

PH5170
QUANTUM MECHANICS II
January–May 2020

Lecture schedule and meeting hours

- The course will consist of about 42 lectures, including about 8–10 tutorial sessions. However, note that there will be no separate tutorial sessions, and they will be integrated with the lectures.
 - The duration of each lecture will be 50 minutes. We will be meeting in HSB 310.
 - The first lecture will be on Tuesday, January 14, and the last lecture will be on Thursday, April 23.
 - We will meet thrice a week. The lectures are scheduled for 11:00–11:50 AM on Tuesdays, 10:00–10:50 AM on Wednesdays, and 8:00–8:50 AM on Thursdays.
 - We may also meet during 5:00–5:50 PM on Fridays for either the quizzes or to make up for lectures that I may have to miss due to, say, travel. Changes in schedule, if any, will be notified sufficiently in advance.
 - If you would like to discuss with me about the course outside the lecture hours, you are welcome to meet me at my office (in HSB 202) during 12:00–1:00 PM on Tuesdays. In case you are unable to find me in my office on more than occasion, please send me an e-mail at sriram@physics.iitm.ac.in.
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Information about the course

- I will be distributing hard copies containing information such as the schedule of the lectures, the structure and the syllabus of the course, suitable textbooks and additional references at the start of the course. They will also be available on the course's page on Moodle at the following URL:
<https://courses.iitm.ac.in/>
 - The exercise sheets and other additional material will be made available on Moodle.
 - A PDF file containing these information as well as completed quizzes will also made be available at the link on this course at the following URL:
<http://www.physics.iitm.ac.in/~sriram/professional/teaching/teaching.html>
I will keep updating this file and the course's page on Moodle as we make progress.
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Quizzes, end-of-semester exam and grading

- The grading will be based on three scheduled quizzes and an end-of-semester exam.
 - I will consider the best two quizzes for grading, and the two will carry 25% weight each.
 - The three quizzes will be held on February 7, March 6 and April 3. The three dates are Fridays, and the quizzes will be held during 5:00–6:30 PM on these days.
 - The end-of-semester exam will be held during 1:00–4:00 PM on Friday, May 8, and the exam will carry 50% weight.
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Syllabus and structure

Quantum Mechanics II

1. Warming up with basic quantum mechanics [~ 6 lectures]

- (a) The Schrodinger equation
- (b) Uncertainty principle – Commutation relations – Simultaneous observables
- (c) Essential mathematical formalism
- (d) Simple problems in one dimension – Particle in a box – Harmonic oscillator
- (e) The Schrodinger equation in three dimensions – The angular and the radial equations

Exercise sheet 1

2. Rotation and angular momentum [~ 12 lectures]

- (a) Rotation – Orbital and spin angular momentum
- (b) Angular momentum algebra – Eigenstates and eigenvalues of angular momentum
- (c) Addition of angular momenta – Clebsch-Gordan coefficients
- (d) Measurements of spin correlations – Bell's inequality
- (e) Irreducible tensor operators – Wigner-Eckart theorem

Exercise sheets 2, 3 and 4

Quiz I

Additional exercises I

3. Systems of identical particles [~ 8 lectures + ~ 3 recorded lectures]

- (a) Symmetric and antisymmetric wave functions
- (b) Bosons and Fermions – Pauli's exclusion principle – The Slater determinant
- (c) Ground state of Helium
- (d) The free electron gas – Degeneracy pressure
- (e) Bloch's theorem – Band structure in solids
- (f) Quantum statistical mechanics – Black body radiation

Exercise sheets 5 and 6

Quiz II

4. Theory of scattering [~ 4 recorded lectures]

- (a) Scattering in non-relativistic quantum mechanics
- (b) Scattering amplitude and cross-section
- (c) Partial wave analysis – Phase shifts
- (d) The integral equation form of the Schrodinger equation – Born approximation
- (e) The optical theorem

Exercise sheets 7 and 8

Additional exercises II

5. **Relativistic quantum mechanics** [~ 4 recorded lectures]

- (a) Elements of relativistic quantum mechanics
- (b) The Klein-Gordon equation
- (c) The Dirac equation – Dirac matrices – Spinors
- (d) Positive and negative energy solutions – Physical interpretation
- (e) Non-relativistic limit of the Dirac equation
- (f) Covariant form of the Dirac equation – **Bilinear invariants**
- (g) **The hydrogen atom**

Exercise sheets 9 and 10

Assignment in lieu of Quiz III

End-of-semester exam

Advanced problems

Note: The topics/plans in red could not be covered/implemented for want of time.

Basic textbooks

1. J. Bjorken and S. Drell, *Relativistic Quantum Mechanics* (McGraw-Hill, New York, 1965).
2. L. D. Landau and E. M. Lifshitz, *Quantum Mechanics* (Course of Theoretical Physics, Volume 3), Third Edition (Pergamon Press, New York, 1977).
3. D. J. Griffiths, *Introduction to Elementary Particles* (John Wiley, New York, 1987).
4. J. J. Sakurai, *Modern Quantum Mechanics* (Addison-Wesley, Singapore, 1994).
5. D. J. Griffiths, *Introduction to Quantum Mechanics*, Second Edition (Pearson Education, Delhi, 2005).

Additional references

1. P. A. M. Dirac, *The Principles of Quantum Mechanics*, Fourth Edition (Oxford University Press, Oxford, 1958).
 2. A. Messiah, *Quantum Mechanics*, Volumes 1 and 2 (North Holland, Amsterdam, 1961).
 3. J. J. Sakurai, *Advanced Quantum Mechanics* (Addison-Wesley, Singapore, 1967).
 4. F. Halzen and A. D. Martin, *Quarks and Leptons: An Introductory Course in Modern Particle Physics* (John Wiley, New York, 1984).
 5. R. W. Robinett, *Quantum Mechanics*, Second Edition (Oxford University Press, Oxford, 2006).
 6. F. Dyson, *Advanced Quantum Mechanics* (World Scientific, Singapore, 2007).
 7. R. Shankar, *Principles of Quantum Mechanics*, Second Edition (Springer, Delhi, 2008).
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Exercise sheet 1

Warming up with basic quantum mechanics

Some classical mechanics

1. Non-relativistic particle in an electromagnetic field: A non-relativistic particle that is moving in an electromagnetic field described by the scalar potential ϕ and the vector potential \mathbf{A} is governed by the Lagrangian

$$L = \frac{m \mathbf{v}^2}{2} + q \left(\frac{\mathbf{v}}{c} \cdot \mathbf{A} \right) - q \phi,$$

where m and q are the mass and the charge of the particle, while c denotes the velocity of light. Show that the equation of motion of the particle is given by

$$m \frac{d\mathbf{v}}{dt} = q \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right),$$

where \mathbf{E} and the \mathbf{B} are the electric and the magnetic fields given by

$$\mathbf{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} \quad \text{and} \quad \mathbf{B} = \nabla \times \mathbf{A}.$$

Note: The scalar and the vector potentials, viz. ϕ and \mathbf{A} , are dependent on time *as well as* space. Further, given two vectors, say, \mathbf{C} and \mathbf{D} , one can write,

$$\nabla(\mathbf{C} \cdot \mathbf{D}) = (\mathbf{D} \cdot \nabla) \mathbf{C} + (\mathbf{C} \cdot \nabla) \mathbf{D} + \mathbf{D} \times (\nabla \times \mathbf{C}) + \mathbf{C} \times (\nabla \times \mathbf{D}).$$

Also, since \mathbf{A} depends on time as well as space, we have,

$$\frac{d\mathbf{A}}{dt} = \frac{\partial \mathbf{A}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{A}.$$

2. Period associated with bounded, one-dimensional motion: Determine the period of oscillation as a function of the energy, say, E , when a particle of mass m moves in a field governed by the potential $V(x) = V_0 |x|^n$, where V_0 is a constant and n is a positive integer.
3. Poisson brackets: Recall that the Poisson bracket $\{A, B\}$ between two observables $A(q_i, p_i)$ and $B(q_i, p_i)$ is defined as

$$\{A, B\} = \sum_{i=1}^N \left(\frac{\partial A}{\partial q_i} \frac{\partial B}{\partial p_i} - \frac{\partial A}{\partial p_i} \frac{\partial B}{\partial q_i} \right),$$

where q_i and p_i denote the generalized coordinates and the corresponding conjugate momenta, respectively, while N denotes the number of degrees of freedom of the system.

- (a) Poisson brackets for position and momenta: Establish the following relations: $\{q_i, q_j\} = 0$, $\{p_i, p_j\} = 0$, $\{q_i, p_j\} = \delta_{ij}$, $\dot{q}_i = \{q_i, H\}$ and $\dot{p}_i = \{p_i, H\}$, where H represents the Hamiltonian of the system, while δ_{ij} represents the standard Kronecker symbol.
- (b) Poisson brackets for angular momenta: Let (L_x, L_y, L_z) be the Cartesian components of the angular momentum vector \mathbf{L} , and let $L^2 = L_x^2 + L_y^2 + L_z^2$. Evaluate the following Poisson brackets: $\{L_x, L_y\}$, $\{L_y, L_z\}$ and $\{L_z, L_x\}$, $\{L_x, L^2\}$, $\{L_y, L^2\}$ and $\{L_z, L^2\}$.
Hint: For the final set of Poisson brackets, it may be useful to establish and use the result $\{A, BC\} = \{A, B\}C + B\{A, C\}$.

4. Phase portraits: Draw the phase portraits of a particle moving in the following one dimensional potentials: (a) $U(x) = \alpha |x|^n$, (b) $U(x) = \alpha x^2 - \beta x^3$, (c) $U(x) = \alpha (x^2 - \beta^2)^2$, and (d) $U(\theta) = -\alpha \cos \theta$, where $(\alpha, \beta) > 0$ and $n > 2$.

Some basics of quantum mechanics

5. Operators, expectation values and properties: Recall that the expectation value of an operator \hat{A} is defined as

$$\langle \hat{A} \rangle = \int dx \Psi^* \hat{A} \Psi.$$

- (a) Hermitian operators: An operator \hat{A} is said to be hermitian if

$$\langle \hat{A} \rangle = \langle \hat{A} \rangle^*.$$

Show that the position, the momentum and the Hamiltonian operators are hermitian.

- (b) Motivating the momentum operator: Using the time-dependent Schrodinger equation, show that

$$\frac{d\langle x \rangle}{dt} = -\frac{i\hbar}{m} \int dx \Psi^* \frac{\partial \Psi}{\partial x} = \langle \hat{p}_x \rangle,$$

a relation which can be said to motivate the expression for the momentum operator, viz. that $\hat{p}_x = -i\hbar \partial/\partial x$.

- (c) Ehrenfest's theorem: Show that

$$\frac{d\langle \hat{p}_x \rangle}{dt} = -\left\langle \frac{\partial V}{\partial x} \right\rangle,$$

a relation that is often referred to as the Ehrenfest's theorem.

- (d) Virial theorem: Let x and p_x denote the position and momentum of a particle moving in a given time-independent potential $V(x)$. Show that the system satisfies the equation

$$\frac{\langle \hat{p}_x^2 \rangle}{2m} = \frac{1}{2} \left\langle \hat{x} \frac{\partial V}{\partial x} \right\rangle,$$

where the expectation values are evaluated in a *stationary* state described by a *real* wavefunction.

Note: The above relation is often referred to as the virial theorem.

Hint: Establish the virial theorem in the following two steps. To begin with, show that

$$\int_{-\infty}^{\infty} dx \psi x \frac{dV}{dx} \psi = -\langle \hat{V} \rangle - 2 \int_{-\infty}^{\infty} dx \frac{d\psi}{dx} x V \psi$$

and then, using the time-independent Schrodinger equation, illustrate that

$$-2 \int_{-\infty}^{\infty} dx \frac{d\psi}{dx} x V \psi = E + \frac{\hbar^2}{2m} \int_{-\infty}^{\infty} dx \left(\frac{d\psi}{dx} \right)^2.$$

6. Spreading of wave packets: A free particle has the initial wave function

$$\Psi(x, 0) = A e^{-a x^2},$$

where A and a are constants, with a being real and positive.

