# **GRAVITY: THE INSIDE STORY**

T. Padmanabhan

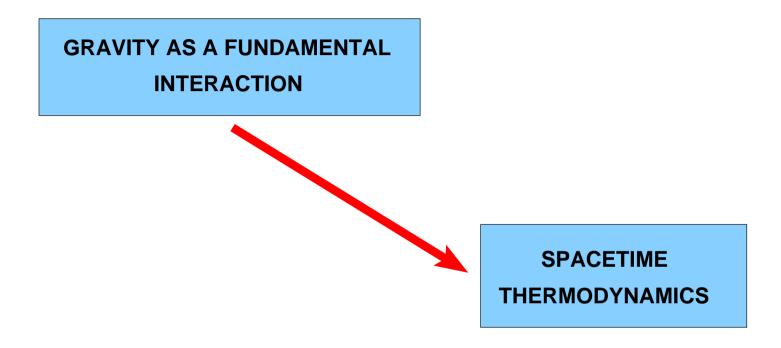
(IUCAA, Pune, India)

VR Lecture, IAGRG Meeting Kolkatta, 28 Jan 09

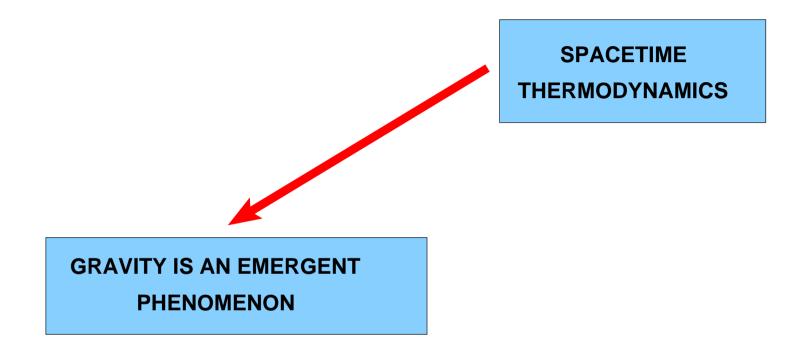
# CONVENTIONAL VIEW

# GRAVITY AS A FUNDAMENTAL INTERACTION

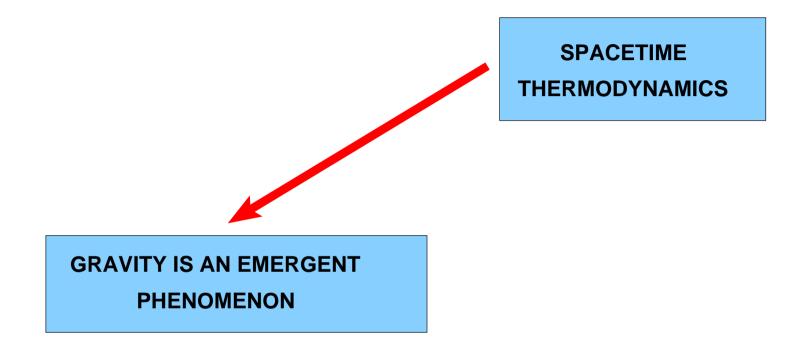
# CONVENTIONAL VIEW



# **NEW PERSPECTIVE**



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GRAVITY IS THE THERMODYNAMIC LIMIT OF THE STATISTICAL MECHANICS OF 'ATOMS OF SPACETIME'

**SOLIDS** 

**SPACETIME** 

Mechanics; Elasticity  $(\rho, \mathbf{v} \dots)$ 

Einstein's Theory  $(g_{ab} ...)$ 

Statistical Mechanics

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of atoms/molecules

of "atoms of spacetime"

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- Thermodynamics offers a connection between the two though the form of entropy functional,  $S[\xi]$ . No microstructure, no thermodynamics!
- You never took a course in 'quantum thermodynamics'.

