

# Understanding our universe – From the early epochs to late times –

L. Sriramkumar

Department of Physics, Indian Institute of Technology Madras, Chennai

IGCAR, Kalpakkam

November 19, 2015

# Plan of the talk

## 1 A survey of the universe



# Plan of the talk

- 1 A survey of the universe
- 2 The composition and evolution of the smooth universe



# Plan of the talk

- 1 A survey of the universe
- 2 The composition and evolution of the smooth universe
- 3 The origin and evolution of perturbations



# Plan of the talk

- 1 A survey of the universe
- 2 The composition and evolution of the smooth universe
- 3 The origin and evolution of perturbations
- 4 The standard model of cosmology

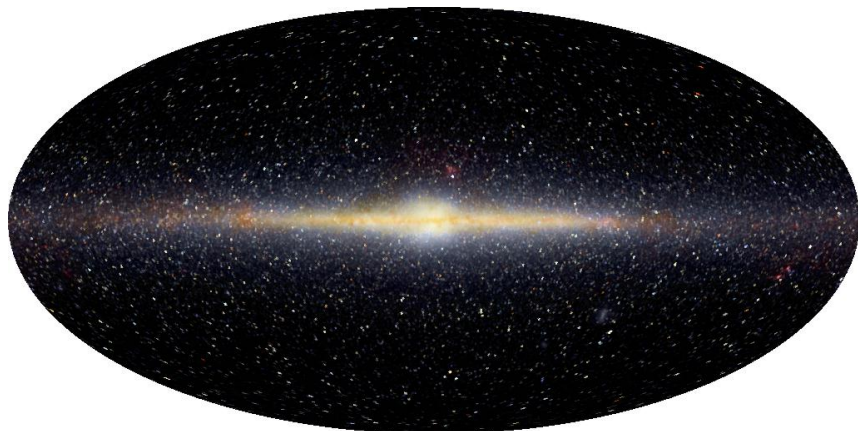


# Plan of the talk

- 1 A survey of the universe
  - Our galaxy and the local group
  - The local cluster, and beyond
  - A global view of the universe
- 2 The composition and evolution of the smooth universe
- 3 The origin and evolution of perturbations
- 4 The standard model of cosmology



# An infrared image of our galaxy



Our galaxy – the Milky Way – as observed by the **COsmic Background Explorer (COBE)** satellite at the infrared wavelengths<sup>1</sup>. The diameter of the disc of our galaxy is, approximately,  $45 \times 10^3$  ly or 15 kpc (*i.e.* a kilo parsec). It contains about  $10^{11}$  stars such as the Sun, and its mass is about  $2 \times 10^{12} M_{\odot}$ .

<sup>1</sup> Image from [http://aether.lbl.gov/www/projects/cobe/cobe\\_pics.html](http://aether.lbl.gov/www/projects/cobe/cobe_pics.html).



# Our galactic neighbors and the local group<sup>2</sup>



**Left:** The Andromeda galaxy and its two companion galaxies. The Andromeda galaxy is very similar to our galaxy and is located at a distance of about **700** kpc.

---

<sup>2</sup>Images from <http://www.seds.org/messier/m/m031.html> and <http://www.seds.org/messier/m/m033.html>.





# Our galactic neighbors and the local group<sup>2</sup>



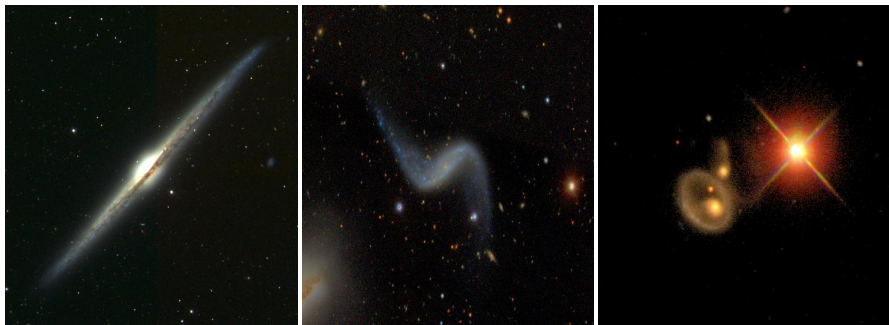
**Left:** The Andromeda galaxy and its two companion galaxies. The Andromeda galaxy is very similar to our galaxy and is located at a distance of about **700** kpc.

**Right:** The Triangulum galaxy. These galaxies, along with our galaxy, are major members of a local group of about **30** galaxies that are bound gravitationally. The size of the local group is estimated to be about **1.3** Mpc.

<sup>2</sup>Images from <http://www.seds.org/messier/m/m031.html> and <http://www.seds.org/messier/m/m033.html>.



# Varieties of galaxies<sup>3</sup>

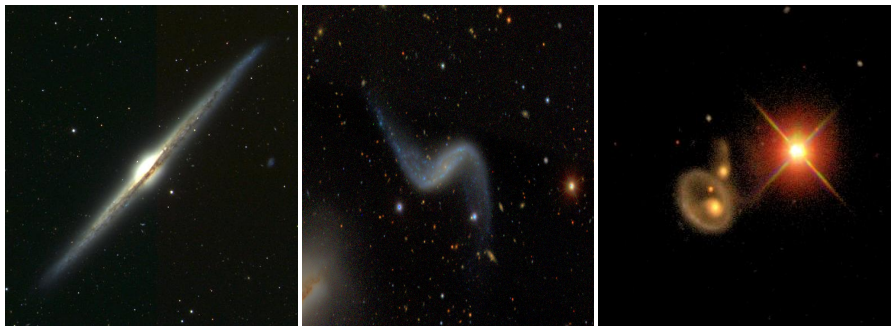


**Left:** The disk galaxy NGC 4565 seen edge on in this image from the **Sloan Digital Sky Survey (SDSS)**. The galaxy has a clear bulge, but little light can be seen from its halo.

<sup>3</sup>Images from <http://www.sdss.org/iotw/archive.html> and <http://cosmo.nyu.edu/hogg/rc3>.



# Varieties of galaxies<sup>3</sup>



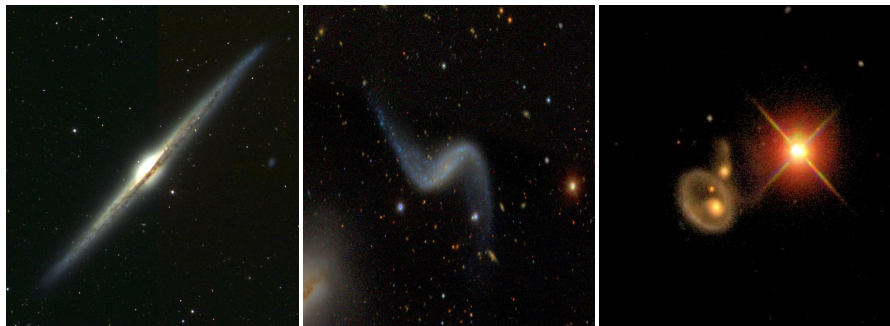
**Left:** The disk galaxy NGC 4565 seen edge on in this image from the **Sloan Digital Sky Survey (SDSS)**. The galaxy has a clear bulge, but little light can be seen from its halo.

**Center:** An image of the spiral galaxy NGC 3187 from SDSS.

<sup>3</sup>Images from <http://www.sdss.org/iotw/archive.html> and <http://cosmo.nyu.edu/hogg/rc3>.



# Varieties of galaxies<sup>3</sup>



**Left:** The disk galaxy NGC 4565 seen edge on in this image from the **Sloan Digital Sky Survey (SDSS)**. The galaxy has a clear bulge, but little light can be seen from its halo.

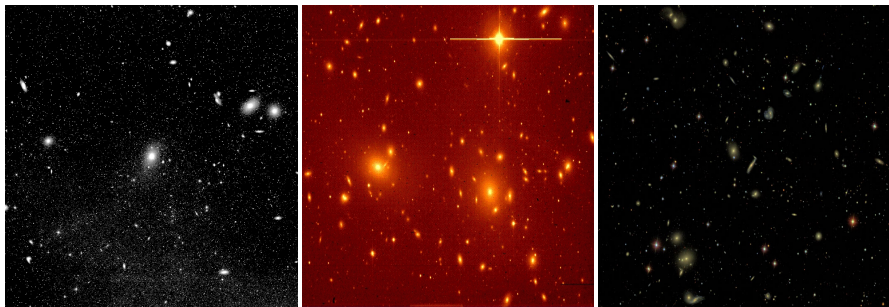
**Center:** An image of the spiral galaxy NGC 3187 from SDSS.

**Right:** CGCG 180-023 is a superb example of a ring galaxy. Ring galaxies are believed to form when a compact smaller galaxy plunges through the center of a larger more diffuse rotating disk galaxy.

<sup>3</sup>Images from <http://www.sdss.org/iotw/archive.html> and <http://cosmo.nyu.edu/hogg/rc3>.



