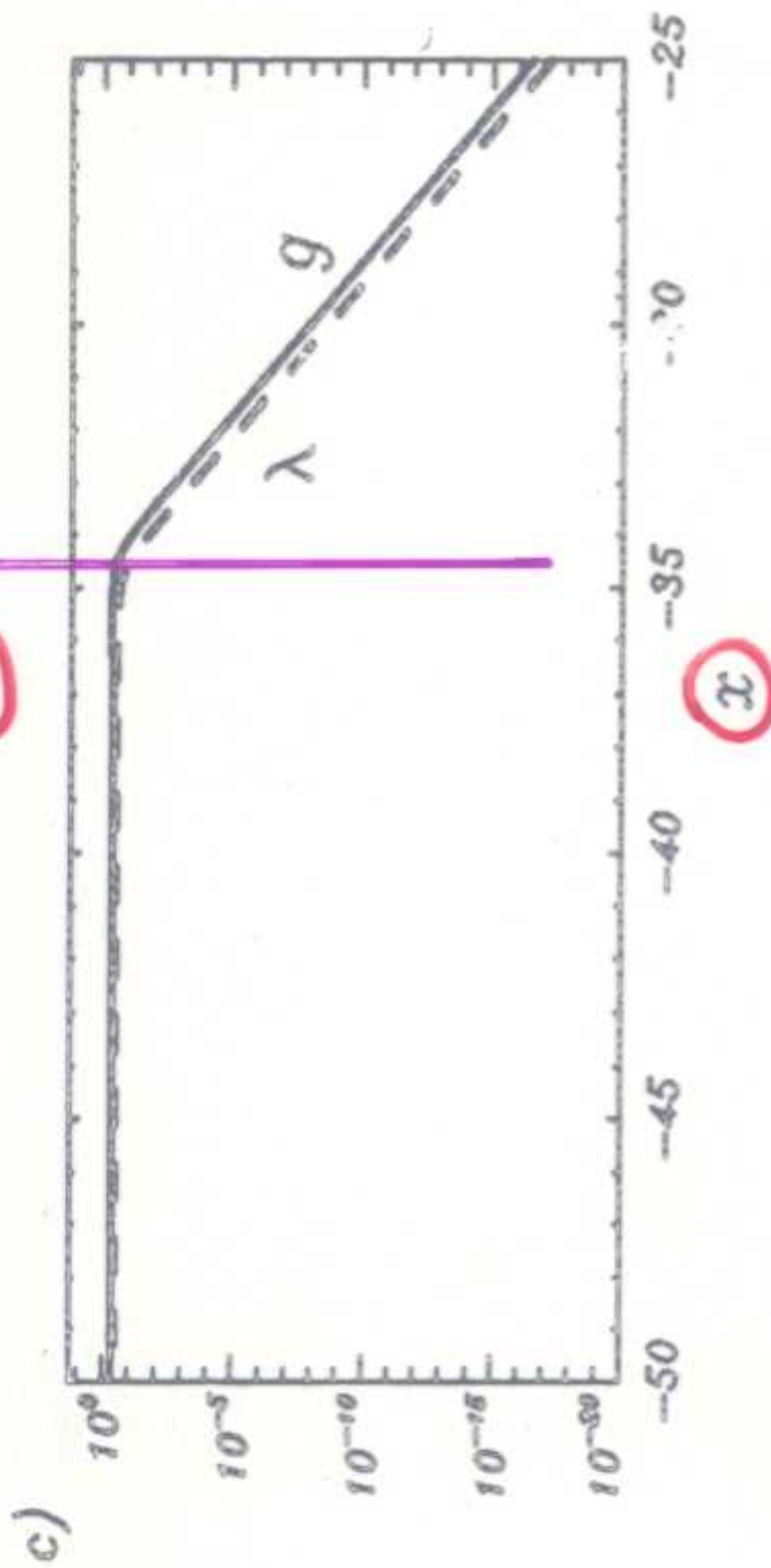
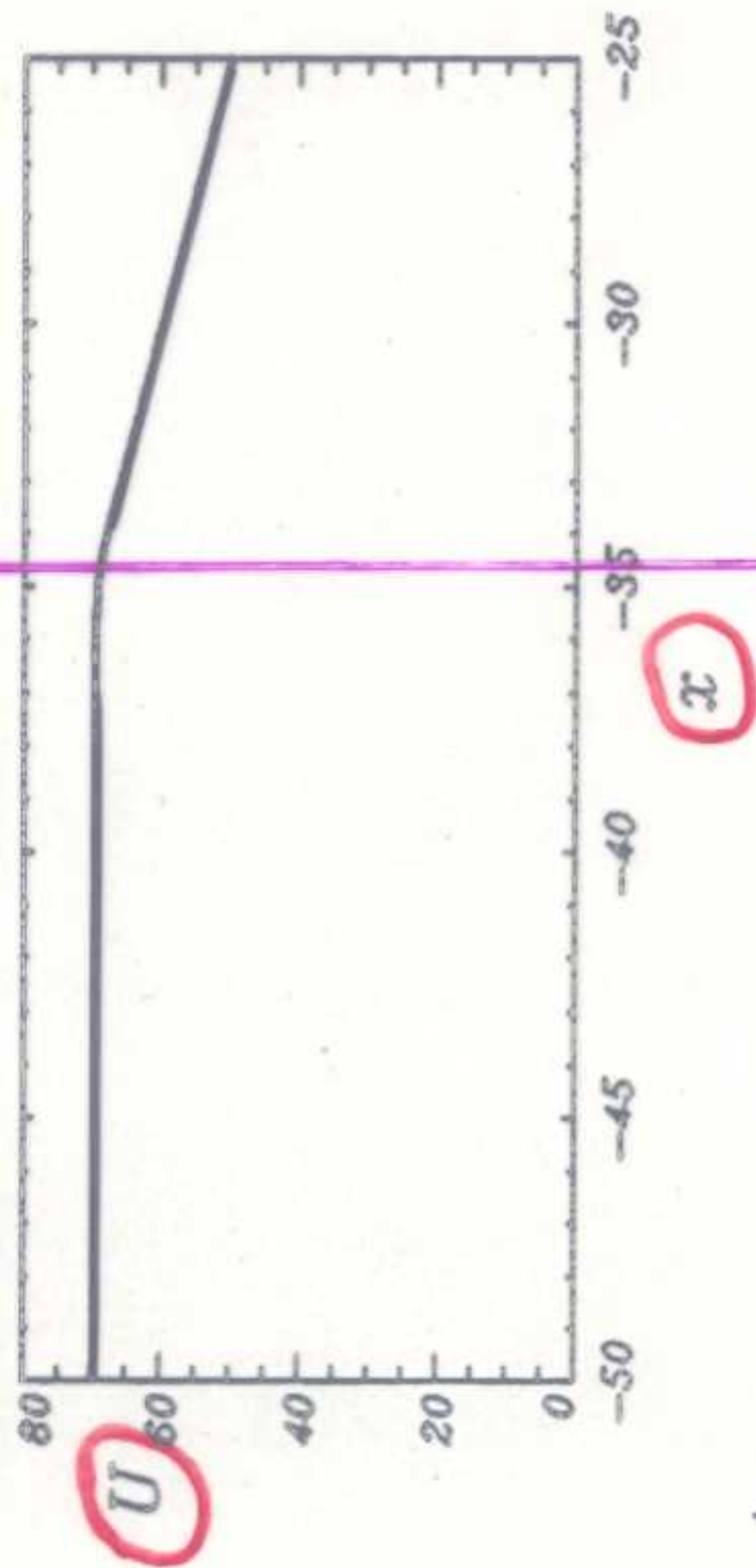


Complete Cosmology:

$$\Omega_{\Lambda}^* = 0.98$$

$$(\alpha = 25)$$

a) NGFP sol. $\approx dS$ $\leftarrow$ $\rightarrow$ class. FRW with $\Omega_{\Lambda}^{FRW} \approx 0$



$$U \equiv \ln \frac{H}{H_T}$$

$$x \equiv \ln \frac{a}{a_T}$$

Conclusion

Cosmology is a natural laboratory for confronting the predictions of asymptotically safe QEG with observations.

Candidates for observable/observed phenomena possibly explained by QEG :

- entropy of matter
- primordial density perturbations
 - NGFP regime is a "critical phenomenon"
 - Λ -driven inflation
- -
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