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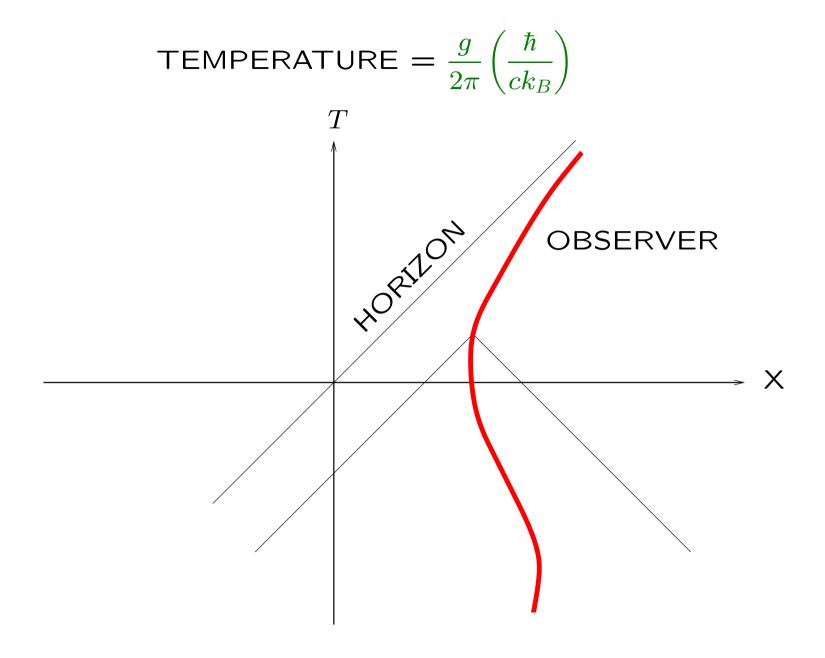
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- Rindler horizons have a temperature (1975-76)

TEMPERATURE = 
$$\frac{g}{2\pi} \left( \frac{\hbar}{ck_B} \right)$$

OBSERVER



Works for Blackholes, deSitter, Rindler

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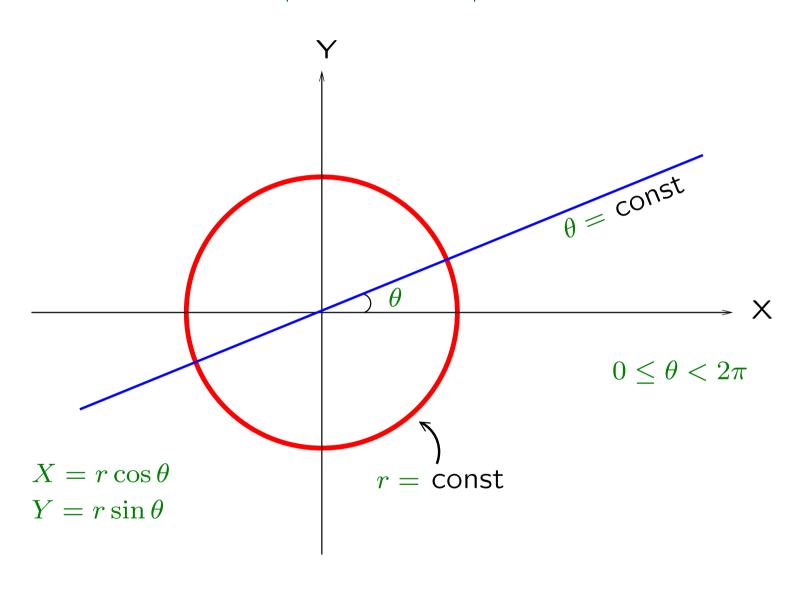
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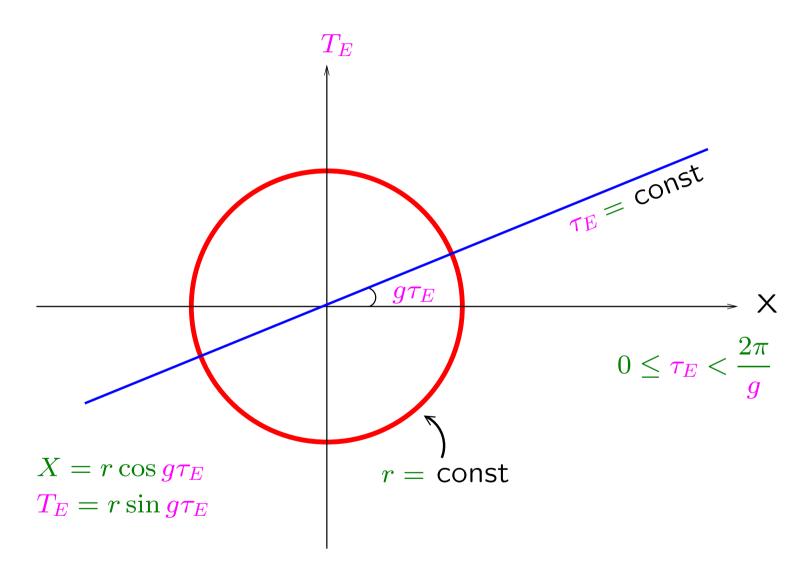
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SPACETIMES WITH HORIZONS EXHIBIT PERIODICITY IN IMAGINARY TIME  $\Longrightarrow$  TEMPERATURE