# The Virgo, the Coma and the Hercules cluster of galaxies<sup>4</sup>

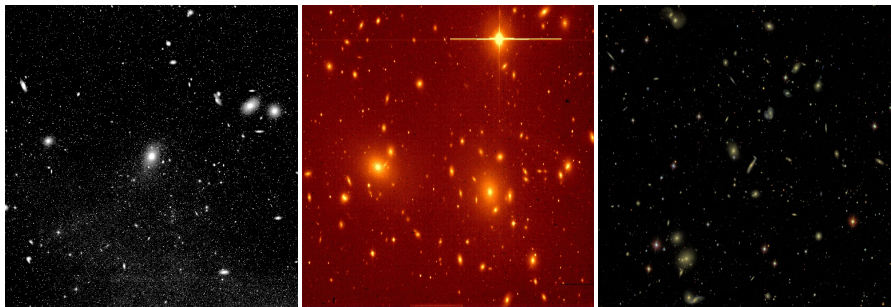


**Left:** The Virgo cluster, whose center is considered to be located at a distance of about 20 Mpc. Consisting of over 100 galaxies, it strongly influences the nearby galaxies and galaxy groups gravitationally due to its enormous mass.

<sup>4</sup> Images from <http://apod.nasa.gov/apod/ap000220.html>, <http://www.astr.ua.edu/gifimages/coma.html> and <http://www.sdss.org/iotw/archive.html>.



# The Virgo, the Coma and the Hercules cluster of galaxies<sup>4</sup>



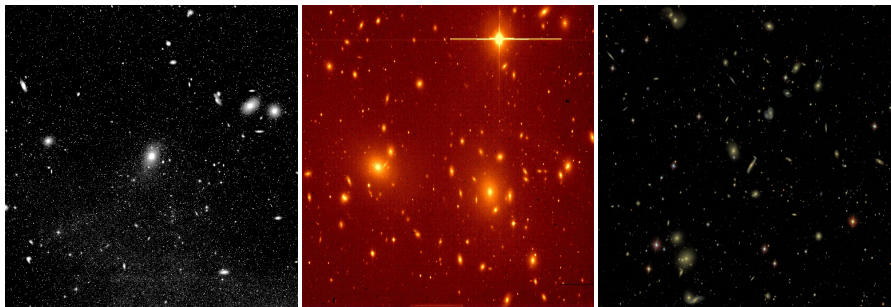
**Left:** The Virgo cluster, whose center is considered to be located at a distance of about 20 Mpc. Consisting of over 100 galaxies, it strongly influences the nearby galaxies and galaxy groups gravitationally due to its enormous mass.

**Center:** The Coma cluster of galaxies. The cluster is nearly spherical in shape and contains more than 1000 bright galaxies. It is about 20 Mpc across, and is located at a distance of about 100 Mpc.

<sup>4</sup> Images from <http://apod.nasa.gov/apod/ap000220.html>, <http://www.astr.ua.edu/gifimages/coma.html> and <http://www.sdss.org/iotw/archive.html>.



# The Virgo, the Coma and the Hercules cluster of galaxies<sup>4</sup>



**Left:** The Virgo cluster, whose center is considered to be located at a distance of about **20** Mpc. Consisting of over **100** galaxies, it strongly influences the nearby galaxies and galaxy groups gravitationally due to its enormous mass.

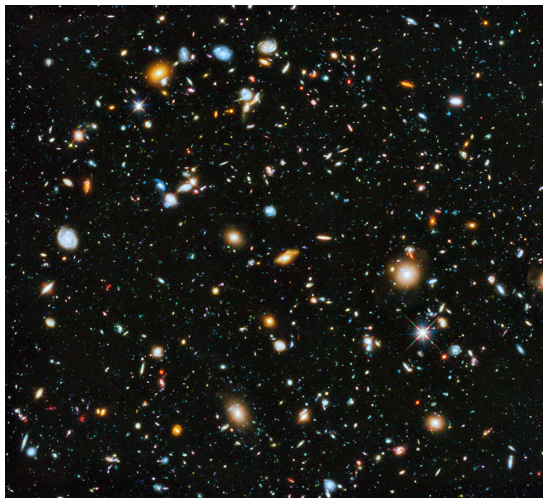
**Center:** The Coma cluster of galaxies. The cluster is nearly spherical in shape and contains more than **1000** bright galaxies. It is about **20** Mpc across, and is located at a distance of about **100** Mpc.

**Right:** An SDSS image of the Hercules galaxy cluster that is located at a distance of about **100** Mpc from us.

<sup>4</sup> Images from <http://apod.nasa.gov/apod/ap000220.html>, <http://www.astr.ua.edu/gifimages/coma.html> and <http://www.sdss.org/iotw/archive.html>.



# Deepest views in space



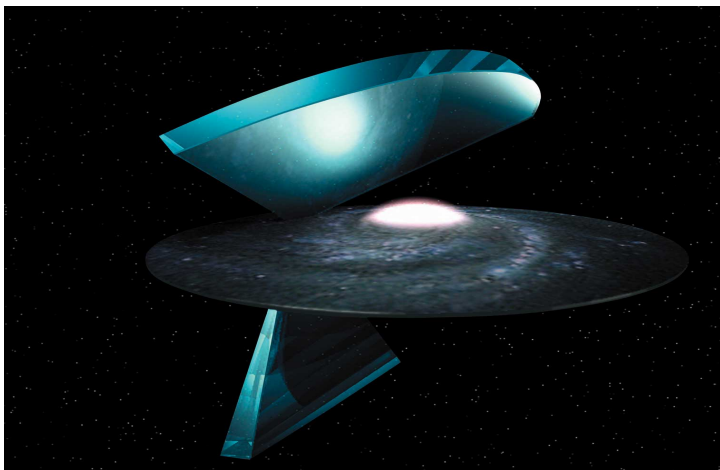
An ultra deep field image from the **Hubble Space Telescope (HST)**. The image contains a bewildering variety of galaxy shapes and colors<sup>5</sup>.

<sup>5</sup>Image from <http://hubblesite.org/newscenter/archive/releases/2014/27>.





# Surveying the universe

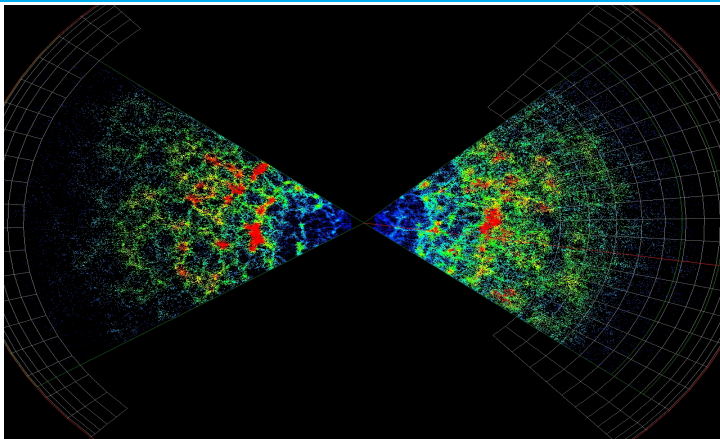


A schematic drawing showing the directions of the regions observed by the **2 degree field (2dF) redshift survey** with respect to our galaxy<sup>6</sup>. The survey regions actually extend more than  $10^5$  times further than shown here.

<sup>6</sup>Image from <http://magnum.anu.edu.au/~TDFgg/Public/Pics/2dF3D.jpg>.



# Distribution of galaxies in the universe



The distribution of more than two million galaxies as observed by the 2dF redshift survey<sup>7</sup>. (Note that each dot in the picture represents a galaxy.) The density and the 'radius' of the universe are estimated to be about  $10^{-28}$  kg/m<sup>3</sup> and 3000 Mpc, respectively.

<sup>7</sup>Image from [http://magnum.anu.edu.au/~TDFgg/Public/Pics/2dFGRS\\_top\\_view.gif](http://magnum.anu.edu.au/~TDFgg/Public/Pics/2dFGRS_top_view.gif).



# The Sloan digital sky survey

- The **Sloan Digital Sky Survey (SDSS)** is one of the most ambitious and influential surveys in the history of astronomy.



# The Sloan digital sky survey

- The **Sloan Digital Sky Survey (SDSS)** is one of the most ambitious and influential surveys in the history of astronomy.
- Over eight years of operations, it has obtained deep, multi-color images covering more than a quarter of the sky and created three-dimensional maps containing more than **930,000** galaxies and more than **120,000** quasars.

▶ [Play SDSS movie](#)



# Plan of the talk

- 1 A survey of the universe
- 2 **The composition and evolution of the smooth universe**
  - The expanding universe and the Hubble's law
  - Describing and characterizing the universe
  - The cosmic microwave background
  - The hot big bang model
  - Decoupling of matter and radiation
  - The baryon content of the universe
  - Why do we require dark energy?
  - The composition of the universe
- 3 The origin and evolution of perturbations
- 4 The standard model of cosmology



# Continuous, emission and absorption spectra<sup>8</sup>

A typical continuous spectrum from an opaque hot body:



---

<sup>8</sup>Images from <http://hea-www.harvard.edu/~efortin/thesis/html/Spectroscopy.shtml>.

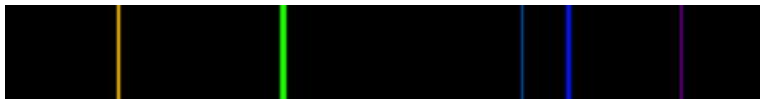


# Continuous, emission and absorption spectra<sup>8</sup>

A typical continuous spectrum from an opaque hot body:



Emission spectrum, as from a given element:



<sup>8</sup>Images from <http://hea-www.harvard.edu/~efortin/thesis/html/Spectroscopy.shtml>.