- (a) Normalize $\Psi(x, 0)$.
 (b) Find $\Psi(x, t)$.
 (c) Plot $\Psi(x, t)$ at $t = 0$ and for large t . Determine qualitatively what happens as time goes on?
 (d) Find $\langle \hat{x} \rangle$, $\langle \hat{x}^2 \rangle$, $\langle \hat{p} \rangle$, $\langle \hat{p}^2 \rangle$, Δx and Δp .

- (e) Does the uncertainty principle hold? At what time does the system have the minimum uncertainty?

Particle in a box

7. Particle in a box: Consider a particle of mass m that confined is to a box with its walls at $x = 0$ and $x = a$.

- (a) Particle in a box I: At time $t = 0$, the particle is equally likely to be found over the domain $0 < x < a/2$.
- What is the initial wave function of the system, i.e. $\Psi(x, 0)$?
Hint: Assume that the wave function is real, and make sure you normalize the wave function.
 - What is the probability that a measurement of the energy of the particle would yield the value $E_1 = \pi^2 \hbar^2 / (2 m a^2)$?
- (b) Particle in a box II: At time $t = 0$, the wave function of the particle is given by

$$\Psi(x, 0) = A [\psi_1(x) + \psi_2(x)],$$

where ψ_1 and ψ_2 are the ground state and the first excited state of the system.

- Determine the constant A assuming that $\Psi(x, 0)$ is normalized.
 - What is $\Psi(x, t)$ for $t > 0$?
 - Determine the expectation values $\langle \hat{x} \rangle$ and $\langle \hat{p}_x \rangle$ in the state $\Psi(x, t)$. Express the results in terms of the angular frequency $\omega = \pi^2 \hbar / (2 m a^2)$.
 - What are probabilities of the system to be found in the ground and the first excited states? What is the expectation value of the Hamiltonian, i.e. $\langle \hat{H} \rangle$, in the state $\Psi(x, t)$?
- (c) Particle in a box III: The particle is described by the following wave function:

$$\psi(x) = \begin{cases} A (x/a) & \text{for } 0 < x < a/2, \\ A [1 - (x/a)] & \text{for } a/2 < x < a, \end{cases}$$

where A is a real constant. If the energy of the system is measured, what is the probability for finding the energy eigen value to be $E_n = n^2 \pi^2 \hbar^2 / (2 m a^2)$?

8. Quantum revival: Consider an arbitrary wavefunction describing a particle in the infinite square well.
- Show that the wave function will return to its original form after a time $T_Q = 4 m a^2 / (\pi \hbar)$.
Note: The time T_Q is known as the *quantum revival time*.
 - Determine the classical revival time T_C for a particle of energy E bouncing back and forth between the walls.
 - What is the energy for which $T_Q = T_C$?

The simple harmonic oscillator

9. Oscillating charge in an electric field: Let a particle of mass m and charge q be oscillating in the simple harmonic potential $V(x) = m \omega^2 x^2 / 2$. A constant electric field of strength \mathcal{E} is turned on along the positive x -direction.
- What is the complete potential influencing the charge in the presence of the electric field?
 - Draw the classical trajectory of the charge in phase space when the electric field has been turned on. How does it compare with the original trajectory?
 - What are the energy eigen values of the system when it is quantized?

(d) Plot the ground state wave function of the system with and without the electric field.

10. Half-an-oscillator: Determine the energy levels and the corresponding eigen functions of an oscillator which is subjected to the additional condition that the potential is infinite for $x \leq 0$.

Note: You do not have to separately solve the Schrodinger equation. You can easily identify the allowed eigen functions and eigen values from the solutions of the original, complete, oscillator!

11. Wagging the dog: Recall that the time-independent Schrodinger equation satisfied by a simple harmonic oscillator of mass m and frequency ω is given by

$$-\frac{\hbar^2}{2m} \frac{d^2\psi_E}{dx^2} + \frac{1}{2} m \omega^2 x^2 \psi_E = E \psi_E.$$

In terms of the dimensionless variable

$$\xi = \sqrt{\frac{m\omega}{\hbar}} x,$$

the above time-independent Schrodinger equation reduces to

$$\frac{d^2\psi_E}{d\xi^2} + (\mathcal{E} - \xi^2) \psi_E = 0,$$

where \mathcal{E} is the energy expressed in units of $(\hbar\omega/2)$, and is given by

$$\mathcal{E} = \frac{2E}{\hbar\omega}.$$

According to the ‘wag-the-dog’ method, one solves the above differential equation numerically, say, using **Mathematica**, varying \mathcal{E} until a wave function that goes to zero at large ξ is obtained.

Find the ground state energy and the energies of the first two excited states of the harmonic oscillator to five significant digits by the ‘wag-the-dog’ method.

12. Constructing the energy eigen states of the harmonic oscillator: If $\psi_n(x)$ denotes the n -th energy eigen state of the harmonic oscillator and \hat{a} represents the lowering operator, then, recall that

$$\hat{a} \psi_n(x) = \sqrt{n} \psi_{n-1}(x).$$

(a) Using the fact that $\hat{a} \psi_0(x) = 0$, construct the ground state wave function $\psi_0(x)$, ensuring that it is normalized.

(b) Also, recall that

$$\hat{a}^\dagger \psi_n(x) = \sqrt{n+1} \psi_{n+1}(x),$$

where \hat{a}^\dagger is the raising operator. Operate the raising operator on $\psi_0(x)$ successively to construct the wavefunctions describing the first and the second excited states, i.e. $\psi_1(x)$ and $\psi_2(x)$.

13. More about oscillators: Let $|0\rangle$ represent the ground state of a one dimensional quantum oscillator. Show that

$$\langle 0 | e^{ik\hat{x}} | 0 \rangle = \exp - (k^2 \langle 0 | \hat{x}^2 | 0 \rangle / 2),$$

where \hat{x} is the position operator.

14. Coherent states of the harmonic oscillator: Consider states, say, $|\alpha\rangle$, which are eigen states of the annihilation (or, more precisely, the lowering) operator, i.e.

$$\hat{a} |\alpha\rangle = \alpha |\alpha\rangle,$$

where α is a complex number.

Note: The state $|\alpha\rangle$ is called the coherent state.

- (a) Calculate the quantities $\langle \hat{x} \rangle$, $\langle \hat{x}^2 \rangle$, $\langle \hat{p}_x \rangle$ and $\langle \hat{p}_x^2 \rangle$ in the coherent state.
- (b) Also, evaluate the quantities Δx and Δp_x in the state, and show that $\Delta x \Delta p_x = \hbar/2$.
- (c) Like any other general state, the coherent state can be expanded in terms of the energy eigenstates $|n\rangle$ of the harmonic oscillator as follows:

$$|\alpha\rangle = \sum_{n=0}^{\infty} c_n |n\rangle.$$

Show that the quantities c_n are given by

$$c_n = \frac{\alpha^n}{\sqrt{n!}} c_0.$$

- (d) Determine c_0 by normalizing $|\alpha\rangle$.
- (e) Upon including the time dependence, show that the coherent state continues to be an eigenstate of the lowering operator \hat{a} with the eigen value evolving in time as

$$\alpha(t) = \alpha e^{-i\omega t}.$$

Note: Therefore, a coherent state *remains* coherent, and continues to minimize the uncertainty.

- (f) Is the ground state $|0\rangle$ itself a coherent state? If so, what is the eigen value α ?

Some mathematical and conceptual aspects

15. Probabilities in momentum space: A particle of mass m is bound in the delta function well $V(x) = -a\delta(x)$, where $a > 0$. What is the probability that a measurement of the particle's momentum would yield a value greater than $p_0 = ma/\hbar$?
16. The energy-time uncertainty principle: Consider a system that is described by the Hamiltonian operator \hat{H} .

- (a) Given an operator, say, \hat{Q} , establish the following relation:

$$\frac{d\langle \hat{Q} \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{H}, \hat{Q}] \rangle + \left\langle \frac{\partial \hat{Q}}{\partial t} \right\rangle,$$

where the expectation values are evaluated in a specific state.

- (b) When \hat{Q} does not explicitly depend on time, using the generalized uncertainty principle, show that

$$\Delta H \Delta Q \geq \frac{\hbar}{2} \left| \frac{d\langle \hat{Q} \rangle}{dt} \right|.$$

- (c) Defining

$$\Delta t \equiv \frac{\Delta Q}{|d\langle \hat{Q} \rangle/dt|},$$

establish that

$$\Delta E \Delta t \geq \frac{\hbar}{2},$$

and interpret this result.

17. Two-dimensional Hilbert space: Imagine a system in which there are only two linearly independent states, viz.

$$|1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{and} \quad |2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

The most general state would then be a normalized linear combination, i.e.

$$|\psi\rangle = \alpha |1\rangle + \beta |2\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix},$$

with $|\alpha|^2 + |\beta|^2 = 1$. The Hamiltonian of the system can, evidently, be expressed as a 2×2 hermitian matrix. Suppose it has the following form:

$$H = \begin{pmatrix} a & b \\ b & a \end{pmatrix},$$

where a and b are *real* constants. If the system starts in the state $|1\rangle$ at an initial time, say, $t = 0$, determine the state of the system at a later time t .

18. A two level system: The Hamiltonian operator of a certain two level system is given by

$$\hat{H} = E \left(|1\rangle\langle 1| - |2\rangle\langle 2| + |1\rangle\langle 2| + |2\rangle\langle 1| \right),$$

where $|1\rangle$ and $|2\rangle$ form an orthonormal basis, while E is a number with the dimensions of energy.

- Find the eigen values and the normalized eigen vectors, i.e. as a linear combination of the basis vectors $|1\rangle$ and $|2\rangle$, of the above Hamiltonian operator.
- What is the matrix that represents the operator \hat{H} in this basis?

19. A three level system: The Hamiltonian for a three level system is represented by the matrix

$$H = \hbar\omega \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 2 \end{pmatrix}.$$

Two other observables, say, A and B , are represented by the matrices

$$A = \lambda \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 2 \end{pmatrix} \quad \text{and} \quad B = \mu \begin{pmatrix} 2 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix},$$

where ω , λ and μ are positive real numbers.

- Find the eigen values and normalized eigen vectors of H , A , and B .
- Suppose the system starts in the generic state

$$|\psi(t=0)\rangle = \begin{pmatrix} c_1 \\ c_2 \\ c_3 \end{pmatrix},$$

with $|c_1|^2 + |c_2|^2 + |c_3|^2 = 1$. Find the expectation values of H , A and B in the state at $t = 0$.