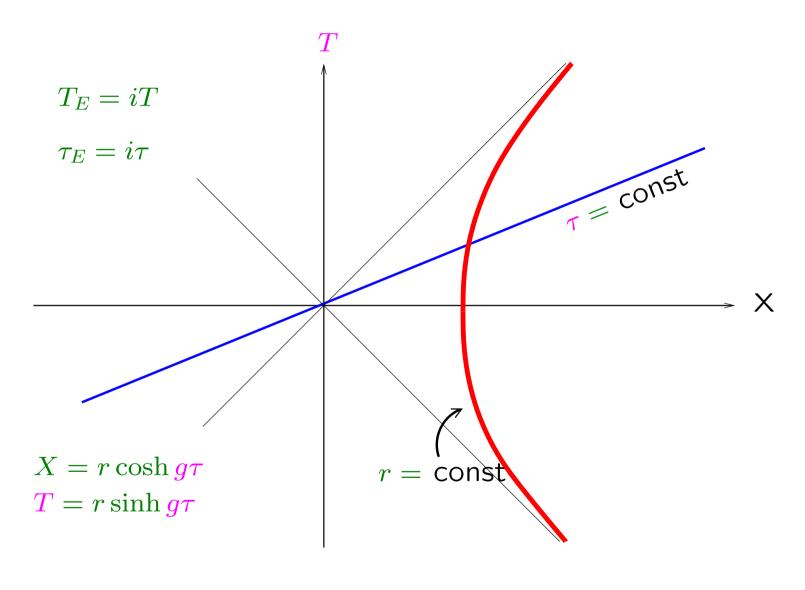
$$ds^2 = dY^2 + dX^2 = r^2 d\theta^2 + dr^2$$



$$ds^2 = dT_E^2 + dX^2 = g^2 r^2 d\tau_E^2 + dr^2$$



$$ds^{2} = -dT^{2} + dX^{2} = -g^{2}r^{2}d\tau^{2} + dr^{2}$$



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# PHYSICS PROGRESSES BY EXPLAINING FEATURES WHICH WE NEVER THOUGHT NEEDED ANY EXPLANATION !!

EXAMPLE:  $m_{inertial} = m_{grav}$ 

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• Read off (with  $L_P^2 \equiv G\hbar/c^3$ ):

[TP, 2002]

$$S = \frac{1}{4L_P^2}(4\pi a^2) = \frac{1}{4}\frac{A_H}{L_P^2}; \quad E = \frac{c^4}{2G}a = \frac{c^4}{G}\left(\frac{A_H}{16\pi}\right)^{1/2}$$

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• Works for Kerr, FRW, ....

[D. Kothawala et al., 06; Rong-Gen Cai, 06, 07]

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Action for gravity has exactly this structure!

[TP, 02, 05]

$$A_{grav} = \int d^4x \, \sqrt{-g} \, R = \int d^4x \, \sqrt{-g} \, [L_{\text{bulk}} + L_{\text{sur}}]$$

$$\sqrt{-g}L_{sur} = -\partial_a \left( g_{ij} \frac{\partial \sqrt{-g}L_{bulk}}{\partial(\partial_a g_{ij})} \right)$$

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  m sur}$ , evaluated on any horizon gives its entropy!
- In fact, one can develop a theory with  $A_{total} = A_{sur} + A_{matter}$  using the virtual displacements of the horizon as key. [TP, 2005]

• In a Riemann normal coordinates around any event  $\mathcal{P}$ , we have  $g \to \eta + R \ x \ x$  and  $L \sim \Gamma^2 + \partial \Gamma \to \partial \Gamma$  - Action is pure surface term! And  $\partial_a P_b \sim R_{ab}$ .