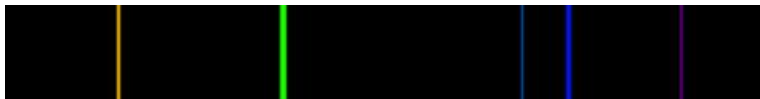


# Continuous, emission and absorption spectra<sup>8</sup>

A typical continuous spectrum from an opaque hot body:



Emission spectrum, as from a given element:



Absorption spectrum, as due to an intervening cool gas:

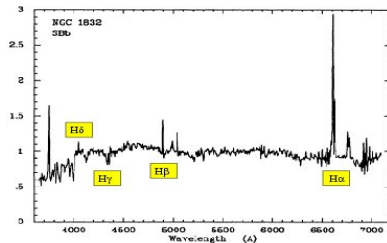
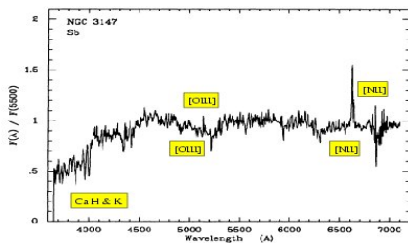
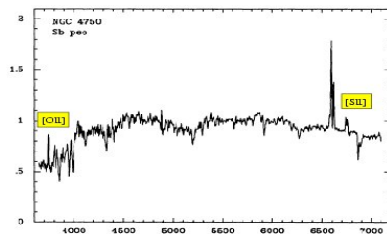
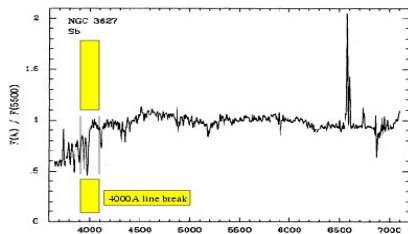


<sup>8</sup>Images from <http://hea-www.harvard.edu/~efortin/thesis/html/Spectroscopy.shtml>.





# Typical spectra of galaxies<sup>9</sup>



Spectra of some spiral galaxies. The spectra usually contain characteristic emission and absorption lines.

<sup>9</sup>Image from <http://astronomy.nmsu.edu/nicole/teaching/ASTR505/lectures/lecture26/slide01.html>.



# The 'Doppler effect' and redshift<sup>10</sup>

If the source is receding, the spectrum will be red-shifted



<sup>10</sup>Images from <http://www.astronomynotes.com/light/s10.htm>.



# The 'Doppler effect' and redshift<sup>10</sup>

If the source is receding, the spectrum will be red-shifted



when compared to the spectrum in the source's frame



<sup>10</sup>Images from <http://www.astronomynotes.com/light/s10.htm>.



# The 'Doppler effect' and redshift<sup>10</sup>

If the source is receding, the spectrum will be red-shifted



when compared to the spectrum in the source's frame



The redshift  $z$  of the receding source is defined as:

$$1 + z = \frac{\lambda_O}{\lambda_E} = \frac{\omega_E}{\omega_O},$$

where  $\lambda_O$  and  $\omega_O$  denote the observed wavelength and frequency of the source, while  $\lambda_E$  and  $\omega_E$  denote its emitted wavelength and frequency, respectively.

<sup>10</sup>Images from <http://www.astronomynotes.com/light/s10.htm>.

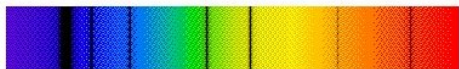


# Runaway galaxies – A schematic diagram<sup>11</sup>

A distant galaxy,  $z = 0.25$



A farther galaxy,  $z = 0.05$



A nearby galaxy,  $z = 0.01$



A galactic star,  $z = 0$



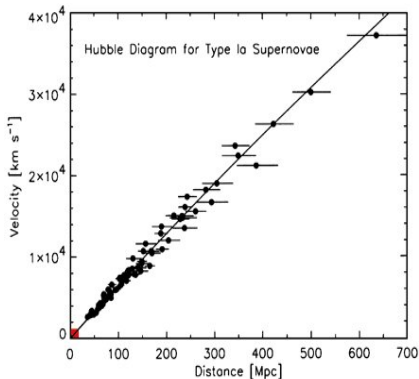
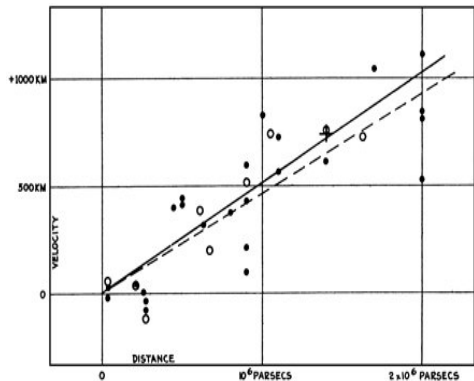
In the above spectrum of the galactic star, the wavelengths of the absorption lines are **393** and **397** nm from Ca II (ionized calcium); **410**, **434**, **486** and **656** nm from H I (atomic hydrogen); **518** nm from Mg I (neutral magnesium); and **589** nm from Na I (neutral sodium).

<sup>11</sup> Image from <http://www.astro.ucla.edu/~wright/doppler.htm>.





# Relation between the velocity and the distance of galaxies<sup>13</sup>

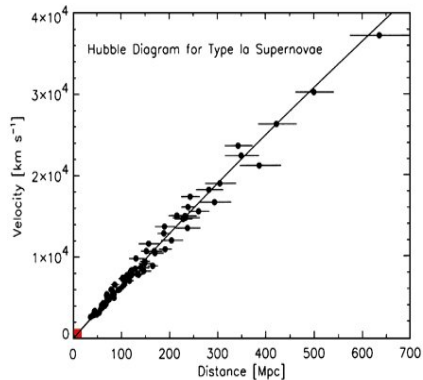
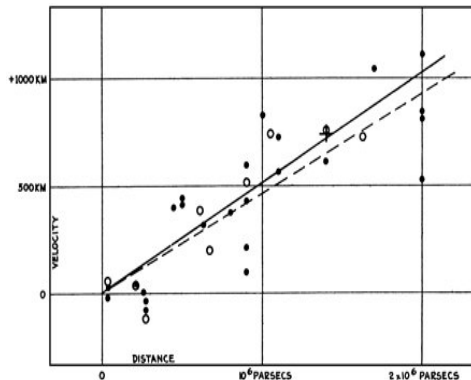


**Left:** The original Hubble data. The slope of the two fitted lines are about **500** km/sec/Mpc and **530** km/sec/Mpc.

<sup>13</sup>R. Kirshner, Proc. Natl. Acad. Sci. USA **101**, 8 (2004).



# Relation between the velocity and the distance of galaxies<sup>13</sup>



**Left:** The original Hubble data. The slope of the two fitted lines are about **500** km/sec/Mpc and **530** km/sec/Mpc.

**Right:** A more recent Hubble diagram. The slope of the straight line is found to be about **72** km/sec/Mpc. The small red region in the lower left marks the span of Hubble's original diagram.

<sup>13</sup>R. Kirshner, Proc. Natl. Acad. Sci. USA **101**, 8 (2004).





# The Friedmann-Lemaître-Robertson-Walker metric

The homogeneous, isotropic and expanding universe can be described by the following Friedmann-Lemaître-Robertson-Walker (FLRW) line element:

$$ds^2 = dt^2 - a^2(t) \left[ \frac{dr^2}{(1 - \kappa r^2)} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right],$$

where  $t$  is the cosmic time and  $a(t)$  denotes the scale factor, while  $\kappa = 0, \pm 1$ .

---

<sup>14</sup>Image from [http://abyss.uoregon.edu/~js/lectures/cosmo\\_101.html](http://abyss.uoregon.edu/~js/lectures/cosmo_101.html).



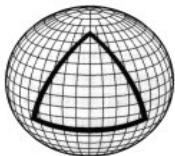
# The Friedmann-Lemaître-Robertson-Walker metric

The homogeneous, isotropic and expanding universe can be described by the following Friedmann-Lemaître-Robertson-Walker (FLRW) line element:

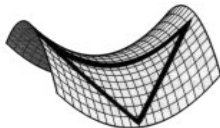
$$ds^2 = dt^2 - a^2(t) \left[ \frac{dr^2}{(1 - \kappa r^2)} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right],$$

where  $t$  is the cosmic time and  $a(t)$  denotes the scale factor, while  $\kappa = 0, \pm 1$ .

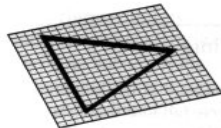
The quantity  $\kappa$  denotes the spatial geometry of the universe. It can be flat ( $\kappa = 0$ ), closed ( $\kappa = 1$ ) or open ( $\kappa = -1$ ) depending on the total energy density of matter present in the universe<sup>14</sup>.



**Positive Curvature**



**Negative Curvature**



**Flat Curvature**

<sup>14</sup>Image from [http://abyss.uoregon.edu/~js/lectures/cosmo\\_101.html](http://abyss.uoregon.edu/~js/lectures/cosmo_101.html).