- What is $|\psi(t)\rangle$ for $t > 0$? If you measure the energy of the state at a time t , what are the values of energies that you will get and what would be the probability for obtaining each of the values?
 - Also, arrive at the corresponding answers for the quantities A and B .
20. Expectation values in momentum space: Given a wave function, say, $\Psi(x, t)$, in the position space, the corresponding wave function in momentum space is given by

$$\Phi(p, t) = \int_{-\infty}^{\infty} \frac{dp}{\sqrt{2\pi\hbar}} \Psi(x, t) e^{-ipx/\hbar}.$$

- (a) Show that, if the wavefunction $\Psi(x, t)$ is normalized in position space, it is normalized in momentum space as well, i.e.

$$\int_{-\infty}^{\infty} dx |\Psi(x, t)|^2 = \int_{-\infty}^{\infty} dp |\Phi(p, t)|^2 = 1.$$

- (b) Recall that, in the position representation, the expectation value of the momentum operator can be expressed as follows:

$$\langle \hat{p} \rangle = -i\hbar \int_{-\infty}^{\infty} dx \Psi^*(x, t) \frac{\partial \Psi(x, t)}{\partial x}.$$

Show that, in momentum space, it can be expressed as

$$\langle \hat{p} \rangle = \int_{-\infty}^{\infty} dp p |\Phi(p, t)|^2.$$

- (c) Similarly, show that the following expectation value of the position operator in position space

$$\langle \hat{x} \rangle = \int_{-\infty}^{\infty} dx x |\Psi(x, t)|^2,$$

can be written as

$$\langle \hat{x} \rangle = i\hbar \int_{-\infty}^{\infty} dp \Phi^*(p, t) \frac{\partial \Phi(p, t)}{\partial p}.$$

21. The Schrodinger equation in momentum space: Assuming that the potential $V(x)$ can be expanded in a Taylor series, show that the time-dependent Schrodinger equation in momentum space can be written as

$$\frac{p^2}{2m} \Phi(p, t) + V \left(i\hbar \frac{\partial}{\partial p} \right) \Phi(p, t) = i\hbar \frac{\partial \Phi(p, t)}{\partial t}.$$

22. Sequential measurements: An operator \hat{A} , representing the observable A , has two normalized eigenstates ψ_1 and ψ_2 , with eigen values a_1 and a_2 . Operator \hat{B} , representing another observable B , has two normalized eigenstates ϕ_1 and ϕ_2 , with eigen values b_1 and b_2 . These eigen states are related as follows:

$$\psi_1 = \frac{3}{5} \phi_1 + \frac{4}{5} \phi_2 \quad \text{and} \quad \psi_2 = \frac{4}{5} \phi_1 - \frac{3}{5} \phi_2.$$

- (a) Observable A is measured and the value a_1 is obtained. What is the state of the system immediately after this measurement?
 (b) If B is now measured, what are the possible results, and what are their probabilities?
 (c) Immediately after the measurement of B , A is measured again. What is the probability of getting a_1 ?

The Schrodinger equation in three dimensions

23. Particle in a three dimensional box: Consider a particle that is confined to a three dimensional box of side, say, a . In other words, the particle is free inside the box, but the potential energy is infinite on the walls of the box, thereby confining the particle to the box.

- (a) Determine the energy eigen functions and the corresponding energy eigen values.
 (b) Does there exist degenerate energy eigen states? Identify a few of them.

24. Particle in a spherical well: Consider a particle that is confined to the following spherical well:

$$V(r) = \begin{cases} 0 & \text{for } r < a, \\ \infty & \text{for } r \geq a. \end{cases}$$

Find the energy eigen functions and the corresponding energy eigen values of the particle.

25. Expectation values in the energy eigen states of the hydrogen atom: Recall that, the normalized wavefunctions that describe the energy eigen states of the electron in the hydrogen atom are given by

$$\psi_{nlm}(r, \theta, \phi) = \left[\left(\frac{2}{na_0} \right)^3 \frac{(n-l-1)!}{2n[(n+l)!]^3} \right]^{1/2} e^{-r/(na_0)} \left(\frac{2r}{na_0} \right)^l L_{n-l-1}^{2l+1}(2r/na_0) Y_l^m(\theta, \phi),$$

where $L_p^q(x)$ and Y_l^m represent the associated Laguerre polynomials and the spherical harmonics, respectively, while a_0 denotes the Bohr radius.

- Evaluate $\langle \hat{r} \rangle$ and $\langle \hat{r}^2 \rangle$ for the electron in the ground state of the hydrogen atom, and express it in terms of the Bohr radius.
- Find $\langle \hat{x} \rangle$ and $\langle \hat{x}^2 \rangle$ for the electron in the ground state of hydrogen.
Hint: Express r^2 as $x^2 + y^2 + z^2$ and exploit the symmetry of the ground state.
- Calculate $\langle \hat{x}^2 \rangle$ in the state $n = 2$, $l = 1$ and $m = 1$.

Note: This state is not symmetrical in x , y and z . Use $x = r \sin \theta \cos \phi$.

Exercise sheet 2
Spin

1. Velocity on the surface of a spinning electron: Consider the electron to be a classical solid sphere. Assume that the radius of the electron is given by the classical electron radius, viz.

$$r_c = \frac{e^2}{4\pi\epsilon_0 m_e c^2}$$

where e and m_e denote the charge and the mass of the electron, while c represents the speed of light. Also, assume that the angular momentum of the electron is $\hbar/2$. Evaluate the speed on the surface of the electron under these conditions.

2. Orienting spin along an arbitrary direction: Let $\hat{\mathbf{n}}$ be a three-dimensional unit vector whose polar angle (i.e. the angle with respect to the z -axis) is θ and the azimuthal angle (i.e. the angle with respect to the x -axis, when the unit vector $\hat{\mathbf{n}}$ has been projected on to the x - y plane) is ϕ , i.e.

$$\hat{\mathbf{n}} = \sin\theta \cos\phi \hat{i} + \sin\theta \sin\phi \hat{j} + \cos\theta \hat{k}.$$

- (a) Obtain the eigen values of the operator $\hat{\mathbf{S}} \cdot \hat{\mathbf{n}}$ describing a spin- $\frac{1}{2}$ particle.
 (b) Construct the corresponding eigen vectors.
 (c) Show that the eigen vectors are orthonormal.
3. Larmor precession: Consider a charged, spin- $\frac{1}{2}$ particle which is at rest in an external and uniform magnetic field, say, \mathbf{B} , that is oriented along the z -direction, i.e. $\mathbf{B} = B\hat{k}$, where B is a constant. The Hamiltonian of the particle is then given by

$$\hat{H} = -\gamma B \hat{S}_z,$$

where γ is known as the gyromagnetic ratio of the particle.

- (a) Determine the most general, time dependent, solution that describes the state of the particle.
 (b) Evaluate the expectation values of the operators \hat{S}_x , \hat{S}_y and \hat{S}_z in the state.
 (c) Show that the expectation value of the operator $\hat{\mathbf{S}} = \hat{S}_x \hat{i} + \hat{S}_y \hat{j} + \hat{S}_z \hat{k}$ is tilted at a constant angle with respect to the direction of the magnetic field and precesses about the field at the so-called Larmor frequency $\omega = \gamma B$.
4. Mean values and uncertainties associated with spin operators: An electron is in the spin state

$$\chi = A \begin{pmatrix} 3i \\ 4 \end{pmatrix}.$$

- (a) Determine the normalization constant A .
 (b) Find the expectation values of the operators \hat{S}_x , \hat{S}_y and \hat{S}_z in the above state.
 (c) Evaluate the corresponding uncertainties, i.e. ΔS_x , ΔS_y and ΔS_z .
 (d) Examine if the products of any two of these quantities are consistent with the corresponding uncertainty principles.
5. Spin- $\frac{3}{2}$ particle: Consider a particle with spin $\frac{3}{2}$.

- (a) What are the eigen values and eigen states of the \hat{S}_z operator for the system?
 (b) Determine the effects of the operators \hat{S}_+ and \hat{S}_- on these eigen states.
 (c) Construct the matrices describing the operators \hat{S}_+ and \hat{S}_- .
 (d) Obtain the matrix describing the operator \hat{S}_x .
-

Quiz I

From basic quantum mechanics to spin

1. Absence of degenerate bound states in one spatial dimension: Two or more quantum states are said to be degenerate if they are described by distinct solutions to the time-independent Schrodinger equation corresponding to the *same* energy. For example, the free particle states are doubly degenerate—one solution describes motion to the right and the other motion to the left. It is not an accident that we have not encountered normalizable degenerate solutions in one spatial dimension. By following the steps listed below, prove that there are no degenerate bound states in one spatial dimension.

- (a) Suppose that there are two solutions, say, ψ_1 and ψ_2 , with the same energy E . Using the time-independent Schrodinger equation, show that the quantity

$$\mathcal{W} = \psi_1 \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_1}{dx}$$

is a constant independent of x .

3 marks

- (b) Argue that, since the wavefunctions ψ_1 and ψ_2 are normalizable, the quantity \mathcal{W} defined above must vanish.

3 marks

- (c) If $\mathcal{W} = 0$, integrate the above equation to show that ψ_2 is a multiple of ψ_1 and hence the solutions are not distinct.

4 marks

2. The translation operator: Recall that a wavefunction can be translated by the vector \mathbf{a} in three dimensions using the operator

$$\hat{\mathcal{T}}(\mathbf{a}) = \exp - \left(\frac{i \hat{\mathbf{p}} \cdot \mathbf{a}}{\hbar} \right),$$

where $\hat{\mathbf{p}}$ denotes the momentum operator. Let $|\psi\rangle$ be a state vector describing a system and, upon translation, let

$$|\psi\rangle \rightarrow |\bar{\psi}\rangle = \hat{\mathcal{T}}(\mathbf{a})|\psi\rangle.$$

- (a) Evaluate the commutation relation: $[\hat{x}_i, \hat{\mathcal{T}}(\mathbf{a})]$.

7 marks

- (b) Using the result of the commutation relation, evaluate $\langle \bar{\psi} | \hat{\mathbf{x}} | \bar{\psi} \rangle$ and express it in terms of $\langle \psi | \hat{\mathbf{x}} | \psi \rangle$.

3 marks

3. Commuting operators and simultaneous eigenkets: Consider a three-dimensional ket space. If a certain set of orthonormal kets, say, $|1\rangle$, $|2\rangle$ and $|3\rangle$, are used as the base kets, the operators \hat{A} and \hat{B} are found to be represented by the matrices

$$A = \begin{pmatrix} a & 0 & 0 \\ 0 & -a & 0 \\ 0 & 0 & -a \end{pmatrix}, \quad B = \begin{pmatrix} b & 0 & 0 \\ 0 & 0 & -ib \\ 0 & ib & 0 \end{pmatrix}$$

with a and b both real.

- (a) Do the operators \hat{A} and \hat{B} represent observables?

1 marks

- (b) If they do, can the observables be measured simultaneously?

3 marks

- (c) What is the spectrum of the two operators? Are they degenerate?

2 marks

- (d) If the \hat{A} and \hat{B} represent observables and, if they can be measured simultaneously, construct a new set of orthonormal kets which are simultaneous eigenkets of both the operators. What are the eigen values of \hat{A} and \hat{B} for these eigenkets?

3+1 marks

4. Wave function describing the coherent state: Earlier, that we had arrived at the wave function $\psi_0(x) = \langle x|0\rangle$ describing the ground state of the oscillator using the following definition of the ground state $|0\rangle$:

$$\hat{a}|0\rangle = 0$$

and the position representation of the lowering operator \hat{a} . Recall that the coherent state $|\alpha\rangle$ is defined through the relation

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle,$$

where α is a complex number.

- (a) Evaluate $\langle \hat{x} \rangle$ and $\langle \hat{p}_x \rangle$ in the state $|\alpha\rangle$. 3 marks
- (b) Using the above definition of $|\alpha\rangle$, obtain the wave function $\psi_\alpha(x) = \langle x|\alpha\rangle$. Express the wavefunction $\psi_\alpha(x)$ in terms of $\langle \hat{x} \rangle$ and $\langle \hat{p}_x \rangle$. 6+1 marks
5. Maximizing uncertainty: Recall that, according to the generalized uncertainty principle, the uncertainty associated with two observables represented by the operators \hat{A} and \hat{B} is given by

$$\Delta A^2 \Delta B^2 \geq \frac{1}{4} |[\hat{A}, \hat{B}]|^2,$$

where the uncertainty, say, ΔA^2 , is defined as $\Delta A^2 = \langle \hat{A}^2 \rangle - \langle \hat{A} \rangle^2$ and the angular brackets as usual represent expectation values evaluated in a given state. Let, as usual, $|+\rangle$ and $|-\rangle$ represent the eigenkets of the \hat{S}_z operator describing a spin- $\frac{1}{2}$ system.

