ATOMS OF SPACE AND THE COSMOS

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The Major Challenges to GR

SINGULARITIES: BLACK HOLES, UNIVERSE

COSMOLOGICAL CONSTANT

THE THERMODYNAMIC CONNECTION

These challenges involve \hbar !

$$A_{Planck}=rac{G\hbar}{c^3}; \ \ \Lambda\left(rac{G\hbar}{c^3}
ight)pprox 10^{-123}; \ \ k_BT=rac{\hbar}{c}\left(rac{g}{2\pi}
ight)$$

HOW DO WE PUT TOGETHER THE PRINCIPLES OF QUANTUM THEORY AND GRAVITY?

Everybody Wants To Quantize Gravity!

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Hathaway's Einstein theories

O fashion magazines for Hollywood actress Anne Hathaway, she spends her free time studying books on physics to enhance her knowledge on the universe.

The Devil Weers Prada star admits she shuns fashion magazines and instead stocks up on books by scientist Albert Einstein and physics textbooks in a bid to better understand the universe, reported a wolstie. "I'm interested in elementary particles. What I like thinking about is how time and space exist in the universe and how we understand it. Any spare time Liave, I bury my head in a physic sextook. The reading a lot about Einstein, I'like theories and I want to understand string theory," she said.

- IANS

... But Nobody Has Succeeded!

- The perturbative approach does not work
- Virtually every interesting question about gravity is non-perturbative by nature
- Non perturbative approches are too simplistic and not general enough (e.g., minisuperspace QC, CDT...)
- No guiding principle; spacetime metric is assumed to be a quantum variable

GR: The Next 100 Years

Needs another paradigm shift!

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Study spacetime dynamics the way physicists studied fluids before knowing the atomic structure of matter

The Importance Of Being Hot

You could have figured out that water is made of discrete atoms without ever probing it at Angstrom scales!

The Importance Of Being Hot

Boltzmann: If you can heat it, it must have micro-structure!

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You can even count the number of atoms $N=rac{E}{(1/2)k_BT}$

Microphysics leaves a signature at macro-scales

More IS Different

The key new variable which distinguishes thermodynamics from point mechanics

Heat $Density = \mathcal{H} = \frac{Q}{V} = \frac{TS}{V} = \frac{1}{V}(E - F)$ $\frac{TS}{V} = Ts = p + \rho$

Normal matter has a heat density

The Fluid called Spacetime

Spacetime also has a heat density!

One can associate a T and s with every event in spacetime, just as you could with a glass of water! Spacetime also has a heat density!

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Note: This fact transcends black hole physics and Einstein gravity

Macroscopic Nature Of Gravity

Field equations arise from maximizing entropy/heat density of gravity plus matter on all null surfaces.

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Works for a wide class gravitational theories; entropy decides the theory.

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 $Time\ evolution \propto (N_{
m sur} - N_{
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All static geometries have

 $N_{
m sur} = N_{
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A new perspective!

Gravity responds to heat density $(Ts = p + \rho)$ — not energy density!

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Cosmological constant arises as an integration constant Gravity responds to heat density $(Ts = p + \rho)$ — not energy density!

Cosmological constant arises as an integration constant

Its value is determined by a new conserved quantity for the universe!

The Nature Of Gravity

WHY?

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How do we understand this emergent nature of gravitational dynamics at a deeper level?

If gravity is immune to zero level of energy it **must** have a thermodynamic interpretation!

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Connects two features usually thought to be completely separate!

[TP, arXiv:1508.06286]

Since spacetime can be hot, it must have microstructure

We can count the atoms of spacetime without doing Planck scale experiments

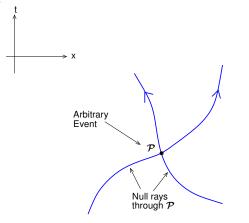
The Atoms Of Space

The distribution function for 'atoms of space' provides the microscopic origin for the variational principle

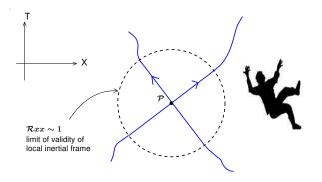
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Points in a renormalized spacetime has zero volume but finite area!

Spacetime in Arbitrary Coordinates



Local Inertial Observers



Validity of laws of SR \Rightarrow How gravity affects matter

Matter equations of motion $\Leftrightarrow \nabla_a T_b^a = 0$



Democracy of all observers

Regions of spacetime can be inaccessible to certain class of observers in any spacetime!