# The Friedmann equations

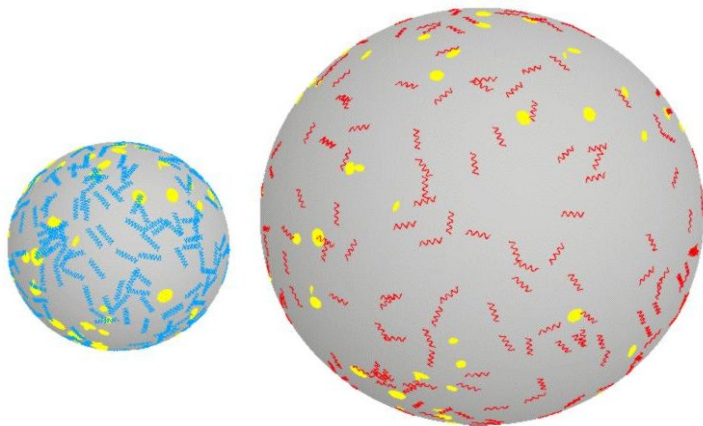
If  $\rho$  and  $p$  denote the energy density and pressure of the smooth component of the matter field that is driving the expansion, then the Einstein's equations for the FLRW metric lead to the following equations for the scale factor  $a(t)$ :

$$H^2 + \frac{\kappa}{a^2} = \frac{8\pi G}{3} \rho \quad \text{and} \quad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p),$$

where  $H = \dot{a}/a$  is the Hubble parameter.



# Visualizing the expanding universe



A two-dimensional analogy for the expanding universe<sup>15</sup>. The yellow blobs on the expanding balloon denote the galaxies. Note that the galaxies themselves do not grow, but the distance between the galaxies grows and the wavelengths of the photons shift from blue to red as the universe expands.

<sup>15</sup>Image from <http://www.astro.ucla.edu/~wright/balloon0.html>.



# The cosmological redshift

Recall that, we had defined the redshift  $z$  of a receding source as follows:

$$1 + z = \frac{\omega_E}{\omega_O},$$

where  $\omega_O$  and  $\omega_E$  denote the observed and emitted frequencies, respectively.



# The cosmological redshift

Recall that, we had defined the redshift  $z$  of a receding source as follows:

$$1 + z = \frac{\omega_E}{\omega_O},$$

where  $\omega_O$  and  $\omega_E$  denote the observed and emitted frequencies, respectively.

In an expanding universe, it can be shown that the frequency of electromagnetic radiation decreases with the expansion as follows:

$$\omega(t) \propto \frac{1}{a(t)},$$

where  $a(t)$  is the scale factor that characterizes the expansion.



# The cosmological redshift

Recall that, we had defined the redshift  $z$  of a receding source as follows:

$$1 + z = \frac{\omega_E}{\omega_O},$$

where  $\omega_O$  and  $\omega_E$  denote the observed and emitted frequencies, respectively.

In an expanding universe, it can be shown that the frequency of electromagnetic radiation decreases with the expansion as follows:

$$\omega(t) \propto \frac{1}{a(t)},$$

where  $a(t)$  is the scale factor that characterizes the expansion.

Therefore, in terms of the scale factor, the cosmological redshift  $z$  is given by

$$\frac{a_0}{a(t)} = 1 + z,$$

where  $a_0$  denotes the value of the scale factor *today* (i.e. at  $t = t_0$ ).



# The cosmological parameters

In terms of the redshift  $z$ , the first of the Friedmann equations can be written as

$$\left[ \frac{H(z)}{H_0} \right]^2 = \Omega_{\text{NR}} (1+z)^3 + \Omega_{\text{R}} (1+z)^4 + \Omega_{\Lambda} - (\Omega - 1) (1+z)^2,$$

where  $H_0 \equiv (\dot{a}/a)_{t=t_0}$  is the Hubble constant,  $\Omega_i = \rho_i/\rho_c$  with  $\rho_c$  being the critical density given by

$$\rho_c = \frac{3 H_0^2}{8 \pi G}$$

and  $\Omega = \Omega_{\text{NR}} + \Omega_{\text{R}} + \Omega_{\Lambda}$ .





# The cosmological parameters

In terms of the redshift  $z$ , the first of the Friedmann equations can be written as

$$\left[ \frac{H(z)}{H_0} \right]^2 = \Omega_{\text{NR}} (1+z)^3 + \Omega_{\text{R}} (1+z)^4 + \Omega_{\Lambda} - (\Omega - 1) (1+z)^2,$$

where  $H_0 \equiv (\dot{a}/a)_{t=t_0}$  is the Hubble constant,  $\Omega_i = \rho_i/\rho_c$  with  $\rho_c$  being the critical density given by

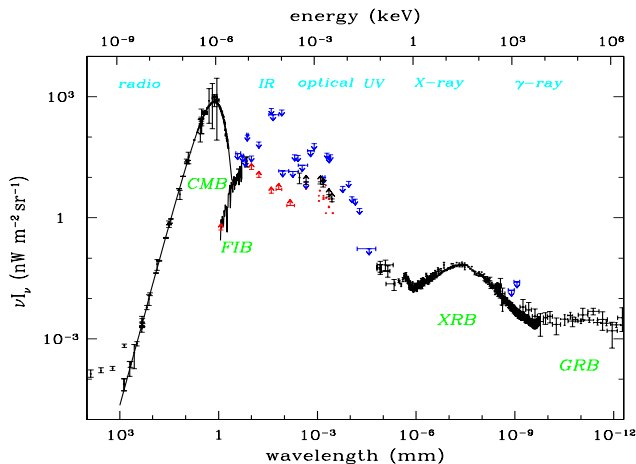
$$\rho_c = \frac{3 H_0^2}{8 \pi G}$$

and  $\Omega = \Omega_{\text{NR}} + \Omega_{\text{R}} + \Omega_{\Lambda}$ .

The quantities  $H_0$ ,  $\Omega_{\text{NR}}$ ,  $\Omega_{\text{R}}$  and  $\Omega_{\Lambda}$  are four of the cosmological parameters that are to be determined by observations.



# The Cosmic Microwave Background (CMB)

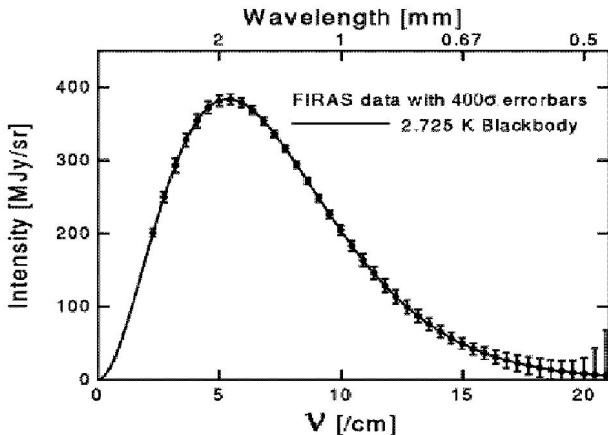


The energy density spectrum of cosmological background radiation has been plotted as a function of wavelength<sup>16</sup>. Note that the CMB contributes the most to the overall background radiation.

<sup>16</sup>Figure from, [D. Scott, arXiv:astro-ph/9912038](https://arxiv.org/abs/astro-ph/9912038).



# The spectrum of the CMB

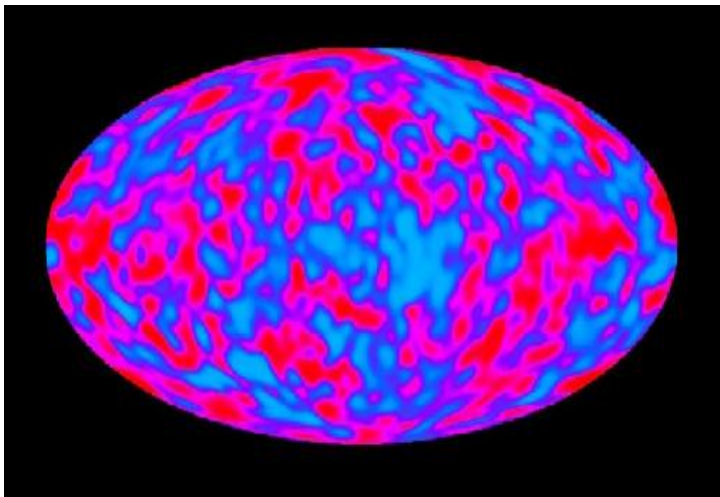


The spectrum of the CMB as measured by the **COBE satellite**<sup>17</sup>. It is such a perfect Planck spectrum (corresponding to a temperature of **2.725° K**) that it is unlikely to be bettered in the laboratory. The error bars in the graph above have been amplified **400** times so that they can be seen!

<sup>17</sup>Image from [http://www.astro.ucla.edu/~wright/cosmo\\_01.htm](http://www.astro.ucla.edu/~wright/cosmo_01.htm).



# The extent of isotropy of the CMB



The fluctuations in the temperature of the CMB as seen by COBE<sup>18</sup>. The CMB turns out to be isotropic to one part in  $10^5$ .