- (a) Find the linear combination of the kets $|+\rangle$ and $|-\rangle$ that maximizes the following uncertainty product: 7 marks

$$\Delta S_x^2 \Delta S_y^2.$$

Note: It will be convenient to express the general state describing the spin- $\frac{1}{2}$ system as follows:

$$|\chi\rangle = \cos(\theta/2)|+\rangle + \sin(\theta/2)e^{i\phi}|-\rangle.$$

- (b) Verify explicitly that, for the state $|\chi\rangle$ you have found, the uncertainty relation for the operators \hat{S}_x and \hat{S}_y is not violated. 3 marks

Exercise sheet 3

Orbital angular momentum

1. Operators describing orbital angular momentum: Show that

$$(a) \hat{L}_x = -i\hbar \left(-\sin\phi \frac{\partial}{\partial\theta} - \cot\theta \cos\phi \frac{\partial}{\partial\phi} \right),$$

$$(b) \hat{L}_y = -i\hbar \left(\cos\phi \frac{\partial}{\partial\theta} - \cot\theta \sin\phi \frac{\partial}{\partial\phi} \right),$$

$$(c) \hat{L}_{\pm} = \pm\hbar e^{\pm i\phi} \left(\frac{\partial}{\partial\theta} \pm i \cot\theta \frac{\partial}{\partial\phi} \right),$$

$$(d) \hat{L}_z = -i\hbar \frac{\partial}{\partial\phi},$$

$$(e) \hat{L}^2 = -\hbar^2 \left[\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2}{\partial\phi^2} \right].$$

2. Ehrenfest's theorem for angular momentum: Consider a particle moving in a potential $V(\mathbf{r})$.

(a) If $\hat{\mathbf{L}}$ and $\hat{\mathbf{N}}$ are the operators representing angular momentum and torque, show that

$$\frac{d\langle\hat{\mathbf{L}}\rangle}{dt} = \langle\hat{\mathbf{N}}\rangle,$$

where $\mathbf{N} = \mathbf{r} \times (-\nabla V)$.

(b) Show that $d\langle\hat{\mathbf{L}}\rangle/dt = 0$ for any spherically symmetric potential, i.e. when $V(\mathbf{r}) = V(r)$.

3. Matrices representing the angular momentum operators: Earlier, we had discussed the matrices describing the operators \hat{S}_x , \hat{S}_y and \hat{S}_z of a spin- $\frac{1}{2}$ system. Consider the matrices describing the angular momentum operators \hat{L}_x , \hat{L}_y and \hat{L}_z corresponding to $l = 1$.

(a) List the allowed states $|l, m\rangle$ for $l = 1$. What is the rank of the matrices in such a Hilbert space?

(b) Construct the matrices representing \hat{L}_z and \hat{L}^2 .

(c) Construct the matrices describing \hat{L}_+ and \hat{L}_- , and thereby \hat{L}_x and \hat{L}_y .

Hint: Recall the effects of the operators \hat{L}_z , \hat{L}^2 , \hat{L}_+ and \hat{L}_- on the states $|l, m\rangle$.

4. Spherical harmonics: Recall that, $\langle\theta, \phi|l, m\rangle = Y_l^m(\theta, \phi)$. Using the relations $\hat{L}_+|l, l\rangle = 0$ and $\hat{L}_z|l, l\rangle = l\hbar|l, l\rangle$, determine the functional form of $Y_l^l(\theta, \phi)$.

Hint: Using the representations of \hat{L}_+ and \hat{L}_z as differential operators, write the two relations above as differential equations and solve them.

Note: You need not have to normalize the wavefunction.

5. Particle in a central potential: The wavefunction of a particle in a central potential $V(r)$ is given by

$$\psi(\mathbf{x}) = (x + y + 3z) f(r).$$

(a) Is this wavefunction an eigenfunction of \hat{L}^2 ? If so, what is the value of l ? If not, what are the possible values of l we may obtain if \hat{L}^2 is measured?

(b) What are the probabilities for the particle to be found in the various m_l states?

(c) Suppose it is known somehow that the above wavefunction is an energy eigenfunction with eigenvalue E . How does one find the potential $V(r)$?

Exercise sheet 4

Addition of angular momentum

1. Commutation relations involving angular momentum: Given $\hat{\mathbf{J}} = \hat{\mathbf{J}}_1 + \hat{\mathbf{J}}_2$, show that

(a) $[\hat{\mathbf{J}}^2, \hat{\mathbf{J}}_1^2] = [\hat{\mathbf{J}}^2, \hat{\mathbf{J}}_2^2] = 0$,

(b) $[\hat{\mathbf{J}}^2, \hat{\mathbf{J}}_1] = 2i\hbar(\hat{\mathbf{J}}_1 \times \hat{\mathbf{J}}_2)$ and $[\hat{\mathbf{J}}^2, \hat{\mathbf{J}}_2] = 2i\hbar(\hat{\mathbf{J}}_2 \times \hat{\mathbf{J}}_1)$, so that $[\hat{\mathbf{J}}^2, \hat{\mathbf{J}}] = 0$,

(c) $[\hat{\mathbf{J}}^2, \hat{J}_{1z}] \neq 0$, $[\hat{\mathbf{J}}^2, \hat{J}_{2z}] \neq 0$, but $[\hat{\mathbf{J}}^2, \hat{J}_z] = 0$,

(d) $[\hat{\mathbf{J}}_1 \cdot \hat{\mathbf{J}}_2, \hat{\mathbf{J}}_1] = i\hbar(\hat{\mathbf{J}}_1 \times \hat{\mathbf{J}}_2)$ and $[\hat{\mathbf{J}}_1 \cdot \hat{\mathbf{J}}_2, \hat{\mathbf{J}}_2] = i\hbar(\hat{\mathbf{J}}_2 \times \hat{\mathbf{J}}_1)$, so that $[\hat{\mathbf{J}}_1 \cdot \hat{\mathbf{J}}_2, \hat{\mathbf{J}}] = 0$.

2. Limits on j : Starting from $j_{\max} = j_1 + j_2$ and

$$\sum_{j=j_{\min}}^{j_{\max}} (2j+1) = (2j_1+1)(2j_2+1),$$

prove that $j_{\min} = |j_1 - j_2|$.

3. Spins of mesons and baryons: As you may know, mesons (such as pions or kaons) and baryons (such as protons or neutrons) consist of two and three quarks, respectively. The quarks are known to carry spin- $\frac{1}{2}$. Assume that the quarks are in their ground state so that their orbital angular momentum is zero.

(a) What are the possible spins for the mesons?

(b) What are the possible spins for the baryons?

4. Determining the Clebsch-Gordan coefficients: Given that $(j_1 = 1, j_2 = 1, j = 1, m = -1)$, express the state $|j_1 j_2; j m\rangle$ in terms of the states $|j_1 j_2; m_1 m_2\rangle = |j_1 m_1\rangle \otimes |j_2 m_2\rangle$.

5. Using the Clebsch-Gordan coefficients: Consider two spin-1 particles that occupy the state

$$|s_1 s_2; m_1 m_2\rangle = |1 1; 1 0\rangle.$$

(a) What is the probability of finding the system in an eigenstate of the total spin $\hat{\mathbf{S}}^2$ with quantum number $s = 1$?

(b) What is the probability for $s = 2$?

Additional exercises I

Angular momentum

1. Hamiltonian and angular momentum of a free particle: Consider a free particle in three dimensions whose position is described in terms of the spherical polar coordinates (r, θ, ϕ) . Let the conjugate momenta corresponding to these coordinates be p_r , p_θ and p_ϕ . Show that the Hamiltonian of the free particle can be expressed as

$$H = \frac{p_r^2}{2m} + \frac{L^2}{2m r^2},$$

where $\mathbf{L} = \mathbf{r} \times \mathbf{p}$ is the angular momentum of the particle. Also, show that

$$L^2 = p_\theta^2 + \frac{p_\phi^2}{\sin^2 \theta}.$$

2. An operator relation involving angular momentum: Show that

$$\hat{\mathbf{L}}^2 = \hat{\mathbf{x}}^2 \hat{\mathbf{p}}^2 - (\hat{\mathbf{x}} \cdot \hat{\mathbf{p}})^2 + i \hbar \hat{\mathbf{x}} \cdot \hat{\mathbf{p}}.$$

3. A model of angular momentum: Let

$$\hat{L}_\pm = (\hat{a}_\pm^\dagger \hat{a}_\mp) \hbar, \quad \hat{L}_z = (\hat{a}_+^\dagger \hat{a}_+ - \hat{a}_-^\dagger \hat{a}_-) \hbar/2 \quad \text{and} \quad \hat{N} = \hat{a}_+^\dagger \hat{a}_+ + \hat{a}_-^\dagger \hat{a}_-,$$

where \hat{a}_\pm and \hat{a}_\pm^\dagger are the annihilation and the creation operators of two *independent* simple harmonic oscillators satisfying the usual commutation relations. Show that

- (a) $[\hat{L}_z, \hat{L}_\pm] = \pm \hbar \hat{L}_\pm$,
 (b) $\hat{\mathbf{L}}^2 = \hat{N} \left[(\hat{N}/2) + 1 \right] (\hbar^2/2)$,
 (c) $[\hat{\mathbf{L}}^2, \hat{L}_z] = 0$.

Note: This representation of the angular momentum operators in terms of creation and the annihilation operators of oscillators is known as the Schwinger model.

4. Particle in a central potential: A particle in a spherically symmetric potential is known to be in an eigenstate of $\hat{\mathbf{L}}^2$ and \hat{L}_z with eigenvalues $l(l+1)\hbar^2$ and $m\hbar$, respectively. Prove that the expectation values in the $|l, m\rangle$ states satisfy

$$\langle \hat{L}_x \rangle = \langle \hat{L}_y \rangle = 0, \quad \langle \hat{L}_x^2 \rangle = \langle \hat{L}_y^2 \rangle = \frac{\hbar^2}{2} [l(l+1) - m^2].$$

5. Hamiltonian involving angular momentum: The Hamiltonian of a system is described in terms of the angular momentum operators as follows:

$$\hat{H} = \frac{\hat{L}_x^2}{2I_1} + \frac{\hat{L}_y^2}{2I_2} + \frac{\hat{L}_z^2}{2I_3}.$$

- (a) What are the eigen values of the Hamiltonian when $I_1 = I_2$?
 (b) What are the eigen values of the Hamiltonian if the angular momentum of the system is unity and $I_1 \neq I_2$?
6. Rotation of a spin- $\frac{1}{2}$ system: Consider a sequence of Euler rotations represented by

$$\mathcal{D}^{(1/2)}(\alpha, \beta, \gamma) = \exp - \left(\frac{i \sigma_3 \alpha}{2} \right) \exp - \left(\frac{i \sigma_2 \beta}{2} \right) \exp - \left(\frac{i \sigma_3 \gamma}{2} \right).$$

Since the rotations form a group, we expect the above sequence of rotations to be equivalent to a single rotation by an angle, say, θ , about some axis. Find the angle θ .