Democracy of all observers

Regions of spacetime can be inaccessible to certain class of observers in any spacetime!

Take non-inertial frames seriously: not "just coordinate relabeling".

The most beautiful result in the interface of quantum theory and gravity

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OBSERVERS WHO PERCEIVE A HORIZON ATTRIBUTE A TEMPERATURE TO SPACETIME

$$k_BT=rac{\hbar}{c}\left(rac{g}{2\pi}
ight)$$

[Davies (1975), Unruh (1976)]

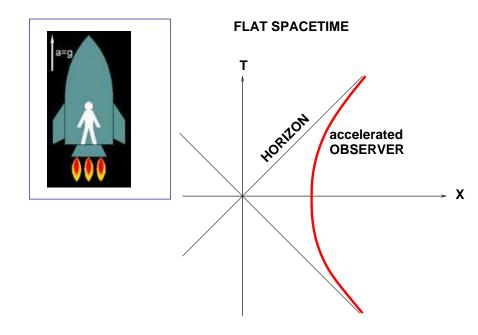
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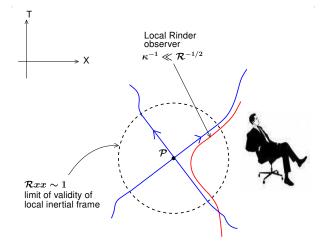
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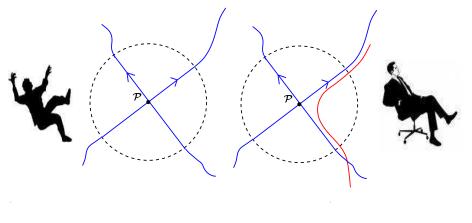
[Davies (1975), Unruh (1976)]

This allows you to associate a heat density $\mathcal{H} = Ts$ with every event of spacetime!



Local Rindler Observers

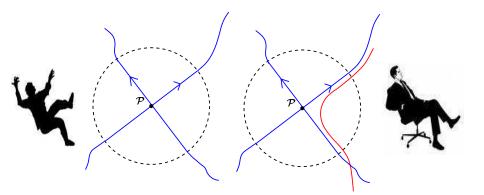




Vacuum fluctuations

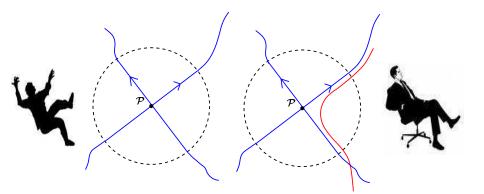
 \iff

Thermal fluctuations



Vacuum fluctuations \iff Thermal fluctuations

A VERY NON-TRIVIAL EQUIVALENCE!



Vacuum fluctuations 🛛 🔶 Thermal fluctuations

A VERY NON-TRIVIAL EQUIVALENCE!

QFT in FFF introduces \hbar ; we now have (\hbar/c) in the temperature

Local Rindler Horizon

Heat transfered due to matter crossing a null surface: [T. Jacobson, gr-qc/9504004]

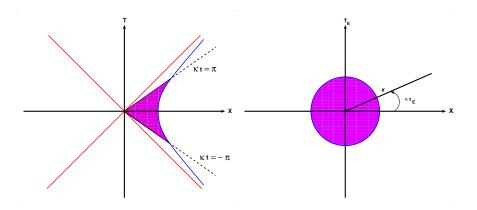
$$Q_m = \int d {\cal V} \, (T_{ab} \ell^a \ell^b); \ \ {\cal H}_m \equiv T_{ab} \ell^a \ell^b$$

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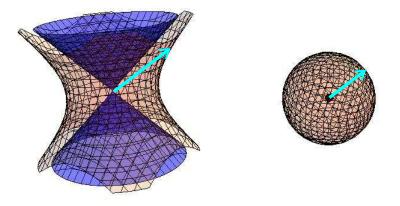
► Note: Null horizon ⇔ Euclidean origin

$$X^2-T^2=0 \Leftrightarrow X^2+T_E^2=0$$



 $T = x \sinh \kappa t, \ X = x \cosh \kappa t$ $T_E = x \sin \kappa t_E, \ X = x \cos \kappa t_E$

$$X^2-T^2=0 \Leftrightarrow X^2+T^2_E=0$$



 $X^2-T^2=\sigma^2 \Leftrightarrow X^2+T^2_E=\sigma^2$

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Guiding Principle For Dynamics

Matter equations of motion remain invariant when a constant is added to the Lagrangian

Gravity must respect this symmetry

The variational principle for the dynamics of spacetime must be invariant under

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The variational principle cannot have metric as the dynamical variable!