<sup>18</sup>Image from [http://aether.lbl.gov/www/projects/cobe/COBE\\_Home/DMR\\_Images.html](http://aether.lbl.gov/www/projects/cobe/COBE_Home/DMR_Images.html).



# Equilibrium between matter and radiation at early epochs

In an evolving universe, the temperature of the CMB goes as

$$T \propto \frac{1}{a(t)},$$

so that the energy density of radiation behaves as

$$\rho_{\text{R}} \propto \frac{1}{a^4(t)}.$$



# Equilibrium between matter and radiation at early epochs

In an evolving universe, the temperature of the CMB goes as

$$T \propto \frac{1}{a(t)},$$

so that the energy density of radiation behaves as

$$\rho_{\text{R}} \propto \frac{1}{a^4(t)}.$$

In contrast, the energy density of non-relativistic (*i.e.* pressureless) matter goes as

$$\rho_{\text{NR}} \propto \frac{1}{a^3(t)}.$$



## Equilibrium between matter and radiation at early epochs

In an evolving universe, the temperature of the CMB goes as

$$T \propto \frac{1}{a(t)},$$

so that the energy density of radiation behaves as

$$\rho_{\text{R}} \propto \frac{1}{a^4(t)}.$$

In contrast, the energy density of non-relativistic (*i.e.* pressureless) matter goes as

$$\rho_{\text{NR}} \propto \frac{1}{a^3(t)}.$$

Observations indicate that, today,

$$\rho_{\text{R}} \simeq \frac{\rho_{\text{NR}}}{10^4}.$$



## Equilibrium between matter and radiation at early epochs

In an evolving universe, the temperature of the CMB goes as

$$T \propto \frac{1}{a(t)},$$

so that the energy density of radiation behaves as

$$\rho_{\text{R}} \propto \frac{1}{a^4(t)}.$$

In contrast, the energy density of non-relativistic (*i.e.* pressureless) matter goes as

$$\rho_{\text{NR}} \propto \frac{1}{a^3(t)}.$$

Observations indicate that, today,

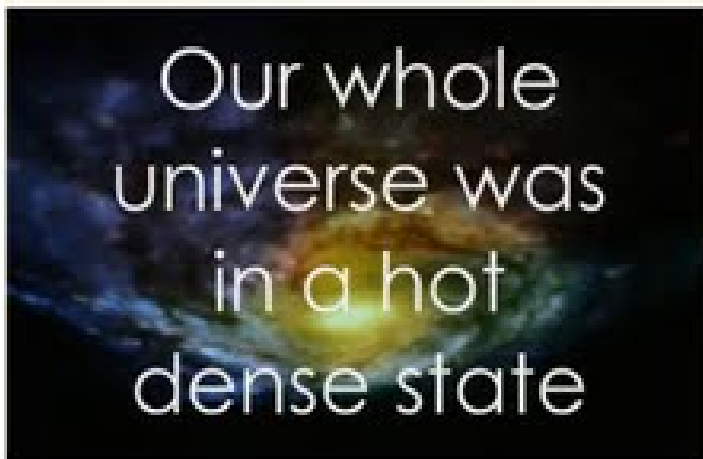
$$\rho_{\text{R}} \simeq \frac{\rho_{\text{NR}}}{10^4}.$$

This points to the fact that matter and radiation would have interacted strongly and, hence would have been in thermal equilibrium, when the universe was about  $10^4$  times smaller.





# The big bang model seems popular!

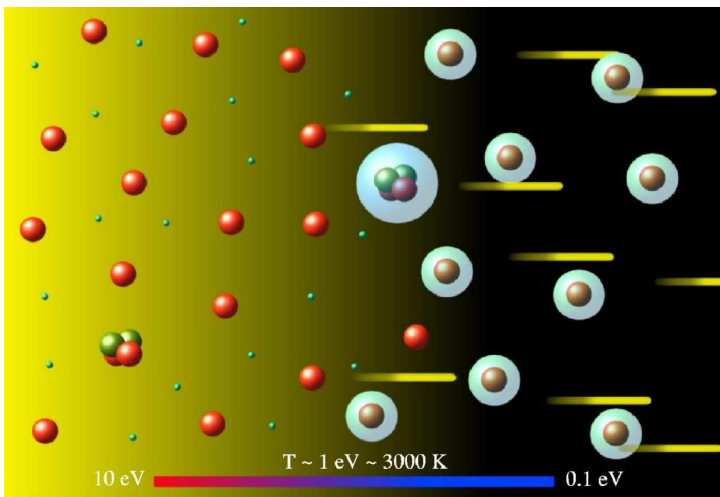


The current view of the universe, encapsulated in the hot big bang model, seems popular. The above image is a screen grab from the theme song of the recent American sitcom 'The Big Bang Theory'<sup>19</sup>!

<sup>19</sup>See [http://www.cbs.com/shows/big\\_bang\\_theory/](http://www.cbs.com/shows/big_bang_theory/).



# Decoupling of matter and radiation<sup>20</sup>

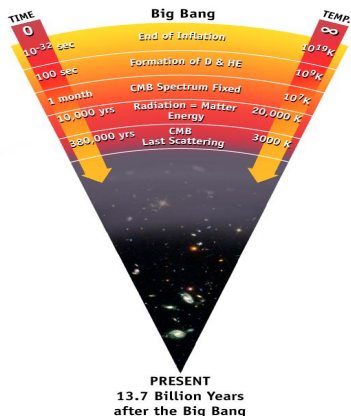


Matter and radiation cease to interact at a temperature of about  $T \simeq 3000^\circ \text{ K}$ , which corresponds to a redshift of about  $z \simeq 1000$ .

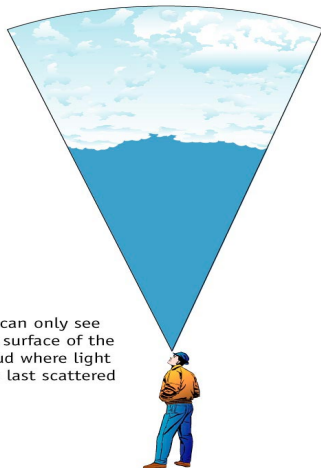
<sup>20</sup>Image from W. H. Kinney, [arXiv:astro-ph/0301448v2](https://arxiv.org/abs/astro-ph/0301448v2).



# The last scattering surface and the freestreaming CMB photons



The cosmic microwave background Radiation's "surface of last scatterer" is analogous to the light coming through the clouds to our eye on a cloudy day.



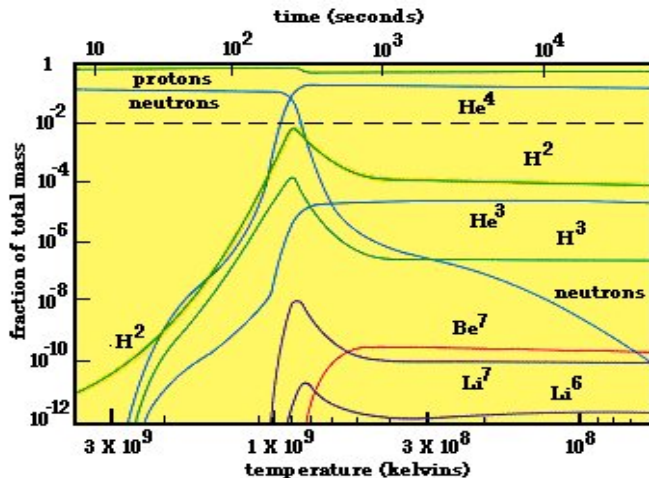
We can only see the surface of the cloud where light was last scattered

The CMB photons streams to us freely from the last scattering surface when radiation decoupled from matter<sup>21</sup>.

<sup>21</sup> Image from <http://map.gsfc.nasa.gov/media/990053/990053.jpg>.



# The abundance of light elements – Theory

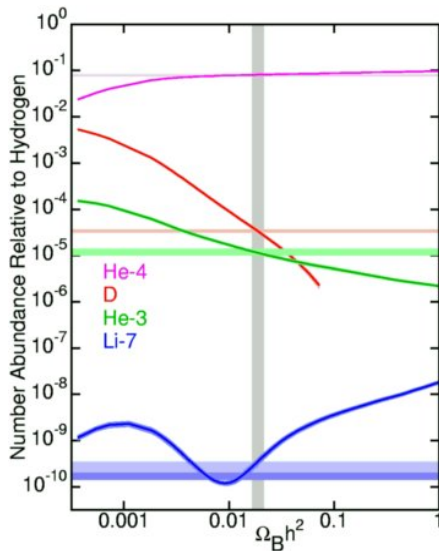


The relative abundances of the light elements in the early radiation dominated epoch have been plotted as a function of temperature<sup>22</sup>.

<sup>22</sup>Image from <http://www.astro.ucla.edu/~wright/BBNS.html>.



# Abundance of light elements – Observations<sup>23</sup>



The graph to the left contains the theoretically predicted abundance versus the density for the light elements as curves, the observed abundances as horizontal stripes and the derived baryon density as the vertical stripe. Note that a single value of the baryon density fits all the four abundances, and it is found that  $\Omega_B h^2 \simeq 0.022$ , where  $H_0 = 100 h$  km/sec/Mpc.

<sup>23</sup>Image from <http://www.astro.ucla.edu/~wright/BBNS.html>.