7. Rotating an eigen state of orbital angular momentum: Consider the orbital angular momentum eigenstate $|l = 2, m = 0\rangle$. If this state is rotated by an angle β about the y -axis, determine the probability for the new state to be found in $m = (0, \pm 1, \pm 2)$.
8. Measuring components of angular momentum: A beam of particles is subject to the simultaneous measurement of angular momentum variables L^2 and L_z . The measurement yields the pairs of values $(l = 0, m = 0)$ and $(l = 1, m = -1)$ with probabilities $3/4$ and $1/4$, respectively.
- Reconstruct the state of the beam immediately before the measurement.
 - The particles in the beam with $(l = 1, m = -1)$ are separated out and subjected to the measurement of L_x . What are the possible outcomes and their probabilities?
 - Construct the spatial wavefunctions of the states that could arise from second measurement.
9. (a) Adding spins: A particle of spin-1 and a particle of spin-2 are at rest in a configuration such that the total spin is 3, and its z -component is 1 (i.e. the eigenvalue of \hat{S}_z is \hbar). If you measured the z -component of the angular momentum of the spin-2 particle, what values might you get, and what is the probability of each one?
- (b) Adding orbital and spin angular momenta: An electron with spin down is in the state ψ_{510} of the hydrogen atom. If you could measure the total angular momentum squared of the electron alone (not including the proton spin), what values might you get, and what is the probability of each?
10. Interacting spins: A system of two particles with spins $s_1 = \frac{3}{2}$ and $s_2 = \frac{1}{2}$ is described by the Hamiltonian

$$\hat{H} = \alpha \hat{\mathbf{S}}_1 \cdot \hat{\mathbf{S}}_2,$$

with α being a given constant. The system is initially (say, at $t = 0$) in the following eigenstate of $\hat{\mathbf{S}}_1^2$, $\hat{\mathbf{S}}_2^2$, \hat{S}_{1z} and \hat{S}_{2z} :

$$|s_1 s_2; m_1 m_2\rangle = \left| \frac{3}{2} \frac{1}{2}; \frac{1}{2} \frac{1}{2} \right\rangle.$$

- Find the state of the system at times $t > 0$.
- What is the probability of finding the system in the state $\left| \frac{3}{2} \frac{3}{2}; \frac{1}{2} - \frac{1}{2} \right\rangle$?

Exercise sheet 5

Systems of identical particles – Concepts

1. *Exchange forces:* Consider a pair of free identical particles of mass m . For simplicity, let us assume that they are moving in one dimension. Also, let us ignore their spin. The particles are localized around the points $+a$ and $-a$ and are described by the following real wave functions:

$$\psi_{\pm}(x) = \left(\frac{\beta}{\pi}\right)^{1/4} \exp - [\beta (x \mp a)^2].$$

Evidently, a well-localized state corresponds to $\beta \gg a^{-2}$.

- (a) Write down the wave function of the system for the cases wherein the identical particles are bosons or fermions.
 - (b) Calculate the expectation value of the energy and show that, if the two particles are fermions, then there is an effective repulsion between them.
 - (c) Compare the result with the case of two identical bosons.
2. *Wavefunctions of three bosons or fermions:* Let $\psi_{f_i}(\xi)$ be the wavefunctions of single particle states normalized to unity, where f_i are the quantum numbers of some complete set.
- (a) Write down the normalized wavefunctions for states of a system consisting of three identical weakly interacting bosons which can occupy single particle states with given quantum numbers f_1, f_2 and f_3 .
 - (b) Also, write down the corresponding wavefunctions if the particles are fermions.
3. *Slater determinant:* Consider a solution of the Schrodinger equation of the following form:

$$\hat{H} \psi_E(x_1, x_2, \dots, x_N) = E \psi_E(x_1, x_2, \dots, x_N),$$

which describes N particles. If the total Hamiltonian is given by

$$\hat{H} = \hat{H}_1 + \hat{H}_2 + \dots + \hat{H}_N,$$

where \hat{H}_i denotes the Hamiltonian of the i -th particle, then we can write

$$\psi_E(x_1, x_2, \dots, x_N) = \psi_{E_1}(x_1) \psi_{E_2}(x_2) \dots \psi_{E_N}(x_N),$$

with $E = E_1 + E_2 + \dots + E_N$. Prove, say, by induction, that a completely antisymmetric wavefunction describing the system of N particles can be written as a determinant as follows:

$$\psi(x_1, x_2, \dots, x_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \psi_{E_1}(x_1) & \psi_{E_1}(x_2) & \psi_{E_1}(x_3) & \dots & \psi_{E_1}(x_N) \\ \psi_{E_2}(x_1) & \psi_{E_2}(x_2) & \psi_{E_2}(x_3) & \dots & \psi_{E_2}(x_N) \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \psi_{E_N}(x_1) & \psi_{E_N}(x_2) & \psi_{E_N}(x_3) & \dots & \psi_{E_N}(x_N) \end{vmatrix}.$$

Note: The above determinant is often referred to as the Slater determinant.

4. *Ground state energy of a pair of interacting electrons:* Consider a pair of electrons constrained to move in one dimension in the total spin $s = 1$ state. The electrons interact through the following attractive potential:

$$V(x) = \begin{cases} 0 & \text{for } |x_1 - x_2| > a, \\ -V_0 & \text{for } |x_1 - x_2| \leq a, \end{cases}$$

where $V_0 > 0$. Find the lowest energy eigenvalue in the case where the total momentum of the two electrons vanishes.

5. Energy of the ground state of Helium: Earlier, when we had discussed the energy of the ground state of the Helium atom, we had ignored the interaction between the two electrons.
- (a) Calculate the energy of the ground state of the system when the interaction between the two electrons is accounted for. How does the result compare with the experimentally observed value?
 - (b) Can you account for the difference?
-

Quiz II

Angular momentum

1. Commutation relations involving angular momentum: Evaluate or establish the following commutation relations involving angular momentum:

(a) $[\hat{L}_i, \hat{x}_j]$, 2 marks

(b) $\hat{\mathbf{L}} \times \hat{\mathbf{r}} + \hat{\mathbf{r}} \times \hat{\mathbf{L}} = 2i\hbar \hat{\mathbf{r}}$, 3 marks

(c) $[\hat{\mathbf{L}}^2, \hat{\mathbf{r}}] = -2i\hbar \hat{\mathbf{\Theta}}$, where $\hat{\mathbf{\Theta}} = \hat{\mathbf{L}} \times \hat{\mathbf{r}} - i\hbar \hat{\mathbf{r}}$ 3 marks

(d) $[\hat{L}_x^2, \hat{L}_y^2] = [\hat{L}_y^2, \hat{L}_z^2] = [\hat{L}_z^2, \hat{L}_x^2]$. 2 marks

Note: $\epsilon_{ijk} \epsilon_{ljk} = 2\delta_{il}$.

2. Rotation matrix for $j = 1$: Consider a system with $j = 1$.

(a) Construct the matrix describing \hat{J}_y . 4 marks

(b) Using the matrix representation, show that $\hat{J}_y^3 = \hat{J}_y$. 2 marks

(c) As a result, when $j = 1$, show that we can write 4 marks

$$\exp - \left(\frac{i \hat{J}_y \alpha}{\hbar} \right) = 1 - i \sin \alpha \frac{\hat{J}_y}{\hbar} - (1 - \cos \alpha) \frac{\hat{J}_y^2}{\hbar^2}.$$

3. (a) Energy of a spin- $\frac{3}{2}$ particle: A spin- $\frac{3}{2}$ particle is described by the Hamiltonian

$$\hat{H} = \frac{\alpha}{\hbar^2} \left(\hat{S}_x^2 + \hat{S}_y^2 - 2\hat{S}_z^2 \right) - \frac{\beta}{\hbar} \hat{S}_z,$$

where α and β are real constants. Determine the energy eigen values of the particle. 5 marks

- (b) Particle in a central potential: A particle in a spherically symmetric potential is known to be in the eigenstate $|l, m\rangle$ of the operators $\hat{\mathbf{L}}^2$ and \hat{L}_z with eigenvalues $l(l+1)\hbar^2$ and $m\hbar$, respectively.

i. Evaluate $\langle \hat{L}_x \rangle$, $\langle \hat{L}_y \rangle$ and $\langle \hat{L}_z \rangle$ in the state $|l, m\rangle$. 2 marks

ii. Similarly, evaluate $\langle \hat{L}_x^2 \rangle$, $\langle \hat{L}_y^2 \rangle$ and $\langle \hat{L}_z^2 \rangle$ in $|l, m\rangle$. 3 marks

4. Expectation value in the singlet state: Suppose two spin- $\frac{1}{2}$ particles are known to be in the singlet configuration. Let S_{1a} be the component of the spin angular momentum of the first particle in the direction defined by the unit vector $\hat{\mathbf{a}}$. Similarly, let S_{2b} be the component of second particle's spin angular momentum in the direction $\hat{\mathbf{b}}$. Evaluate the expectation value of the operator $\hat{S}_{1a} \hat{S}_{2b}$ in the singlet state. 10 marks

5. Addition of spin angular momentum: Consider a system composed of two particles, one with spin- $\frac{1}{2}$ and another with spin-1.

(a) List all the allowed spin states of the composite system. 3 marks

(b) Express all the spin states of the composite system in terms of the spin states of the individual systems. 7 marks

Exercise sheet 6

Systems of identical particles – Applications

1. Fermi energy of copper: The density of copper is 8.96 gm/cm^3 , and its atomic weight is 63.5 gm/mole .
 - (a) Assume $q = 1$, calculate the Fermi energy for copper and express it in electron volts.
 - (b) What is the corresponding velocity of the electrons if one assumes that the Fermi energy can be written as $E_F = m v^2/2$? Is it safe to assume that the electrons in copper are nonrelativistic?
 - (c) At what temperature would the characteristic thermal energy (i.e. $k_B T$, where k_B is the Boltzmann constant and T is the temperature in Kelvin) be equal to the Fermi energy, for copper?
 Note: This is called the Fermi temperature. As long as the actual temperature is substantially below the Fermi temperature, the material can be regarded as ‘cold’, with most of the electrons in the ground-state configuration. Since the melting point of copper is 1356 K , solid copper is always cold.
 - (d) Calculate the degeneracy pressure of copper in the electron gas model.
2. White dwarfs: Certain cold stars (called white dwarfs) are stabilized against gravitational collapse by the degeneracy pressure of their electrons. Assuming constant density, the radius R of such an object can be calculated as follows:
 - (a) Write the total electron energy in terms of the radius R , the number of nucleons (protons and neutrons) N , the number of electrons per nucleon q , and the mass of the electron m_e .
 - (b) Calculate the gravitational energy of a uniformly dense sphere. Express the result in terms of G (the constant of gravitation), R , N and the mass of a nucleon, say, m_n .
 Note: The gravitational energy will be negative.
 - (c) Find the radius for which the sum of the electron and the gravitational energies is a minimum.
 - (d) Express the radius R in terms of the number of particles N .
 - (e) Determine the radius, in kilometers, of a white dwarf with the mass of the sun.
 - (f) Determine the Fermi energy, in electron volts, for the white dwarf and compare it with the rest energy of an electron. What do you notice?
3. Chandrasekhar limit: We can extend the theory of a free electron gas to the relativistic domain by replacing the classical kinetic energy $E = p^2/(2m)$, with the relativistic formula $E^2 = (p^2 c^2 + m^2 c^4) - m^2 c^4$. Note that the momentum is related to the wave vector in the usual way, viz. $\mathbf{p} = \hbar \mathbf{k}$. Specifically, in the extreme relativistic limit, we have $E \simeq pc = \hbar ck$.
 - (a) Replace $\hbar^2 k^2/(2m)$ in the energy associated with the electrons by the ultra-relativistic expression $\hbar ck$ and calculate the total energy of the relativistic electrons.
 - (b) Repeat the arguments of the previous exercise concerning the electron and gravitational energies for the case of the ultra-relativistic electron gas.
 - (c) In this case, one finds that, there is no stable minimum, regardless of the radius R . Find the critical number of nucleons N_c such that gravitational collapse occurs for $N > N_c$.
 Note: This is called the Chandrasekhar limit.
 - (d) Express the corresponding stellar mass in terms of the mass of the Sun.
4. Bose-Einstein condensation: Consider a collection of bosons at a finite temperature T . Let $\mu(T)$ be the chemical potential for the system.