Variational principle must have the form

$$Q = \int dV ~ {\cal H}[T^a_b,g_{ab},q_A]$$

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How can \mathcal{H} depend on T^a_b but yet be invariant under $T^a_b \to T^a_b$ + (constant) δ^a_b ?

The Variational Principle

Minimal possibility: We must have

$$Q=\int dV \{ \mathcal{H}_g[g_{ab},n_a]+T^a_bn_an^b \}$$

where n_a is an auxiliary null vector field

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Can one find such a $\mathcal{H}_{g}[g_{ab}, n_{a}]$?

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

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Choose

$$\mathcal{H}_g = -\left(rac{1}{16\pi L_P^2}
ight) (4P^{ab}_{cd}
abla_a n^c
abla_b n^d)$$

$$P^{ab}_{cd} \propto \delta^{aba_2b_2...a_mb_m}_{cdc_2d_2...c_md_m} R^{c_2d_2}_{a_2b_2} \dots R^{c_md_m}_{a_mb_m}$$

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In
$$d=4$$
, it leads uniquely to GR $G^a_b=(8\pi L_P^2)T^a_b+\Lambda\delta^a_b$

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Macroscopically $\mathcal{H}_g(x^i,n_a)$ is the heat density of null surfaces with $n_a \propto \ell_a$

Microscopically $\mathcal{H}_g(x^i, n_a)$ is the 'distribution function for atoms of space with momentum' n_a

The Thermodynamic Connection

• Macroscopically, identify $n_a \leftrightarrow \ell_a$ and

$$Q_{
m tot} \equiv \int \sqrt{\gamma}\, d^2x d\lambda\, ({\cal H}_g[\ell] + {\cal H}_m[\ell])$$

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m loc}\,s) \Big|_{\lambda_1}^{\lambda_2}$$

Extremizing the heat densities of all null surfaces leads to gravitational dynamics!

Atoms of Matter

Boltzmann: If you can heat it, it must have micro-structure!

To store energy ΔE at temperature T, you need

$$\Delta n = rac{\Delta E}{(1/2)k_BT}$$

degrees of freedom. Connects microphysics with thermodynamics!

Boltzmann: If you can heat it, it must have micro-structure!

You can heat up spacetime!

Do we have an equipartition law for the microscopic spacetime degrees of freedom? Can you count the atoms of space?

The Quantum of Area

TP, [gr-qc/0308070]; [0912.3165]; [1003.5665]

Equipartition with a surface-bulk correspondence

$$E_{
m bulk} = \int_{\partial \mathcal{V}} rac{dA}{L_P^2} \left(rac{1}{2} k_B T_{loc}
ight) \equiv rac{1}{2} k_B \int_{\partial \mathcal{V}} dn \, T_{
m loc}$$

Associates $dn = dA/L_P^2$ atoms (microscopic degrees of freedom) with an area dA

We must be able to express — and interpret — the field equation in a purely thermodynamic language !

Geometry \Leftrightarrow **Thermodynamics**

K. Parattu, B.R. Majhi, T.P. [arXiv:1303.1535]

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$$q^{ab}\equiv\sqrt{-g}\,g^{ab}$$
 $p^a_{bc}\equiv-\Gamma^a_{bc}+rac{1}{2}(\Gamma^d_{bd}\delta^a_c+\Gamma^d_{cd}\delta^a_b)$

These variables have a thermodynamic interpretation

 $(q\delta p,p\delta q) \Leftrightarrow (s\delta T,T\delta s)$

T.P., Gen.Rel.Grav (2014) [arXiv:1312.3253]

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$$\int rac{d\Sigma_a}{8\pi L_P^2} \left[q^{\ell m} \partial \; p^a_{\ell m}
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time evolution of spacetime

= heating of spacetime

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time evolution of spacetime = heating of spacetime deviation from holographic equipartition

This replaces the field equation for gravity

T.P. [hep-th/0205278]

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Three constants: \hbar, c, L_P^2

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Gravity, like matter, is intrinsically quantum and cannot exist in the limit of $\hbar \rightarrow 0$!