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- This leads to  $n^a\partial_a(n^bP_b)=n^an^bT_{ab}$  which leads to Einstein's equations! [TP, 2005,08]

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$$\mathcal{L}^{(D)} = Q_a^{\ bcd} R^a_{\ bcd} = \sum_{m=1}^K c_m \mathcal{L}_m^{(D)} \ ; \ \mathcal{L}_m^{(D)} = \frac{1}{16\pi} 2^{-m} \delta^{a_1 a_2 \dots a_{2m}}_{b_1 b_2 \dots b_{2m}} R^{b_1 b_2}_{a_1 a_2} \dots R^{b_{2m-1} b_{2m}}_{a_{2m-1} a_{2m}} \,,$$

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The surface term is closely related to horizon entropy in Lanczos-Lovelock theory.

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- ullet No. The virtual displacement of a null surface should cost entropy,  $S_{grav}.$

## REWRITING HISTORY: GRAVITY - THE 'RIGHT WAY UP'

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- Around any event there exists local inertial frames AND local Rindler frames with a local horizon and temperature.
- Can flow of matter across the local, hot, horizon hide entropy?
- Equivalently, can virtual displacements of a local patch of null surface, leading to flow of energy across a hot horizon allow you to hide entropy?
- ullet No. The virtual displacement of a null surface should cost entropy,  $S_{grav}.$
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- ullet No. The virtual displacement of a null surface should cost entropy,  $S_{grav}.$
- Dynamics should now emerge from maximising  $S_{matter} + S_{grav}$  for all Rindler observers!.
- Leads to gravity being an emergent phenomenon described by Einstein's equations at lowest order with calculable corrections.

Associate with virtual displacements of null surfaces an entropy/ action which is
quadratic in deformation field:
 [T.P, 08; T.P., A.Paranjape, 07]

$$S[\xi] = S[\xi]_{grav} + S_{matt}[\xi]$$

with

$$S_{grav}[\xi] = \int_{\mathcal{V}} d^D x \sqrt{-g} 4P^{abcd} \nabla_c \xi_a \nabla_d \xi_b; \qquad S_{matt} = \int_{\mathcal{V}} d^D x \sqrt{-g} T^{ab} \xi_a \xi_b$$

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- This leads to  $P^{abcd}$  having a (RG-like) derivative expansion in powers of number of derivatives of the metric:

$$P^{abcd}(g_{ij}, R_{ijkl}) = c_1 \stackrel{(1)}{P}^{abcd}(g_{ij}) + c_2 \stackrel{(2)}{P}^{abcd}(g_{ij}, R_{ijkl}) + \cdots,$$

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- Example: The lowest order term is:

$$S_1[\xi] = \int_{\mathcal{V}} rac{d^D x}{8\pi} \left( 
abla_a \xi^b 
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ight)$$

• Demand that  $\delta S=0$  for variations of all null vectors: This leads to Lanczos-Lovelock theory with an arbitrary cosmological constant:

$$16\pi \left[ P_b^{\ ijk} R^a_{\ ijk} - \frac{1}{2} \delta^a_b \mathcal{L}_m^{(D)} \right] = 8\pi T_b^a + \Lambda \delta^a_b,$$

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 In a derivative coupling expansion, Lanczos-Lovelock terms are calculable corrections.

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- Comparison with quasi-normal modes approach shows that it is the gravitational entropy which is quantised in Lanczos-Lovelock theories.

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- The only way out is to have a formalism for gravity which is invariant under  $T_{ab} \to T_{ab} + \rho g_{ab}$ .
- All these have nothing to do with observations of accelerated universe!
   Cosmological constant problem existed earlier and will continue to exist even if all these observations go away!

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- It also makes the variational principle for Lanczos-Lovelock theories well-defined.

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- Introduces a new length scale  $L_H$ . (Observationally,  $L_P/L_H \approx 10^{-60} \approx \exp(-\sqrt{2}\pi^4)$ .)
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Spacetime

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Connection with	Specify the entropy	Specify the entropy
thermodynamics		
Resulting equation	Classical / Quantum	Einsteins theory with
		calculable corrections

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- ullet Connects with  $A_{
  m sur}$  giving the horizon entropy; leads to quantisation of Wald entropy.

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