- (a) Show that, for bosons, the chemical potential must always be less than the minimum allowed energy.
- (b) In particular, for the ideal Bose gas (identical bosons in the three-dimensional infinite square well), $\mu(T) < 0$ for all T . Show that, in this case, $\mu(T)$ monotonically increases as T decreases, assuming that the total number N and the volume V are held constant.
- (c) Bose condensation corresponds to a situation wherein all the particles crowd into the ground state. This occurs as we lower T and $\mu(T)$ hits zero. Evaluate the integral describing the total number of particles N for $\mu = 0$, and determine the critical temperature T_c at which this occurs.
- (d) Find the critical temperature for ${}^4\text{He}$.
Note: The density of ${}^4\text{He}$ at this temperature is 0.15 gm/cm^3 . Also, the experimental value of the critical temperature in ${}^4\text{He}$ is 2.17 K .

5. Wien displacement and Stefan-Boltzmann laws: Consider black body radiation at a temperature T .

- (a) Determine the energy density in the wavelength range $d\lambda$.
- (b) Use this expression to derive the following Wien displacement law for the wavelength λ_{max} at which the energy density of the black body radiation is a maximum:

$$\lambda_{\text{max}} = \frac{b}{T},$$

where the quantity b is known as the Wien's constant.

Note: You will encounter a transcendental equation, which you will need to solve numerically.

- (c) Determine the value of Wien's constant.
- (d) Derive the following Stefan-Boltzmann formula describing the total energy density of black-body radiation:

$$\rho = \frac{4\sigma}{c} T^4,$$

where σ is a constant known as the Stefan's constant.

- (e) Determine the value of Stefan's constant.
-

Exercise sheet 7
Theory of scattering – Essentials

1. Classical scattering by a hard sphere: Consider a small ball which is incident on a large sphere of radius, say, R , and bouncing off elastically.
 - (a) Express the impact parameter b of the incident ball in terms of the scattering angle θ .
 - (b) Evaluate the differential cross-section $D(\theta)$.
 - (c) What is the corresponding total cross-section?

2. Rutherford scattering in classical mechanics: An incident particle of charge q and kinetic energy E scatters off a heavy, stationary particle with charge Q .
 - (a) Derive the expression relating the impact parameter b to the scattering angle θ .
 - (b) Determine the differential scattering cross-section $D(\theta)$.
 - (c) Show that the corresponding total cross-section is infinite.

3. Rayleigh's formula: Establish the following Rayleigh's formula:

$$e^{ikz} = \sum_{l=0}^{\infty} i^l (2l+1) j_l(kr) P_l(\cos\theta),$$

where $z = r \cos\theta$, while $j_l(x)$ and $P_l(x)$ denote the spherical Bessel function and the Legendre polynomial, respectively.

4. Green's function for the Helmholtz equation: Derive the Green's function $G(\mathbf{r})$ associated with the following Helmholtz equation:

$$(\nabla^2 + k^2) G(\mathbf{r}) = \delta^{(3)}(\mathbf{r}),$$

where $\delta^{(3)}(\mathbf{r})$ denotes the three-dimensional Dirac delta function.

5. The Lippmann-Schwinger equation and the integral form of the Schrodinger equation: Consider the Hamiltonian of a system which can be expressed as $H = H_0 + V$, where H_0 is the Hamiltonian of the free particle. Let $|\psi_0\rangle$ denote the energy eigen ket of the free particle with energy E_0 , i.e.

$$\hat{H}_0 |\psi_0\rangle = E_0 |\psi_0\rangle.$$

Also, let $|\psi\rangle$ denote an energy eigen ket of the complete system with eigen value E satisfying the equation

$$(\hat{H}_0 + \hat{V}) |\psi\rangle = E |\psi\rangle.$$

- (a) Argue that we can rewrite the above equation as

$$|\psi_{\pm}\rangle = |\psi_0\rangle + \frac{1}{E - H_0 \pm i\epsilon} \hat{V} |\psi_{\pm}\rangle,$$

where ϵ is an infinitesimal, positive definite quantity.

Note: This equation is known as the Lippmann-Schwinger equation.

- (b) Show that, in the position representation, the above equation corresponds to the following integral form of the Schrodinger equation:

$$\psi(\mathbf{r}) = \psi_0(\mathbf{r}) - \frac{m}{2\pi\hbar^2} \int d^3\mathbf{r}' \frac{e^{-ik|\mathbf{r}-\mathbf{r}'|}}{|\mathbf{r}-\mathbf{r}'|} V(\mathbf{r}') \psi(\mathbf{r}'),$$

where $\psi_0(\mathbf{r}) = \langle \mathbf{r} | \psi_0 \rangle$ and $\psi(\mathbf{r}) = \langle \mathbf{r} | \psi \rangle$.

- (c) For appropriate $V(\mathbf{r})$ and E , explicitly check that the ground state of the hydrogen atom satisfies above the integral form of the Schrodinger equation.
-

Exercise sheet 8

Theory of scattering – Examples

1. Quantum scattering by a ‘hard’ sphere: Consider a ‘hard’ repulsive spherical potential of the following form:

$$V(r) = \begin{cases} \infty & \text{for } r \leq a, \\ 0 & \text{for } r > a. \end{cases}$$

Using the method of partial waves, determine the corresponding scattering amplitude and the total cross-section.

Note: This can be considered to be the quantum mechanical version of classical scattering by a hard sphere.

2. Scattering by a ‘soft’ sphere: Now, consider the case wherein the repulsive spherical potential is ‘softened’ to a finite value in the following manner:

$$V(r) = \begin{cases} V_0 & \text{for } r \leq a, \\ 0 & \text{for } r > a, \end{cases}$$

where $V_0 > 0$.

- (a) Working in the Born approximation, determine the scattering amplitude, the differential cross-section and the total cross-section for the above potential at low energies.
 - (b) What is the corresponding result for arbitrary energies? Does it reduce to the result obtained at low energies?
3. Yukawa and Rutherford scattering: Consider the following spherically symmetric Yukawa potential:

$$V(r) = \beta \frac{e^{-\mu r}}{r},$$

where β and μ are constants.

- (a) Working in the Born approximation, evaluate the corresponding scattering amplitude and differential cross-section.
 - (b) Determine the scattering amplitude and differential cross-section for the case wherein $\beta = Qq/(4\pi\epsilon_0)$ and $\mu = 0$, which corresponds to Rutherford scattering. How does the differential cross-section compare with the classical result?
4. Scattering by a spherical shell: Consider scattering by a spherical shell described by the following potential:

$$V(r) = \alpha \delta^{(1)}(r - r_0),$$

where α and r_0 are constants.

- (a) Using the method of partial waves, calculate the scattering amplitude at low energies.
- (b) Evaluate the corresponding differential and total cross-sections.
- (c) Using the Born approximation, calculate the scattering amplitude and the differential cross-section at low energies. Does the result match the one obtained using the method of partial waves?
- (d) Working in the Born approximation, calculate the scattering amplitude for arbitrary energies.

5. Beyond the first order Born approximation: Earlier, using the first order Born approximation, we had calculated the scattering amplitude at low energies in the following potential:

$$V(r) = \begin{cases} V_0 & \text{for } r \leq a, \\ 0 & \text{for } r > a, \end{cases}$$

where $V_0 > 0$. Evaluate the correction to the scattering amplitude due to the second order Born approximation.

Additional exercises II

From systems of identical particles to the theory of scattering

1. *Possible configurations of identical bosons:* Three identical bosons with spin $s = 1$ are in the same orbital states described by the wavefunction $\psi(\mathbf{r})$.
 - (a) Write down the normalized spin functions for the total system.
 - (b) How many independent states are possible?
 - (c) What are the possible values of the total spin of the system?
2. *Ortho and para states:* Two identical spin- $\frac{1}{2}$ fermions move in one dimension inside an infinite square well described by the potential

$$V(x) = \begin{cases} 0 & \text{for } 0 < x < L, \\ \infty & \text{otherwise.} \end{cases}$$

- (a) Write down the ground state wave function and the ground state energy when the two particles are in the triplet spin state, referred to as the ortho state.
 - (b) What is the ground state wave function and the corresponding energy when they are in the singlet spin state, called the para state?
 - (c) Let us now suppose that the two particles are interacting mutually through the short range attractive potential $V(x) = \lambda \delta(x_1 - x_2)$, with $\lambda > 0$. Use first order perturbation theory to calculate the corrections to the above energies due to this mutual interaction.
3. *Ground state and Fermi energies of fermions:* Consider N identical spin- $\frac{1}{2}$ particles which are subjected to the one-dimensional simple harmonic oscillator potential.
 - (a) What is the Fermi energy of the system?
 - (b) What is the energy of the ground state of the system?
 - (c) What are the ground state and Fermi energies if we ignore the mutual interactions and assume N to be very large?
4. *Neutron stars:* Stars that are heavier than the Chandrasekhar limit will not form white dwarfs, but they will collapse further, finally becoming neutron stars under certain conditions. This is due to the fact that, at very high density, inverse beta decay, i.e. $e + p \rightarrow n + \nu$, converts virtually all of the protons and electrons into neutrons, liberating neutrinos in the process, which carry off energy. Eventually, neutron degeneracy pressure stabilizes the collapse, just as electron degeneracy does for the white dwarfs.
 - (a) Calculate the radius of a neutron star with the mass of the sun.
 - (b) Also calculate the neutron Fermi energy, and compare it with the rest energy of a neutron. Is it reasonable to treat such a star non-relativistically?
5. *Distinguishable particles in a thermal bath:* Consider a collection of N distinguishable, non-interacting, quantum particles which are in thermal equilibrium at a finite temperature T . Assume that each of these particles experiences the simple harmonic potential.
 - (a) Working in one spatial dimension, evaluate the chemical potential and the total energy of the collection of particles.
 - (b) What are the energies in the quantum and classical limits, viz. when $k_B T \ll \hbar\omega$ and $k_B T \gg \hbar\omega$? Can you interpret the results?
 - (c) What are the corresponding results in three spatial dimensions?

6. Phase shift in one-dimensional scattering: A particle of mass m and energy E is incident from the left on the potential

$$V(x) = \begin{cases} 0 & \text{for } x < -a, \\ -V_0 & \text{for } -a \leq x \leq 0, \\ \infty & \text{for } x > 0, \end{cases}$$

where $V_0 > 0$. Consider an incoming wave $A e^{ikx}$, where $k = \sqrt{2mE}/\hbar$ with $E > 0$, which is reflected by the above potential.

- Find the reflected wave.
 - What is the amplitude of the reflected wave? Does the reflection alter the amplitude of incident wave?
 - Assuming that the well is deep, i.e. $-V_0 \ll E$, determine the shift in the phase of the reflected wave.
7. Partial waves and phase shifts: Recall that, using the method of partial waves, we had obtained the scattering amplitude to be

$$f(\theta) = \sum_{l=0}^{\infty} (2l+1) a_l P_l(\cos \theta),$$

where $P_l(x)$ denotes the Legendre polynomials and a_l was called the l -th partial wave amplitude.

- Using the above result, arrive at the corresponding expression for the differential cross-section and show that total cross-section can be written as

$$\sigma = \sum_{l=0}^{\infty} (2l+1) |a_l|^2.$$

- Focusing on a particular l , show that the partial wave amplitude a_l can be expressed in terms of the phase shift δ_l as follows:

$$a_l = \frac{1}{k} e^{i\delta_l} \sin(\delta_l).$$

- Use this form of a_l to arrive an expression for the total cross-section in terms of the phase shifts δ_l .
8. The optical theorem: According to the so-called optical theorem, the imaginary part of the scattering amplitude $f(\theta)$ along the forward direction is related to the total cross-section σ through the following relation:

$$\text{Im. } f(\theta = 0) = \frac{k \sigma}{4\pi}.$$

Use the results from the previous exercise to establish the theorem.