What Next?

T.P. [arXiv:1508.06286]



Understand this from a deeper, microscopic, level:



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Origin of the auxiliary vector field n_a



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- Origin of the auxiliary vector field n_a
- Why are null vectors selected out?
- Determine *H_g*; use alternative, dimensionless, form:

$${\cal H}_g\equiv -{1\over 8\pi}(L_P^2R_{ab}n^an^b)$$

The Challenge

How can we get \mathcal{H}_g from a microscopic theory without knowing the full QG?

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We need to recognize discreteness and yet use continuum mathematics!

Atoms Of A Fluid

• Continuum fluid mechanics: $\rho(x^i)$, $U(x^i)$, ignores discreteness and *velocity dispersion*.

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 - Many atoms with different p_i can exist at same x^i

Atoms of Space

Atoms of Space

The \mathcal{H}_g is proportional to the $f(x^i, n_j)$ for the number of atoms of space "at" x^i with "momentum" n_i . In dimensionless form:

 $rac{d(Q/E_P)}{d(V/L_P^3)}\equiv \mathcal{H}_g(x^i,n_j)=f(x^i,n_j)$

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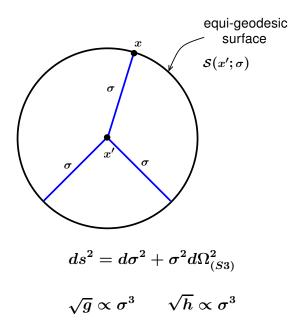
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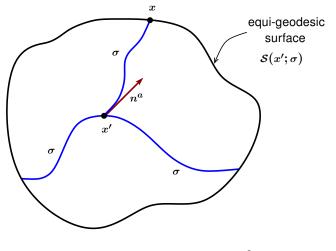
We expect $f(x^i, n_i)$ to be proportional to the volume or the area measure "associated with" the event x^i in the spacetime The \mathcal{H}_g is proportional to the $f(x^i, n_j)$ for the number of atoms of space "at" x^i with "momentum" n_i . In dimensionless form:

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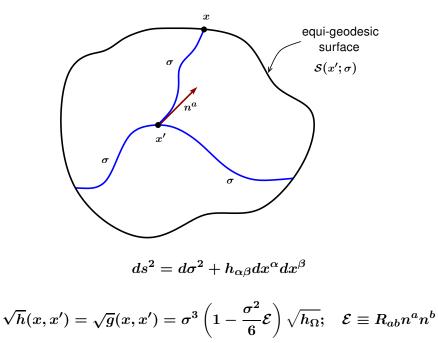
Use equi-geodesic surfaces to make this idea precise





 $ds^2 = d\sigma^2 + h_{lphaeta} dx^lpha dx^eta$

The $\sqrt{g} = \sqrt{h}$ will pick up curvature corrections



Zero-Point Length

T.P. Ann.Phy. (1985), 165, 38; PRL (1997), 78, 1854

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Quantum spacetime has a zero-point length:

$$egin{array}{rcl} \sigma^2(x,x') & o & S(\sigma^2) = \sigma^2(x,x') + L_0^2 \ g_{ab}(x) & o & q_{ab}(x,x';L_0^2) \end{array}$$

D. Kothawala, T.P. [arXiv:1405.4967]; [arXiv:1408.3963]

Origin of Null Vectors

The number of atoms of space at x^i with attribute ("momentum") n_i scales as volume or area measure of the equigeodesic surface in the quantum Euclidean space when $x' \rightarrow x$

 $f(x^i,n_j) \propto \sqrt{g}(x^i,n_j) ~ {
m OR} ~ \sqrt{h}(x^i,n_j)$

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m OR} \; \sqrt{h}(x^i,n_j)$$

The $\sigma^2 \rightarrow 0$ limit picks null vectors! Euclidean origin maps to local Rindler horizons.

