ENTROPY OF SPACETIME AND GRAVITY

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PLAN OF THE TALK

- GRAVITY AND THERMODYNAMICS: HINTS OF A DEEPER CONNECTION
- THE SAKHAROV PARADIGM
- GRAVITY AS AN EMERGENT PHENOMENON
- ENTROPY OF NULL SURFACES AND SPACETIME DYNAMICS
- HORIZON ENTROPY
- PERSPECTIVE AND CONCLUSIONS

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Question:

 Is there a deeper connection between these results and spacetime dynamics ["Matter tells spacetime how to curve"]? What is this telling us?

• Metric:

$$ds^{2} = f(r)dt^{2} - f(r)^{-1}dr^{2} - r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2})$$

$$\equiv f(r)dt^{2} - f(r)^{-1}dr^{2} - dL_{\perp}^{2}$$

• Source:

$$T_t^t = T_r^r = \frac{\epsilon(r)}{8\pi}; \quad T_\theta^\theta = T_\phi^\phi = \frac{\mu(r)}{8\pi}$$

• Einstein's equations:

$$\frac{1}{r^2}(1-f) - \frac{f'}{r} = \epsilon; \quad \nabla^2 f = -2\mu$$

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- Consider the case with horizon at r = a; that is, f = 0 at r = a with f'(a) = B > 0.
- Temperature of horizon: $k_B T = \hbar c B / 4\pi$.

• At $r = a, f(a) = 0, f'(a) \equiv B$. Einstein's equation gives:

$$\frac{c^4}{G}\left[\frac{1}{2}Ba - \frac{1}{2}\right] = -4\pi T_r^r a^2$$

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- Works for time-dependent horizons, Kerr, Lanczos-Lovelock theory; close to Membrane Paradigm in spirit

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Mechanics; Elasticity $(\rho, \mathbf{v} \dots)$

Einstein's Theory $(g_{ab} \dots)$

Statistical Mechanics

of atoms/molecules

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- Strategy: Associate an entropy with a null surfaces. Demand maximisation of entropy of all null surfaces to get the dynamics.
- Note: There is no 'quantum thermodynamics'; for example, TdS = dE + PdV is valid with quantum corrections as well.

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- Our state of knowledge of quantum gravity is worse that that of 18th century physicists about microstructure of solids! So we should not hesitate to be phenomenological.

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- Associate with every vector field an entropy which is quadratic in deformation field:

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• Demand that: (i) P^{abcd} has the symmetries of Riemann tensor; (ii) Covariant divergence on all indices are zero: $\nabla_a P^{abcd} = 0$; $\nabla_a T^{ab} = 0$. [Analogue of elastic "constants"]

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- We expect *P*^{*abcd*} to have a (RG-like) derivative expansion in powers of number of derivatives of the metric:

$$P^{abcd}(g_{ij}, R_{ijkl}) = c_1 \overset{(1)}{P}^{abcd}(g_{ij}) + c_2 \overset{(2)}{P}^{abcd}(g_{ij}, R_{ijkl}) + \cdots,$$

• If we do not use the curvature tensor:

$$\stackrel{(1)}{P}{}^{ab}_{cd} = rac{1}{32\pi} (\delta^a_c \delta^b_d - \delta^a_d \delta^b_c) = rac{1}{16\pi} rac{1}{2} \delta^{ab}_{cd} \, .$$

THE DERIVATIVE EXPANSION FOR P^{ab}_{cd}

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• Linear on curvature:

$$\overset{(2)}{P}{}^{ab}_{cd} = \frac{1}{8\pi} \left(R^{ab}_{cd} - G^a_c \delta^b_d + G^b_c \delta^a_d + R^a_d \delta^b_c - R^b_d \delta^a_c \right) = \frac{1}{16\pi} \frac{1}{2} \delta^{aba_3a_4}_{cd \, b_3 \, b_4} R^{b_3b_4}_{a_3a_4} \,.$$