9. Scattering in one dimension: Consider scattering by a potential $V(x)$ in one spatial dimension. In such a case, we can write the time-independent Schrodinger equation as

$$\frac{d^2\psi}{dx^2} + k^2 \psi = Q,$$

where $k = \sqrt{2mE}/\hbar$ and $Q = (2m/\hbar^2) V \psi$.

- Obtain the Green's function for the above Schrodinger equation.
- Use it to construct the integral form of the equation.
- Develop the Born approximation for one-dimensional scattering and obtain an expression for the reflection coefficient in terms of the potential $V(x)$.

- (d) Use the Born approximation to determine the reflection and transmission coefficients for the cases of the delta function potential

$$V(x) = -\alpha \delta^{(1)}(x),$$

with $\alpha > 0$, and the finite square well potential of the following form:

$$V(x) = \begin{cases} -V_0 & \text{for } -a < x < a, \\ 0 & \text{for } -|x| > a, \end{cases}$$

with $V_0 > 0$.

- (e) Compare the results from the Born approximation with the exact ones.

10. Scattering by an electric dipole: Consider an electric dipole consisting of two electric charges e and $-e$ located at, say, $x = -a$ and $x = a$. A particle of charge e and mass m described by the wave vector $\mathbf{k} = k \hat{\mathbf{z}}$ is incident on the dipole. Calculate the scattering amplitude in the Born approximation. Find the directions at which the differential cross section is maximal.
-

Exercise sheet 9

The Dirac equation and solutions describing the free particle

1. Properties of the Dirac matrices: Recall that the four Dirac matrices α_i with $i = (1, 2, 3)$ and β are defined as

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} \mathcal{I} & 0 \\ 0 & -\mathcal{I} \end{pmatrix},$$

where σ_i and \mathcal{I} are the 2×2 Pauli matrices and the unit matrix given by

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \mathcal{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

- (a) Show that the Dirac matrices satisfy the following relations:

$$\alpha_i \alpha_j + \alpha_j \alpha_i = 2 \delta_{ij}, \quad \alpha_i \beta + \beta \alpha_i = 0, \quad \alpha_i^2 = \beta^2 = 1.$$

- (b) Show that the sum of the eigen values is zero for all the four matrices.
 (c) Determine the eigen values of all the Dirac matrices.

2. Pauli's Hamiltonian and the gyromagnetic ratio of the electron: Consider an electron in an external magnetic field \mathbf{B} . In non-relativistic quantum mechanics, the interaction of the magnetic moment associated with the spin of the electron \mathbf{S} and the magnetic field is described by the following Hamiltonian originally due to Pauli:

$$\hat{H} = -\frac{e g}{2 m c} \hat{\mathbf{S}} \cdot \mathbf{B} = -\frac{e g \hbar}{4 m c} \hat{\boldsymbol{\sigma}} \cdot \mathbf{B},$$

where $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ denotes the Pauli matrices, and the quantity g is known as the gyromagnetic ratio. Evidently, e and m denote the charge and the mass of the electron, while c represents the velocity of light.

Note: The gyromagnetic ratio g of a system is defined as the ratio of its magnetic moment and angular momentum.

- (a) Show that, in the absence of the magnetic field, the Hamiltonian of the free, non-relativistic electron can be expressed as

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2 m} = \frac{1}{2 m} (\boldsymbol{\sigma} \cdot \hat{\mathbf{p}})^2.$$

- (b) In the presence of an external magnetic field, we can replace the momentum operator $\hat{\mathbf{p}}$ by $\hat{\mathbf{p}} - (e/c) \mathbf{A}$, where \mathbf{A} is the vector potential describing the magnetic field. In such a case, the Hamiltonian of the electron can be written as

$$\hat{H} = \frac{1}{2 m} \left[\boldsymbol{\sigma} \cdot \left(\hat{\mathbf{p}} - \frac{e}{c} \mathbf{A} \right) \right]^2.$$

Using the properties of the Pauli matrices, show that the interaction of the electron's spin with the magnetic field is described by Pauli's Hamiltonian above with $g = 2$.

Note: For this reason, the gyromagnetic ratio of the electron is said to be 2.

3. Properties of the gamma matrices: Recall that the gamma matrices $\gamma^\mu = (\gamma^0, \gamma^i)$ are defined as

$$\gamma^0 = \beta, \quad \gamma^i = \beta \alpha_i.$$

Show that the relations for the matrices α_i and β (listed in the previous exercise) can be expressed as

$$\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu = 2 \eta^{\mu\nu} \mathcal{I},$$

where $\eta^{\mu\nu}$ is the metric tensor describing the Minkowski spacetime and \mathcal{I} represents the 4×4 unit matrix.

4. Spinors describing particles and anti-particles: Recall that the solutions to the Dirac equation which governs relativistic free particles can be expressed as

$$\psi(\mathbf{r}, t) = A e^{-(i/\hbar)(Et - \mathbf{p}\cdot\mathbf{r})} u(E, \mathbf{p}),$$

where E and $\mathbf{p} = (p_x, p_y, p_z)$ represent the energy and momentum of the particle, while $u(E, \mathbf{p})$ denotes a bispinor and A is a constant. The spinors $u(E, \mathbf{p})$ that describe particles are given by

$$u^{(1)}(E, \mathbf{p}) = N \begin{pmatrix} 1 \\ 0 \\ \frac{c p_z}{E+m c^2} \\ \frac{c(p_x + i p_y)}{E+m c^2} \end{pmatrix}, \quad u^{(2)}(E, \mathbf{p}) = N \begin{pmatrix} 0 \\ 1 \\ \frac{c(p_x - i p_y)}{E+m c^2} \\ \frac{-c p_z}{E+m c^2} \end{pmatrix},$$

where $E = \sqrt{p^2 c^2 + m^2 c^4}$ and N is the normalization constant. The spinors $v(E, \mathbf{p})$ that describe anti-particles are given by

$$v^{(1)}(E, \mathbf{p}) = N \begin{pmatrix} \frac{c(p_x - i p_y)}{E+m c^2} \\ \frac{-c p_z}{E+m c^2} \\ 0 \\ 1 \end{pmatrix}, \quad v^{(2)}(E, \mathbf{p}) = -N \begin{pmatrix} \frac{c p_z}{E+m c^2} \\ \frac{c(p_x + i p_y)}{E+m c^2} \\ 1 \\ 0 \end{pmatrix}.$$

- (a) Show that, in $u^{(1)}$ and $u^{(2)}$, the bottom two components (i.e. u_B) are smaller than the top two components (i.e. u_A) by the factor of v/c in the non-relativistic limit.
- (b) The scalar product between two spinors, say, u and v , is defined as $u^\dagger v$. Determine the normalization constant N for all the above spinors by setting $u^\dagger u = 2E/c$.
- (c) Show that $u^{(1)}$ and $u^{(2)}$ are orthogonal. Similarly, show that $v^{(1)}$ and $v^{(2)}$ are also orthogonal.
- (d) Are $u^{(1)}$ and $u^{(2)}$ orthogonal to $v^{(1)}$ and $v^{(2)}$?
5. From the Dirac equation to the Klein-Gordon equation: Recall that, in its covariant form, the Dirac equation can be expressed as

$$(i \hbar \gamma^\mu \partial_\mu - m c) \psi = 0,$$

where $\partial_\mu = \partial/\partial x^\mu$ and ψ represents a spinor. By applying the operator $\gamma^\nu \partial_\nu$ on the above equation, show that all the components of the Dirac spinor satisfy the following Klein-Gordon equation:

$$\left(\partial^\mu \partial_\mu + \frac{m^2 c^2}{\hbar^2} \right) \psi = \left(\square + \frac{m^2 c^2}{\hbar^2} \right) \psi = 0.$$

Exercise sheet 10

Spin and helicity of particles described by the Dirac equation

1. Eigen spinors of \hat{S}_z : Recall that the spin operator describing the Dirac particles is given by

$$\hat{S} = \frac{\hbar}{2} \Sigma,$$

where

$$\Sigma = \begin{pmatrix} \boldsymbol{\sigma} & 0 \\ 0 & \boldsymbol{\sigma} \end{pmatrix}.$$

We had seen earlier that the spinors $\{u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)}\}$ (listed in previous exercise sheet) are *not* eigen spinors of \hat{S}_z operator. Show that, when $p_x = p_y = 0$ so that $p_z = p$, these spinors prove to be eigen spinors of \hat{S}_z . What are the corresponding eigen values?

2. Spin of particles described by the Dirac field: Show that every bispinor is an eigen state of the operator \hat{S}^2 . Determine the corresponding eigen value and identify the spin of the particle described by the Dirac equation.
3. Conservation of total angular momentum: Consider a relativistic particle described by the Dirac equation.
- What is the Hamiltonian \hat{H} of the relativistic particle?
 - Evaluate the commutator of \hat{H} with the orbital angular momentum $\hat{L} = \hat{\mathbf{r}} \times \hat{\mathbf{p}}$.
 - Also, calculate the commutator of \hat{H} with the spin angular momentum $\hat{S} = (\hbar/2) \hat{\Sigma}$.
 - Show that, while the orbital and spin angular momenta \hat{L} and \hat{S} are not conserved individually, the total angular momentum $\hat{J} = \hat{L} + \hat{S}$ is conserved.
4. Spinors with definite helicity: Construct normalized spinors, say, $u^{(+)}$ and $u^{(-)}$, representing an electron of momentum \mathbf{p} with helicity ± 1 . In other words, obtain solutions to $(\gamma^\mu p_\mu - mc) u = 0$, with positive energy E , which are eigen spinors of the helicity operator $\mathbf{p} \cdot \hat{\Sigma}$ with eigenvalues ± 1 .
5. Massless Dirac particles and helicity: Consider the case of massless particles described by the Dirac equation.

- Writing the Dirac bispinor ψ in terms of the spinors describing the particles and anti-particles, say, ϕ and χ , arrive at the equations governing these spinors in Fourier space (or, equivalently, in the energy-momentum space).
- Let us define two new spinors (which are linear combinations of the original spinors ϕ and χ) as follows:

$$\psi_\pm = \frac{1}{\sqrt{2}} (\chi \pm \phi).$$

Arrive at the equations describing the spinors ψ_\pm .

- Show that the ψ_\pm spinors describe particles and anti-particles with definite helicity.

Note: Unlike electrons and positrons which can be described by states with different helicities (as we saw in the previous exercise), neutrinos and anti-neutrinos have definite helicities. While neutrinos are left handed (i.e. they have helicity -1), anti-neutrinos are right handed (i.e. they have helicity $+1$). Therefore, they can be described by spinors such as ψ_- and ψ_+ , respectively.

End-of-semester exam I

From basics of quantum mechanics to relativistic quantum mechanics

1. Decomposing a second rank square matrix: An arbitrary second rank square matrix, say, A , can be expanded in terms of the unit matrix \mathcal{I} and the Pauli matrices $\boldsymbol{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$ as follows:

$$A = a_0 \mathcal{I} + a_x \sigma_x + a_y \sigma_y + a_z \sigma_z = a_0 + \mathbf{a} \cdot \boldsymbol{\sigma}.$$

Prove that the coefficients are given by $2a_0 = \text{Tr}. A$ and $2\mathbf{a} = \text{Tr}. (\boldsymbol{\sigma} A)$.

10 marks

2. An operator relation involving angular momentum: Show that

10 marks

$$\hat{\mathbf{L}}^2 = \hat{\mathbf{x}}^2 \hat{\mathbf{p}}^2 - (\hat{\mathbf{x}} \cdot \hat{\mathbf{p}})^2 + i \hbar \hat{\mathbf{x}} \cdot \hat{\mathbf{p}}.$$

3. Interacting spins: A system of two particles with spins $s_1 = \frac{3}{2}$ and $s_2 = \frac{1}{2}$ is described by the Hamiltonian

$$\hat{H} = \alpha \hat{\mathbf{S}}_1 \cdot \hat{\mathbf{S}}_2,$$

with α being a given constant. The system is initially (say, at $t = 0$) in the following eigenstate of $\hat{\mathbf{S}}_1^2$, $\hat{\mathbf{S}}_2^2$, $\hat{\mathbf{S}}_{1z}$ and $\hat{\mathbf{S}}_{2z}$:

$$|s_1 s_2; m_1 m_2\rangle = \left| \frac{3}{2} \frac{1}{2}; \frac{1}{2} \frac{1}{2} \right\rangle.$$

- (a) Find the state of the system at times $t > 0$.