# Supernovae can be as bright as the host galaxy<sup>24</sup>



Supernova 1994D, visible as the bright spot on the lower left, occurred in the outskirts of disk galaxy NGC 4526.

<sup>24</sup>Image from <http://apod.nasa.gov/apod/ap981230.html>.



# A supernova explosion in a distant galaxy<sup>25</sup>



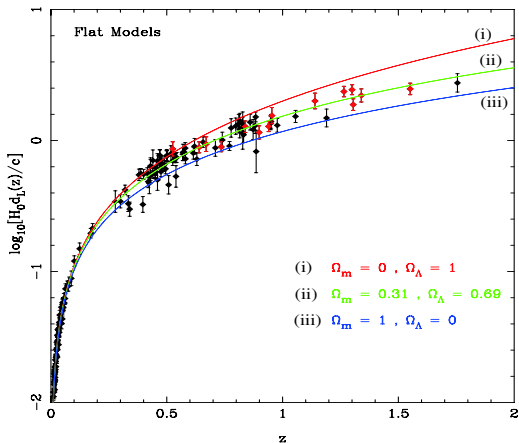
**Left:** A supernova at the redshift of **0.28** caught at maximum light by the **Supernova Legacy Survey (SNLS)**.

**Right:** The supernova after it has faded.

<sup>25</sup>Images from **C. J. Pritchett *et. al.*, arXiv:astro-ph/0406242v1.**



# Supernovae data and the need for a cosmological constant



The luminosity distance  $H_0 d_L$  plotted as a function of the redshift  $z$  for spatially flat cosmological models<sup>26</sup>. The black points are from the 'Gold' data sets and the red points are the data from the **Hubble Space Telescope**<sup>27</sup>.

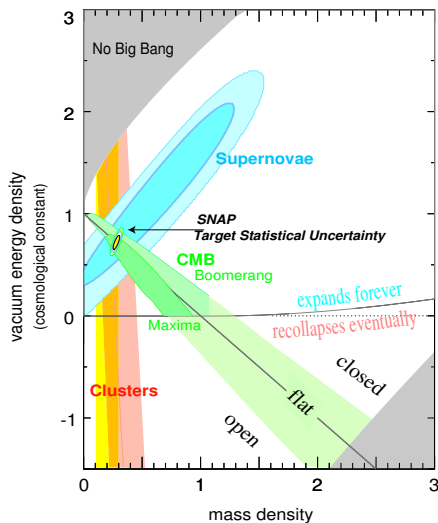
<sup>26</sup>Figure from T. R. Choudhury and T. Padmanabhan, *Astron. Astrophys.* **429**, 807 (2005).

<sup>27</sup>R. A. Knop *et. al.*, *Astrophys. J.* **598**, 102 (2003); A. G. Riess *et. al.*, *Astrophys. J.* **607**, 665 (2004).





# Joint constraints on $\Omega_{NR}$ and $\Omega_{\Lambda}$ <sup>28</sup>

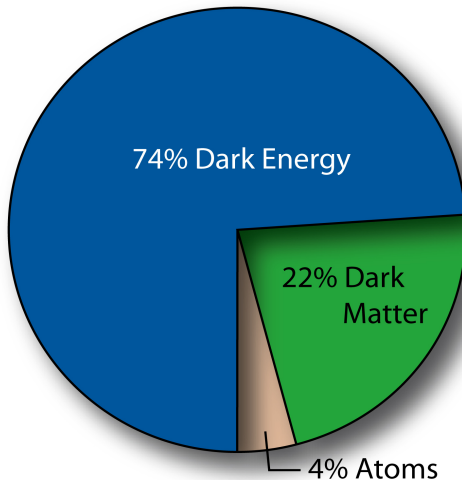


Joint constraints on  $\Omega_{NR}$  and  $\Omega_{\Lambda}$  from the observations of supernovae, CMB and galaxy clustering. Note that a cosmology with  $\Omega_{NR} = 1$  and  $\Omega_{\Lambda} = 0$  is ruled out to 99% confidence level, while a universe with  $\Omega_{NR} \simeq 0.3$  and  $\Omega_{\Lambda} \simeq 0.7$  proves to be a good fit to the data. The figure also contains the constraints that can be expected from the planned **Supernova/Acceleration Probe (SNAP)**.

<sup>28</sup>Figure from G. Aldering *et. al.*, arXiv:astro-ph/0209550v1.



# Matter content of the universe



A pie chart of the matter content of the universe today<sup>29</sup>.

<sup>29</sup>Image from [http://map.gsfc.nasa.gov/media/060916/060916\\_UniversePie300.jpg](http://map.gsfc.nasa.gov/media/060916/060916_UniversePie300.jpg).

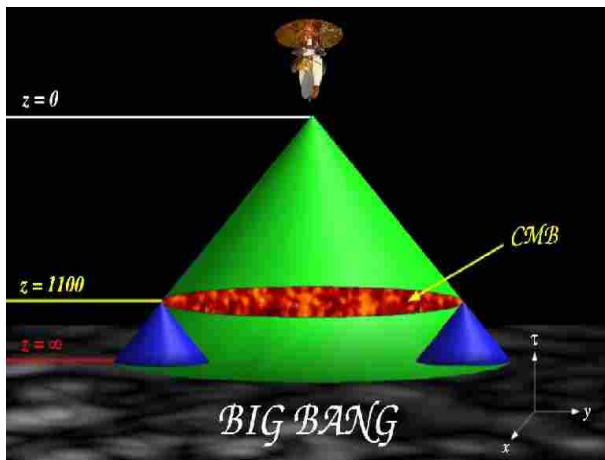


# Plan of the talk

- 1 A survey of the universe
- 2 The composition and evolution of the smooth universe
- 3 The origin and evolution of perturbations**
  - The need for an inflationary epoch
  - The generation and the evolution of perturbations
  - The universe according to WMAP and Planck
  - The formation of large scale structures
- 4 The standard model of cosmology



# The horizon problem

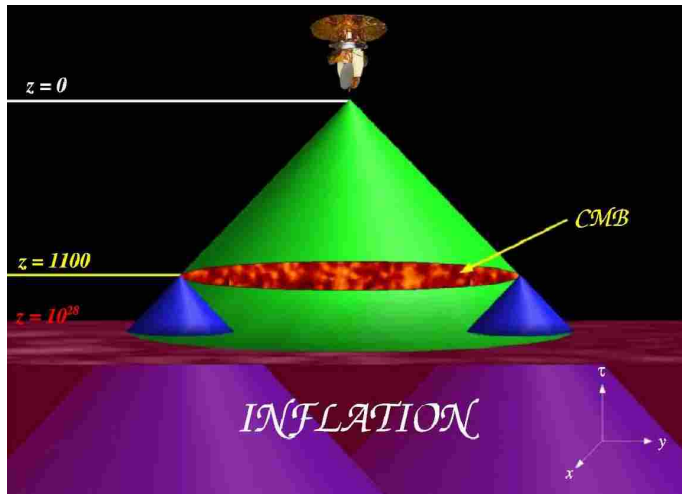


The radiation from the CMB arriving at us from regions separated by more than the Hubble radius at the last scattering surface (which subtends an angle of about  $1^\circ$  today) could not have interacted before decoupling<sup>30</sup>.

<sup>30</sup>Image from W. H. Kinney, [arXiv:astro-ph/0301448v2](https://arxiv.org/abs/astro-ph/0301448v2).



# Inflation resolves the horizon problem



An illustration of how an early and sufficiently long epoch of inflation (*viz.* a phase when  $\ddot{a} > 0$ ) resolves the horizon problem<sup>31</sup>.

<sup>31</sup> Image from W. H. Kinney, [arXiv:astro-ph/0301448v2](https://arxiv.org/abs/astro-ph/0301448v2).



# The origin and the evolution of the perturbations

- Inflation is typically driven with the aid of scalar fields. It is the quantum fluctuations associated with these scalar fields which are responsible for the origin of the perturbations.



# The origin and the evolution of the perturbations

- Inflation is typically driven with the aid of scalar fields. It is the quantum fluctuations associated with these scalar fields which are responsible for the origin of the perturbations.
- These perturbations are amplified during the inflationary epoch, which leave their imprints as anisotropies in the CMB.

[▶ Play movie](#)

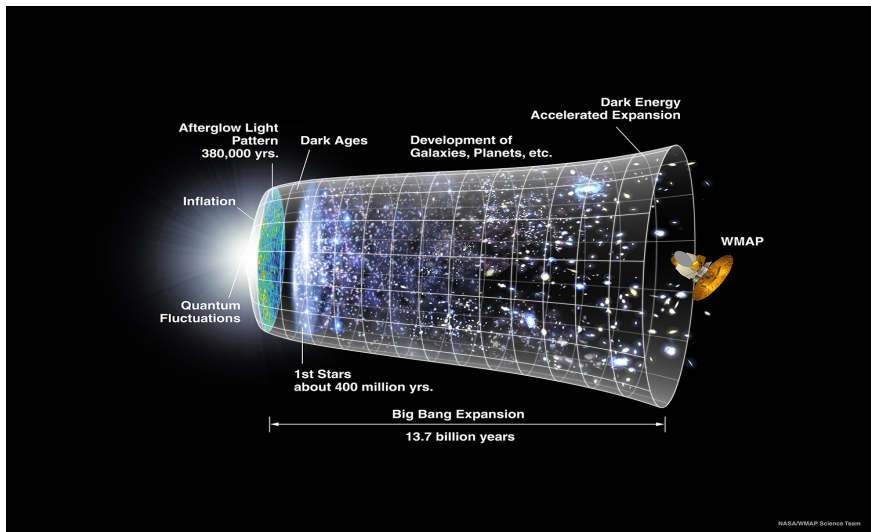
# The origin and the evolution of the perturbations

- Inflation is typically driven with the aid of scalar fields. It is the quantum fluctuations associated with these scalar fields which are responsible for the origin of the perturbations.
- These perturbations are amplified during the inflationary epoch, which leave their imprints as anisotropies in the CMB.
- The fluctuations in the CMB in turn grow in magnitude due to gravitational instability and develop into the structures that we see around us today.