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• The *m*-th order term is unique:

$${}^{(m)}_{P \ cd} \propto \delta^{aba_3...a_{2m}}_{cdb_3...b_{2m}} R^{b_3b_4}_{a_3a_3} \cdots R^{b_{2m-1}b_{2m}}_{a_{2m-1}a_{2m}}$$

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• Explicit form:

$$S_1[\xi] = \int_{\mathcal{V}} \frac{d^D x}{8\pi} \left((\nabla_c \xi^c)^2 - \nabla_a \xi^b \nabla_b \xi^a \right)$$
$$S_2[\xi] = c_2 \int_{\mathcal{V}} d^D x \left(R^{cd}_{ab} \nabla_c \xi^a \nabla_d \xi^b - (G^c_a + R^c_a) (\nabla_c \xi^a \nabla_b \xi^b - \nabla_c \xi^b \nabla_b \xi^a) \right)$$

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• Closely related to the Lanczos-Lovelock Lagrangian:

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• The Lanczos-Lovelock Lagrangian separates to a bulk and surface terms

$$\sqrt{-g}L = 2\partial_c \left[\sqrt{-g}Q_a^{\ bcd}\Gamma^a_{bd}\right] + 2\sqrt{-g}Q_a^{\ bcd}\Gamma^a_{dk}\Gamma^k_{bc} \equiv L_{\rm sur} + L_{\rm bulk}$$

and is 'holographic':

$$[(D/2) - m]L_{sur} = -\partial_i \left[g_{ab} \frac{\delta L_{bulk}}{\delta(\partial_i g_{ab})} + \partial_j g_{ab} \frac{\partial L_{bulk}}{\partial(\partial_i \partial_j g_{ab})} \right]$$

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• The surface term is closely related to horizon entropy in Lanczos-Lovelock theory.

• Demand that $\delta S = 0$ for variations of all null vectors:

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• This leads to Lanczos-Lovelock theory with an arbitrary cosmological constant:

$$16\pi \left[P_b^{\ ijk} R^a_{\ ijk} - \frac{1}{2} \delta^a_b \mathcal{L}_m \right] = 8\pi T^a_b + \Lambda \delta^a_b,$$

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• In a derivative coupling expansion, Lanczos-Lovelock terms are corrections.

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- At the lowest order, we get quarter of transverse area as entropy.
- For *any* solution, in a local Rindler frame, the causal horizons have the correct entropy.

THE REAL TROUBLE COSMOLOGICAL CONSTANT

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- So after you have "solved" the cosmological constant problem, if someone introduces $L_{matter} \rightarrow L_{matter} \rho$, you are in trouble again!
- The only way out is to have a formalism for gravity which is invariant under $T_{ab} \rightarrow T_{ab} + \rho g_{ab}$.

- The matter sector and its equations are invariant under the shift of the Lagrangian by a constant: $L_{matter} \rightarrow L_{matter} \rho$.
- But this changes energy momentum tensor by $T_{ab} \rightarrow T_{ab} + \rho g_{ab}$ and gravity sector is not invariant under this transformation.
- So after you have "solved" the cosmological constant problem, if someone introduces $L_{matter} \rightarrow L_{matter} \rho$, you are in trouble again!
- The only way out is to have a formalism for gravity which is invariant under $T_{ab} \rightarrow T_{ab} + \rho g_{ab}$.
- This will also make gravity immune to bulk vacuum energy density.

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- But not the quantum fluctuations in the energy density.

VACUUM FLUCTUATIONS AND $\rho_{\rm vac}$

T.P, CQG, 22, L107-L110, (2005) hep-th/0406060

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Important: Bulk term
is now ignored by gravity

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- Connects with the radial displacements of horizons and TdS = dE + PdV as the key to obtaining a thermodynamic interpretation of gravitational theories.
- The deep connection between gravity and thermodynamics *goes well beyond Einstein's theory*. Closely related to the holographic structure of Lanczos-Lovelock theories.



OPEN QUESTIONS

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 - Conditions on P_{cd}^{ab} , derivatives of curvature tensor
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 Analogy: horizons in euclidean spacetime ↔ defects in solids. Clue for a new path integral prescription ?

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