7 marks

- (b) What is the probability of finding the system in the state $\left| \frac{3}{2} \frac{1}{2}; \frac{3}{2} - \frac{1}{2} \right\rangle$?

3 marks

4. Scattering amplitude and cross-section in the Born approximation: Using the Born approximation, find the scattering amplitude and the total scattering cross-section of a particle in the following two central potentials: $V(r) = \alpha e^{-\mu r}$ and $V(r) = \alpha/r^2$.

6+4 marks

5. Non-relativistic and relativistic spin- $\frac{1}{2}$ particles: Consider non-relativistic and relativistic spin- $\frac{1}{2}$ particles.

- (a) What is the spin operator $\hat{\mathbf{S}}$ describing a non-relativistic spin- $\frac{1}{2}$ particle? What are the eigen values of the corresponding $\hat{\mathbf{S}}^2$ operator? Show that *all* spinors are eigen functions of the $\hat{\mathbf{S}}^2$ operator?

1+1+1 marks

- (b) What is the spin operator $\hat{\mathbf{S}}$ describing a *relativistic* spin- $\frac{1}{2}$ particle? What are the eigen values of the corresponding $\hat{\mathbf{S}}^2$ operator? Show that *all bispinors* are eigen functions of the $\hat{\mathbf{S}}^2$ operator?

1+3+3 marks

Assignment in lieu of Quiz III

From systems of identical particles to relativistic quantum mechanics

1. Possible configurations of identical bosons: Three identical bosons with spin $s = 1$ are in the same orbital states described by the wavefunction $\psi(\mathbf{r})$.
 - (a) Write down the normalized spin functions for the total system.
 - (b) How many independent states are possible?
 - (c) What are the possible values of the total spin of the system?
2. Neutron stars: Stars that are heavier than the Chandrasekhar limit will not form white dwarfs, but they will collapse further, finally becoming neutron stars under certain conditions. This is due to the fact that, at very high density, inverse beta decay, i.e. $e + p \rightarrow n + \nu$, converts virtually all of the protons and electrons into neutrons, liberating neutrinos in the process, which carry off energy. Eventually, the degeneracy pressure of non-relativistic neutrons stabilizes the collapse, just as electron degeneracy does for the white dwarfs.
 - (a) Calculate the radius of a neutron star with the mass of the sun.
 - (b) Also calculate the neutron Fermi energy, and compare it with the rest energy of a neutron. Is it reasonable to treat such a star non-relativistically?
3. Partial waves and phase shifts: Recall that, using the method of partial waves, we had obtained the scattering amplitude to be

$$f(\theta) = \sum_{l=0}^{\infty} (2l+1) a_l P_l(\cos \theta),$$

where $P_l(x)$ denotes the Legendre polynomials and a_l was called the l -th partial wave amplitude.

- (a) Using the above result, arrive at the corresponding expression for the differential cross-section and show that total cross-section can be written as

$$\sigma = \sum_{l=0}^{\infty} (2l+1) |a_l|^2.$$

- (b) Focusing on a particular l , show that the partial wave amplitude a_l can be expressed in terms of the phase shift δ_l as follows:

$$a_l = \frac{1}{k} e^{i\delta_l} \sin(\delta_l).$$

- (c) Use this form of a_l to arrive an expression for the total cross-section in terms of the phase shifts δ_l .

4. Scattering amplitude and cross-section in the Born approximation: Using the Born approximation, find the scattering amplitude *as well as* the total scattering cross-section of a particle in the following two central potentials: $V(r) = \alpha e^{-\mu r}$ and $V(r) = \alpha/r^2$.
5. Simultaneous diagonalization of $\hat{\mathbf{p}}$, \hat{H} and $\hat{\Sigma}$ describing a Dirac particle: Recall that the Hamiltonian \hat{H} of a Dirac particle is given by

$$\hat{H} = \boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m c^2,$$

where $\hat{\mathbf{p}}$ is the momentum operator. The quantities α_i with $i = (1, 2, 3)$ and β are the Dirac matrices defined as

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} \mathcal{I} & 0 \\ 0 & -\mathcal{I} \end{pmatrix},$$

where σ_i and \mathcal{I} are the 2×2 Pauli matrices and the unit matrix given by

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \mathcal{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

Also, note that the spin operator describing the Dirac particles can be expressed in terms of the Pauli matrices $\boldsymbol{\sigma}$ as follows:

$$\hat{S} = \frac{\hbar}{2} \boldsymbol{\Sigma},$$

where

$$\boldsymbol{\Sigma} = \begin{pmatrix} \boldsymbol{\sigma} & 0 \\ 0 & \boldsymbol{\sigma} \end{pmatrix}.$$

Show that the commutators $[\hat{\boldsymbol{p}}, \hat{H}]$, $[\hat{\boldsymbol{p}}, \hat{\boldsymbol{p}} \cdot \hat{\boldsymbol{\Sigma}}]$ and $[\hat{H}, \hat{\boldsymbol{p}} \cdot \hat{\boldsymbol{\Sigma}}]$ vanish.

End-of-semester exam II

From basics of quantum mechanics to relativistic quantum mechanics

1. Uncertainties in angular momentum: Recall that, according to the generalized uncertainty principle, the uncertainty associated with two observables represented by the operators \hat{A} and \hat{B} is given by

$$\Delta A^2 \Delta B^2 \geq \left(\frac{1}{2i} \langle [\hat{A}, \hat{B}] \rangle \right)^2,$$

where the uncertainty, say, ΔA , is defined through the relation $\Delta A^2 = \langle \hat{A}^2 \rangle - \langle \hat{A} \rangle^2$ and the angular brackets as usual represent expectation values evaluated in a given state.

- (a) Evaluate the uncertainties ΔJ_x and ΔJ_y in the state $|j, m\rangle$. 5 marks
- (b) Show that the products of the uncertainties ΔJ_x and ΔJ_y is consistent with the generalized uncertainty principle. 5 marks
2. Relations involving angular momentum operators: Establish the following operator relations:

(a) $[\hat{\phi}, \hat{L}_z] = i\hbar$, 1 marks

(b) $\exp(i \hat{L}_z \alpha / \hbar) \psi(\phi) = \psi(\phi + \alpha)$, 2 marks

(c) $[\hat{L}^2, [\hat{L}^2, \mathbf{r}]] = 2\hbar^2 (\mathbf{r} \hat{L}^2 + \hat{L}^2 \mathbf{r})$. 7 marks

3. States and energies of distinguishable and identical particles: Consider a system of three non-interacting particles that are confined to a one-dimensional box with its walls at $x = 0$ and $x = L$. As usual, the potential $V(x)$ is assumed to be zero inside the box and infinity outside. Determine the energies and the wavefunctions of the ground state, the first and the second excited states of the system, when the three particles are:

(a) spinless and distinguishable particles with masses $m_1 < m_2 < m_3$, 4 marks

(b) identical bosons. 6 marks

4. Differential and total cross-sections in the Born approximation: Using the Born approximation, find the differential and the total scattering cross-sections of a particle in the following two central potentials: 5+5 marks

$$V(r) = V_0 \left(\frac{a}{r} \right) e^{-r/a} \quad \text{and} \quad V(r) = V_0 e^{-r^2/a^2}.$$

5. Conserved ‘probability’ and current densities: Recall that the Klein-Gordon equation is given by

$$\frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2} - \nabla^2 \Psi + \frac{m^2 c^2}{\hbar^2} \Psi = 0,$$

where Ψ is a complex scalar quantity. In contrast, the Dirac equation is given by

$$i\hbar \frac{\partial \Psi}{\partial t} = \frac{\hbar c}{i} (\boldsymbol{\alpha} \cdot \nabla) \Psi + \beta m c^2 \Psi,$$

where $\boldsymbol{\alpha}$ and β are the Dirac matrices and Ψ represents a four component, complex spinor.

- (a) Construct the conserved ‘probability’ density, say, $\rho(t, \mathbf{x})$, associated with the Klein-Gordon and Dirac equations. 5 marks
- (b) What are the associated current densities, say, $\mathbf{j}(t, \mathbf{x})$? 5 marks

End-of-semester exam III

From basics of quantum mechanics to relativistic quantum mechanics

1. States, measurements and uncertainties: Consider a system that is initially in the state

$$\psi(\theta, \phi) = \frac{1}{\sqrt{5}} Y_{1-1}(\theta, \phi) + \sqrt{\frac{3}{5}} Y_{10}(\theta, \phi) + \frac{1}{\sqrt{5}} Y_{11}(\theta, \phi),$$

where $Y_{lm}(\theta, \phi)$ are the spherical harmonics.

- (a) Evaluate $\langle \hat{L}_+ \rangle$ in the above state. 3 marks
- (b) If \hat{L}_z is measured, what are the values one would obtain and what are the corresponding probabilities? 2 marks
- (c) If, after measuring \hat{L}_z , one obtains $L_z = -\hbar$, calculate the corresponding uncertainties ΔL_x and ΔL_y as well as their product $\Delta L_x \Delta L_y$. 5 marks

2. An operator relation involving angular momentum: Show that 10 marks

$$\hat{L}^2 = \hat{x}^2 \hat{p}^2 - (\hat{x} \cdot \hat{p})^2 + i \hbar \hat{x} \cdot \hat{p}.$$

3. Interacting spins: A system of two particles with spins $s_1 = \frac{3}{2}$ and $s_2 = \frac{1}{2}$ is described by the Hamiltonian

$$\hat{H} = \alpha \hat{S}_1 \cdot \hat{S}_2,$$

with α being a given constant. The system is initially (say, at $t = 0$) in the following eigenstate of \hat{S}_1^2 , \hat{S}_2^2 , \hat{S}_{1z} and \hat{S}_{2z} :

$$|s_1 s_2; m_1 m_2\rangle = \left| \frac{3}{2} \frac{1}{2}; \frac{1}{2} \frac{1}{2} \right\rangle.$$

- (a) Find the state of the system at times $t > 0$. 7 marks
- (b) What is the probability of finding the system in the state $\left| \frac{3}{2} \frac{1}{2}; \frac{3}{2} - \frac{1}{2} \right\rangle$? 3 marks
4. States and energies of non-interacting bosons and fermions: Consider a system of N non-interacting, identical particles that are confined to a one-dimensional box with its walls at $x = 0$ and $x = L$. As usual, the potential $V(x)$ is assumed to be zero inside the box and infinity outside. Determine the energy and the wavefunction of the ground state, when the particles are:

- (a) identical bosons, 3 marks
- (b) identical spin- $\frac{1}{2}$ fermions. 7 marks

5. Non-relativistic and relativistic spin- $\frac{1}{2}$ particles: Consider non-relativistic and relativistic spin- $\frac{1}{2}$ particles.

- (a) What is the spin operator \hat{S} describing a non-relativistic spin- $\frac{1}{2}$ particle? What are the eigen values of the corresponding \hat{S}^2 operator? Show that *all* spinors are eigen functions of the \hat{S}^2 operator? 1+1+1 marks
- (b) What is the spin operator \hat{S} describing a *relativistic* spin- $\frac{1}{2}$ particle? What are the eigen values of the corresponding \hat{S}^2 operator? Show that *all bispinors* are eigen functions of the \hat{S}^2 operator? 1+3+3 marks