[▶ Play movie](#)[▶ Play movie](#)



# The timeline of the universe

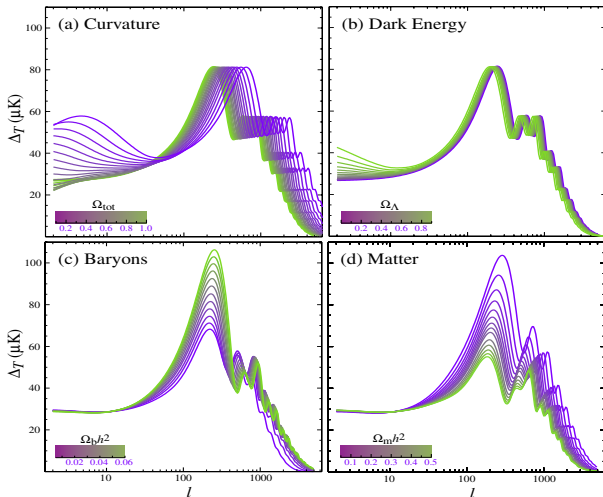


## A pictorial timeline of the universe<sup>32</sup>.

<sup>32</sup> See [http://wmap.gsfc.nasa.gov/media/060915/060915.CMB\\_Timeline150.jpg](http://wmap.gsfc.nasa.gov/media/060915/060915.CMB_Timeline150.jpg).



# 'Effects' of the cosmological parameters on the CMB<sup>33</sup>

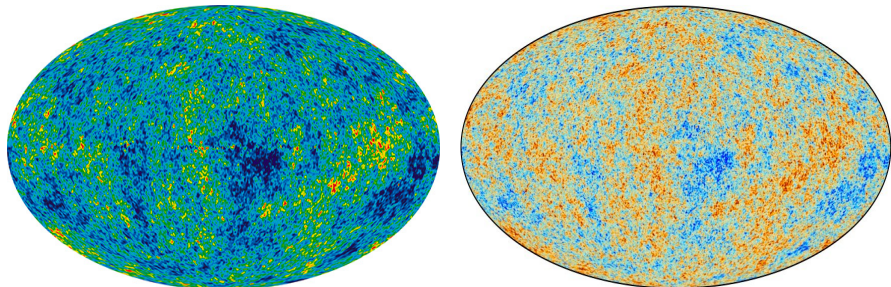


Sensitivity of the CMB angular power spectrum to the four cosmological parameters:  $\Omega$ ,  $\Omega_{\Lambda}$ ,  $\Omega_B h^2$  and the non-relativistic matter density  $\Omega_{\text{NR}} h^2$ .

<sup>33</sup>Figures from W. Hu and S. Dodelson, *Ann. Rev. Astron. Astrophys.* **40**, 171 (2002).



# CMB anisotropies as seen by WMAP and Planck



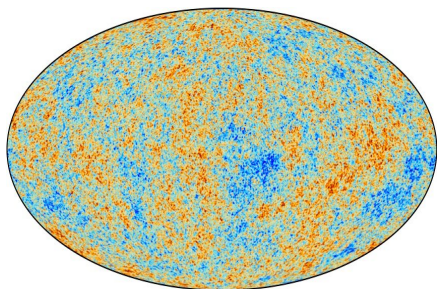
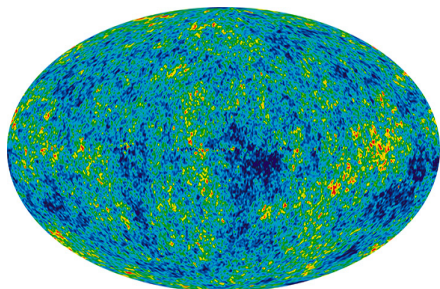
**Left:** All-sky map of the anisotropies in the CMB created from nine years of **Wilkinson Microwave Anisotropy Probe (WMAP)** data<sup>34</sup>.

<sup>34</sup>Image from <http://wmap.gsfc.nasa.gov/media/121238/index.html>.

<sup>35</sup>**P. A. R. Ade *et al.*, arXiv:1502.01582 [astro-ph.CO].**



# CMB anisotropies as seen by WMAP and Planck



**Left:** All-sky map of the anisotropies in the CMB created from nine years of **Wilkinson Microwave Anisotropy Probe (WMAP)** data<sup>34</sup>.

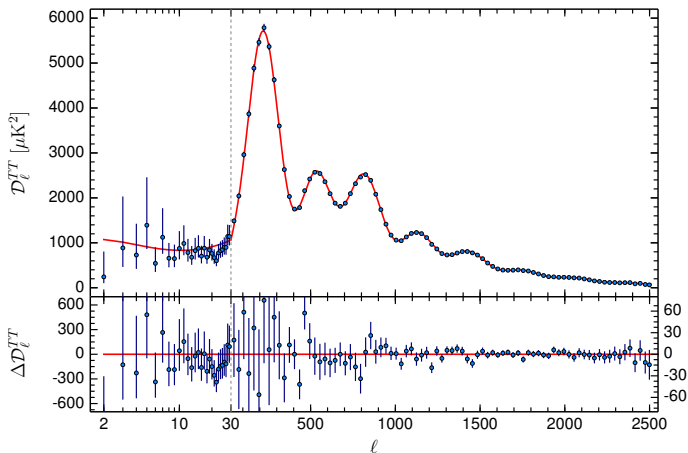
**Right:** CMB intensity map derived from the joint analysis of Planck, WMAP, and 408 MHz observations<sup>35</sup>. The above images show temperature variations (as color differences) of the order of  $200^\circ \mu\text{K}$ . The angular resolution of WMAP was about  $1^\circ$ , while that of Planck was about  $5'$ . These temperature fluctuations correspond to regions of slightly different densities, and they represent the seeds of all the structure around us today.

<sup>34</sup>Image from <http://wmap.gsfc.nasa.gov/media/121238/index.html>.

<sup>35</sup>P. A. R. Ade *et al.*, [arXiv:1502.01582 \[astro-ph.CO\]](https://arxiv.org/abs/1502.01582).



# The CMB angular power spectrum from Planck

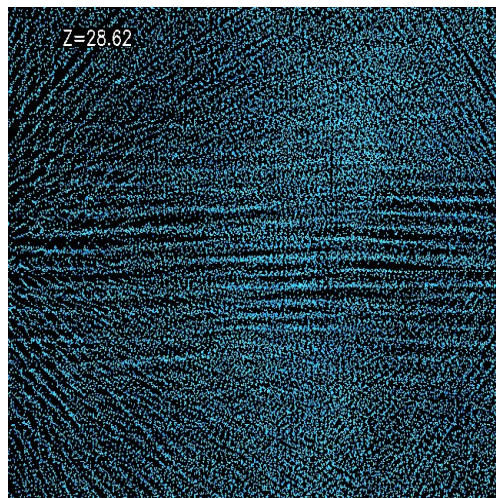


The CMB angular power spectrum from the Planck 2015 data (the blue dots with error bars) and the theoretical, best fit  $\Lambda$ CDM model with a power law primordial spectrum (the solid red curve)<sup>36</sup>.

<sup>36</sup> P. A. R. Ade *et al.*, [arXiv:1502.02114](https://arxiv.org/abs/1502.02114) [astro-ph.CO].