T.P. [arXiv:1508.06286]

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$$\sqrt{q}=\sigma\left(\sigma^{2}+L_{0}^{2}
ight)\left[1-rac{1}{6}\mathcal{\mathcal{E}}\left(\sigma^{2}+L_{0}^{2}
ight)
ight]\sqrt{h_{\Omega}}$$

T.P. [arXiv:1508.06286]

$$egin{aligned} \sqrt{q} &= \sigma \left(\sigma^2 + L_0^2
ight) \left[1 - rac{1}{6} \mathcal{E} \left(\sigma^2 + L_0^2
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T.P. [arXiv:1508.06286]

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Aside: Spacetime becomes two-dimensional at Planck scales

The area measure gives exactly what we need, along with a zero-point-contribution

$$f(x^i,n_a) = 1 - rac{1}{8\pi} \mathcal{E}L_P^2 = 1 - rac{1}{8\pi} L_P^2 R_{ab} n^a n^b$$

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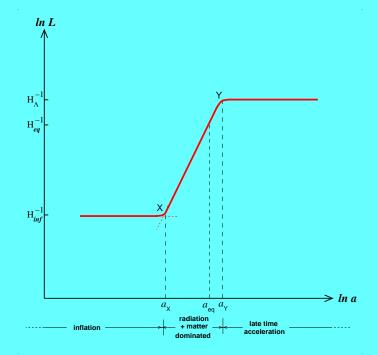
Zero-point contribution is important; degrees of freedom of Planck 2-sphere: $4\pi L_P^2/L_P^2 = 4\pi$

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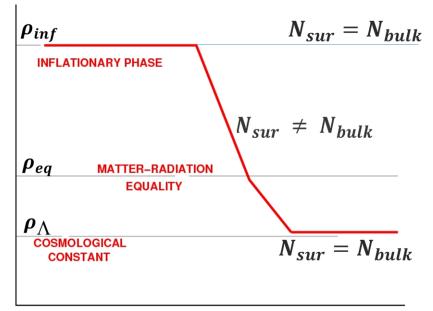
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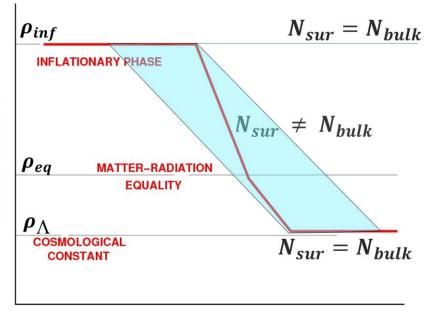
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But we have no clue why $ho_\Lambda L_P^4 pprox 1.4 imes 10^{-123} pprox 1.1 imes e^{-283}.$



SIZE OF THE UNIVERSE



SIZE OF THE UNIVERSE

Value of the Cosmological Constant

Hamsa Padmanabhan, T.P. [arXiv:1302.3226]

 $= \frac{4}{27} \frac{\rho_{inf}^{3/2}}{\rho_{ca}^{1/2}} \exp(-36\pi^2)$



$P(x^i,n_a) \propto \exp[\mu f(x^i,n_a)]$



$P(x^i,n_a) \propto \exp[\mu f(x^i,n_a)] onumber \ \langle n^a n^b angle pprox (4\pi/\mu L_P^2) R_{ab}^{-1}$

Speculation

$$P(x^i,n_a) \propto \exp[\mu f(x^i,n_a)]$$
 $\langle n^a n^b
angle pprox (4\pi/\mu L_P^2) R_{ab}^{-1}$

$2\mu L_P^4 \left< ar{T}_{ab} n^a n^b \right> pprox 2\mu L_P^4 \left< ar{T}_{ab} \right> \left< n^a n^b \right> = 1$

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MAY BE ONE SHOULD NOT THINK OF COSMOLOGY AS PART OF GENERAL RELATIVITY!

Open Questions

Matter and Geometry need to emerge together for proper interpretation of T^{ab}n_an_b at the microscopic scale. How do we do this?

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- Generalisation to Lanczos-Lovelock models with $R_{ab} \rightarrow \mathcal{R}_{ab}$: What happens at microscopic scales?

Summary

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- If we take number of atoms of spacetime at a point to be proportional to the area measure in a spacetime with zero-point length, we get the correct variational principle
- A Planck scale 2-sphere has 4π degrees of freedom which allows the determination of the cosmological constant

References

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Acknowledgements

Sunu Engineer Dawood Kothawala Bibhas Majhi Krishna Parattu Sumanta Chakraborty James Bjorken Aseem Paranjape Hamsa Padmanabhan Donald Lynden-Bell

THANK YOU FOR YOUR TIME!