# Formation of structures due to gravitational instability

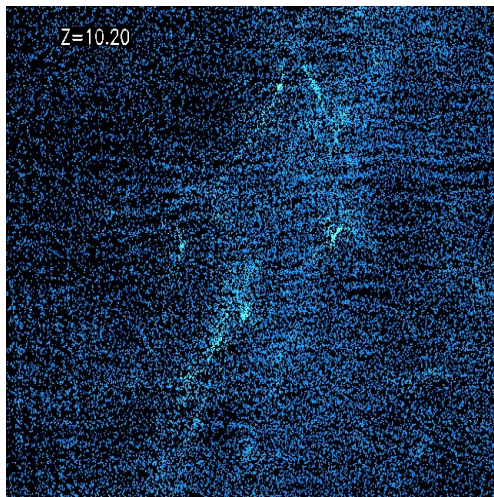


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability

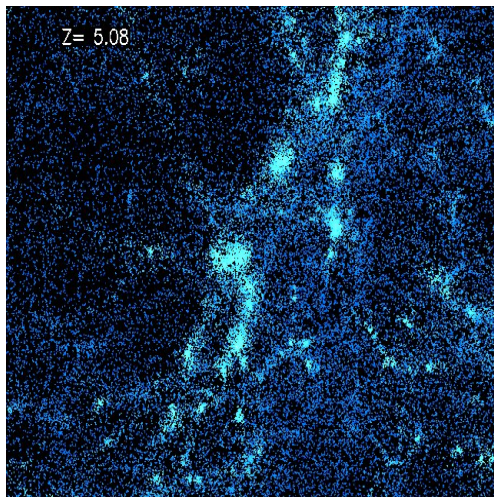


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup>Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability



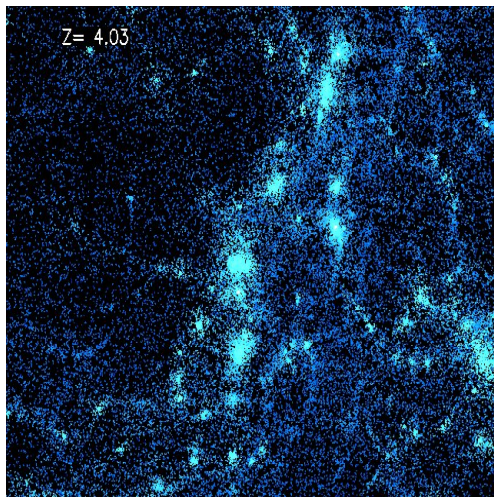
A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.





# Formation of structures due to gravitational instability

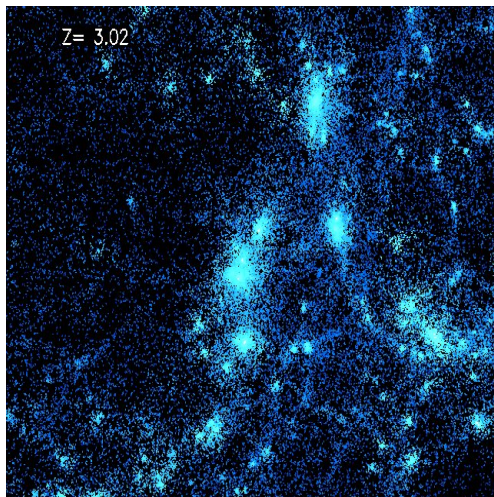


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup>Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability

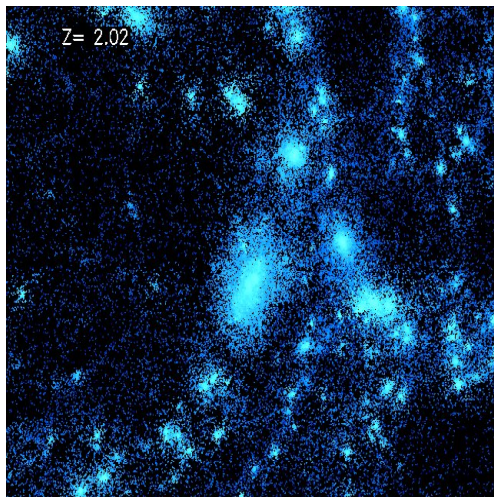


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability

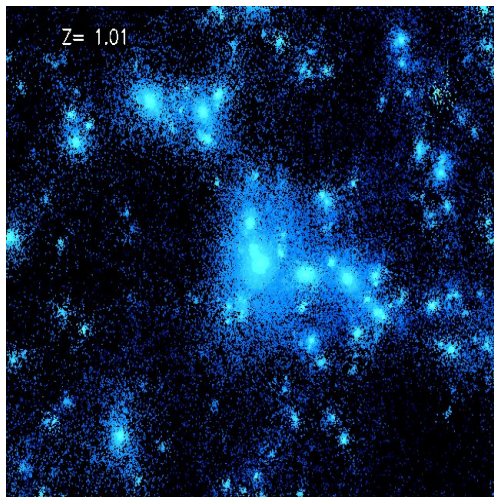


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability

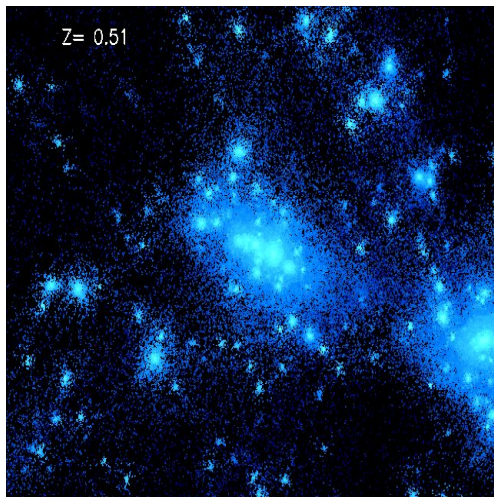


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability

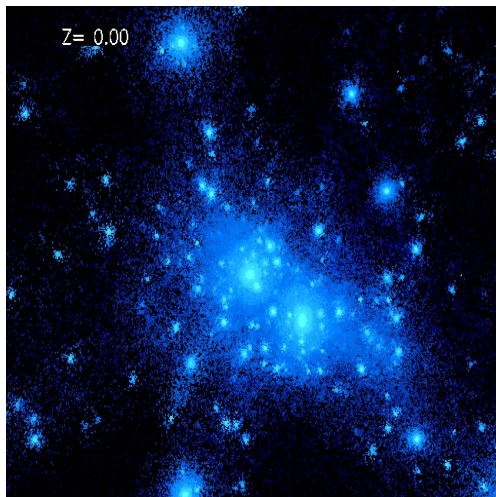


A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# Formation of structures due to gravitational instability



A numerical simulation illustrating the formation of large scale structures due to gravitational instability<sup>37</sup>.

<sup>37</sup> Images from <http://cfcp.uchicago.edu/lss/group.html>.



# The millennium simulation

- The Millennium Run used more than 10 billion particles to trace the evolution of the matter distribution in a cubic region of the universe over 2 billion light years on a side<sup>37</sup>.

---

<sup>37</sup> See <http://www.mpa-garching.mpg.de/galform/virgo/millennium/>.



# The millennium simulation

- The Millennium Run used more than 10 billion particles to trace the evolution of the matter distribution in a cubic region of the universe over 2 billion light years on a side<sup>37</sup>.
- It kept busy the principal supercomputer at the Max Planck Society's Supercomputing Centre in Garching, Germany for more than a month.

▶ Play movie

---

<sup>37</sup> See <http://www.mpa-garching.mpg.de/galform/virgo/millennium/>.





# Plan of the talk

- 1 A survey of the universe
- 2 The composition and evolution of the smooth universe
- 3 The origin and evolution of perturbations
- 4 **The standard model of cosmology**



# The standard model of cosmology

- The universe is homogeneous and isotropic at length scales of the order of 100 Mpc and above.



# The standard model of cosmology

- The universe is homogeneous and isotropic at length scales of the order of 100 Mpc and above.
- Baryons, *i.e.* matter as we know it, contribute less than 5% to the total density of the universe today. Most of the matter today is, in fact, dark and dominated by dark energy. Dark energy and pressureless (*i.e.* cold) dark matter contribute about 70% and 25% to the the density today, respectively.



# The standard model of cosmology

- The universe is homogeneous and isotropic at length scales of the order of **100 Mpc** and above.
- Baryons, *i.e.* matter as we know it, contribute less than **5%** to the total density of the universe today. Most of the matter today is, in fact, dark and dominated by dark energy. Dark energy and pressureless (*i.e.* cold) dark matter contribute about **70%** and **25%** to the the density today, respectively.
- The inflationary epoch magnifies the tiny fluctuations in the quantum fields present at the beginning of epoch into classical perturbations.



# The standard model of cosmology

- The universe is homogeneous and isotropic at length scales of the order of **100 Mpc** and above.
- Baryons, *i.e.* matter as we know it, contribute less than **5%** to the total density of the universe today. Most of the matter today is, in fact, dark and dominated by dark energy. Dark energy and pressureless (*i.e.* cold) dark matter contribute about **70%** and **25%** to the the density today, respectively.
- The inflationary epoch magnifies the tiny fluctuations in the quantum fields present at the beginning of epoch into classical perturbations.
- These inhomogeneities leave their imprints as anisotropies in the CMB.



# The standard model of cosmology

- The universe is homogeneous and isotropic at length scales of the order of **100 Mpc** and above.
- Baryons, *i.e.* matter as we know it, contribute less than **5%** to the total density of the universe today. Most of the matter today is, in fact, dark and dominated by dark energy. Dark energy and pressureless (*i.e.* cold) dark matter contribute about **70%** and **25%** to the the density today, respectively.
- The inflationary epoch magnifies the tiny fluctuations in the quantum fields present at the beginning of epoch into classical perturbations.
- These inhomogeneities leave their imprints as anisotropies in the CMB.
- Gravitational instability then takes over, and converts the tiny perturbations in the CMB into the large scale structures that we see around us today as galaxies and clusters of galaxies.



# Popular books

- 1 S. Weinberg, *The First Three Minutes* (Bantam, New York, 1977).
- 2 J. Silk, *The Big Bang* (W. H. Freeman, San Francisco, 1988).
- 3 S. Singh, *Big Bang* (Harper Collins, New York, 2004).



Thank you for your attention